

Acoustics

The Science of Sound

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Acoustics



Acoustics (from Greek ακουστικός pronounced akoustikos meaning "of or for hearing, ready to hear") is the science that studies sound, in particular its production, transmission, and effects. The science of acoustics has many applications which are dependent upon the nature of the sound that is to be produced, transmitted or controlled. In the case of a desirable sound, such as music, the main application of acoustics is to make the music sound as good as possible. In the case of an undesirable sound, such as traffic noise, the main application of acoustics is in noise reduction. Another major area of acoustics is in the field of ultrasound which has applications in detection, such as sonar systems or non-destructive material testing. The articles in this Wikibook describe the fundamentals of acoustics and some of the major applications.

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Fundamentals

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Fundamentals

Fundamentals of Acoustics



Introduction

Sound is an oscillation of pressure transmitted through a gas, liquid, or solid in the form of a traveling wave, and can be generated by any localized pressure variation in a medium. An easy way to understand how sound propagates is to consider that space can be divided into thin layers. The vibration (the successive compression and relaxation) of these layers, at a certain velocity, enables the sound to propagate, hence producing a wave. The speed of sound depends on the compressibility and density of the medium.

Sound can be more readily understood as this: Kinetic Energy with nowhere to go. The cracking sound of a baseball connecting with a bat is the end result of the imperfect retention of the kinetic energy between the ball and the bat.

In this chapter, we will only consider the propagation of sound waves in an area without any acoustic source, in a homogeneous fluid.

Equation of waves

Sound waves consist in the propagation of a scalar quantity, acoustic over-pressure. The propagation of sound waves in a stationary medium (e.g. still air or water) is governed by the following equation (see wave equation):

$$\nabla^2 p - \frac{1}{c_0^2} \frac{\partial^2 p}{\partial t^2} = 0$$

This equation is obtained using the conservation equations (mass, momentum and energy) and the thermodynamic equations of state of an ideal gas (or of an ideally compressible solid or liquid), supposing that the pressure variations are small, and neglecting viscosity and thermal conduction, which would give other terms, accounting for sound attenuation.

In the propagation equation of sound waves, c_0 is the propagation velocity of the sound wave (which has nothing to do with the vibration velocity of the air layers). This propagation velocity has the following expression:

$$c_0 = \frac{1}{\sqrt{\rho_0 \chi_s}}$$

where ρ_0 is the density and χ_s is the compressibility coefficient of the propagation medium.

Helmholtz equation

Since the velocity field \underline{v} for acoustic waves is irrotational we can define an acoustic potential Φ by:

$$\underline{v} = \text{grad } \Phi$$

Using the propagation equation of the previous paragraph, it is easy to obtain the new equation:

$$\nabla^2 \Phi - \frac{1}{c_0^2} \frac{\partial^2 \Phi}{\partial t^2} = 0$$

Applying the Fourier Transform, we get the widely used Helmholtz equation:

$$\nabla^2 \hat{\Phi} + k^2 \hat{\Phi} = 0$$

where k is the wave number associated with Φ . Using this equation is often the easiest way to solve acoustical problems.

Acoustic intensity and decibel

The acoustic intensity represents the acoustic energy flux associated with the wave propagation:

$$\underline{i}(t) = p \underline{v}$$

We can then define the average intensity:

$$\underline{I} = \langle \underline{i} \rangle$$

However, acoustic intensity does not give a good idea of the sound level, since the sensitivity of our ears is logarithmic. Therefore we define decibels, either using acoustic over-pressure or acoustic average intensity:

$$p^{dB} = 20 \log \left(\frac{p}{p_{\text{ref}}} \right); L_I = 10 \log \left(\frac{I}{I_{\text{ref}}} \right)$$

where $p_{\text{ref}} = 2 \cdot 10^{-5} \text{ Pa}$ for air, or $p_{\text{ref}} = 10^{-6} \text{ Pa}$ for any other media, and $I_{\text{ref}} = 10^{-12} \text{ W/m}^2$.

Solving the wave equation

Plane waves

If we study the propagation of a sound wave, far from the acoustic source, it can be considered as a plane 1D wave.

If the direction of propagation is along the x axis, the solution is:

$$\Phi(x, t) = f \left(t - \frac{x}{c_0} \right) + g \left(t + \frac{x}{c_0} \right)$$

where f and g can be any function. f describes the wave motion toward increasing x, whereas g describes the motion toward decreasing x.

The momentum equation provides a relation between p and \underline{v} which leads to the expression of the specific impedance, defined as follows:

$$\frac{p}{v} = Z = \pm \rho_0 c_0$$

And still in the case of a plane wave, we get the following expression for the acoustic intensity:

$$\underline{i} = \pm \frac{p^2}{\rho_0 c_0} \underline{e}_x$$

Spherical waves

More generally, the waves propagate in any direction and are spherical waves. In these cases, the solution for the acoustic potential Φ is:

$$\Phi(r, t) = \frac{1}{r} f\left(t - \frac{r}{c_0}\right) + \frac{1}{r} g\left(t + \frac{r}{c_0}\right)$$

The fact that the potential decreases linearly while the distance to the source rises is just a consequence of the conservation of energy. For spherical waves, we can also easily calculate the specific impedance as well as the acoustic intensity.

Boundary conditions

Concerning the boundary conditions which are used for solving the wave equation, we can distinguish two situations. If the medium is not absorptive, the boundary conditions are established using the usual equations for mechanics. But in the situation of an absorptive material, it is simpler to use the concept of acoustic impedance.

Non-absorptive material

In that case, we get explicit boundary conditions either on stresses and on velocities at the interface. These conditions depend on whether the media are solids, inviscid or viscous fluids.

Absorptive material

Here, we use the acoustic impedance as the boundary condition. This impedance, which is often given by experimental measurements depends on the material, the fluid and the frequency of the sound wave.

Fundamentals of Room Acoustics



Introduction

Three theories are used to understand room acoustics :

1. The modal theory
2. The geometric theory
3. The theory of Sabine

The modal theory

This theory comes from the homogeneous Helmholtz equation $\nabla^2 \hat{\Phi} + k^2 \hat{\Phi} = 0$. Considering a simple geometry of a parallelepiped (L_1, L_2, L_3), the solution of this problem is with separated variables :

$$P(x, y, z) = X(x)Y(y)Z(z)$$

Hence each function X, Y and Z has this form :

$$X(x) = Ae^{-ikx} + Be^{ikx}$$

With the boundary condition $\frac{\partial P}{\partial x} = 0$, for $x = 0$ and $x = L_1$ (idem in the other directions), the expression of pressure is :

$$P(x, y, z) = C \cos\left(\frac{m\pi x}{L_1}\right) \cos\left(\frac{n\pi y}{L_2}\right) \cos\left(\frac{p\pi z}{L_3}\right)$$

$$k^2 = \left(\frac{m\pi}{L_1}\right)^2 + \left(\frac{n\pi}{L_2}\right)^2 + \left(\frac{p\pi}{L_3}\right)^2$$

where m, n, p are whole numbers

It is a three-dimensional stationary wave. Acoustic modes appear with their modal frequencies and their modal forms. With a non-homogeneous problem, a problem with an acoustic source Q in τ_0 , the final pressure in τ is the sum of the contribution of all the modes described above.

The modal density $\frac{dN}{df}$ is the number of modal frequencies contained in a range of 1Hz. It depends on the frequency f , the volume of the room V and the speed of sound c_0 :

$$\frac{dN}{df} \simeq \frac{4\pi V}{c_0^2} f^2$$

The modal density depends on the square frequency, so it increase rapidly with the frequency. At a certain level of frequency, the modes are not distinguished and the modal theory is no longer relevant.

The geometry theory

For rooms of high volume or with a complex geometry, the theory of acoustical geometry is critical and can be applied. The waves are modelised with rays carrying acoustical energy. This energy decrease with the reflection of the rays on the walls of the room. The reason of this phenomenon is the absorption of the walls.

The problem is this theory needs a very high power of calculation and that is why the theory of Sabine is often chosen because it is easier.

The theory of Sabine

Description of the theory

This theory uses the hypothesis of the diffuse field, the acoustical field is homogeneous and isotropic. In order to obtain this field, the room has to be enough reverberating and the frequencies have to be high enough to avoid the effects of predominating modes.

The variation of the acoustical energy E in the room can be written as :

$$\frac{dE}{dt} = W_s - W_{abs}$$

Where W_s and W_{abs} are respectively the power generated by the acoustical source and the power absorbed by the walls.

The power absorbed is related to the voluminal energy in the room e :

$$W_s = \frac{ec_0}{4}a$$

Where a is the equivalent absorption area defined by the sum of the product of the absorption coefficient and the area of each material in the room :

$$a = \sum_i \alpha_i S_i$$

The final equation is : $V \frac{de}{dt} = W_s - \frac{ec_0}{4}a$

The level of stationary energy is : $e_{sat} = 4 \frac{W_s}{ac_0}$

Reverberation time

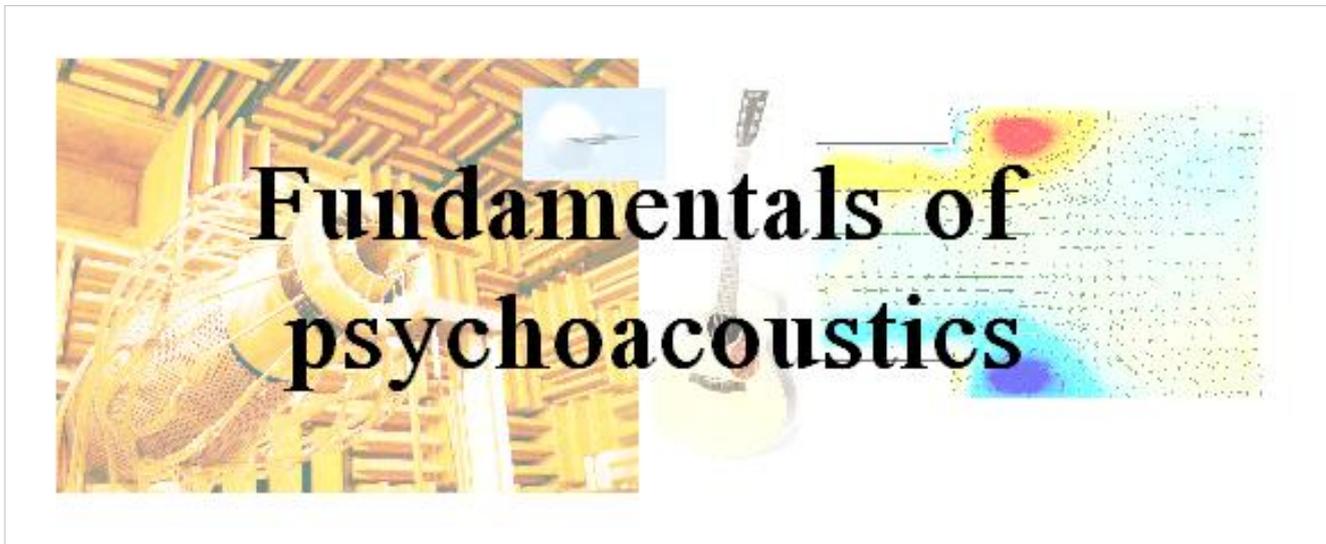
With this theory described, the reverberation time can be defined. It is the time for the level of energy to decrease of 60 dB. It depends on the volume of the room V and the equivalent absorption area a :

$$T_{60} = \frac{0.16V}{a} \text{ Sabine formula}$$

This reverberation time is the fundamental parameter in room acoustics and depends through the equivalent absorption area and the absorption coefficients on the frequency. It is used for several measurement :

- Measurement of an absorption coefficient of a material
- Measurement of the power of a source
- Measurement of the transmission of a wall

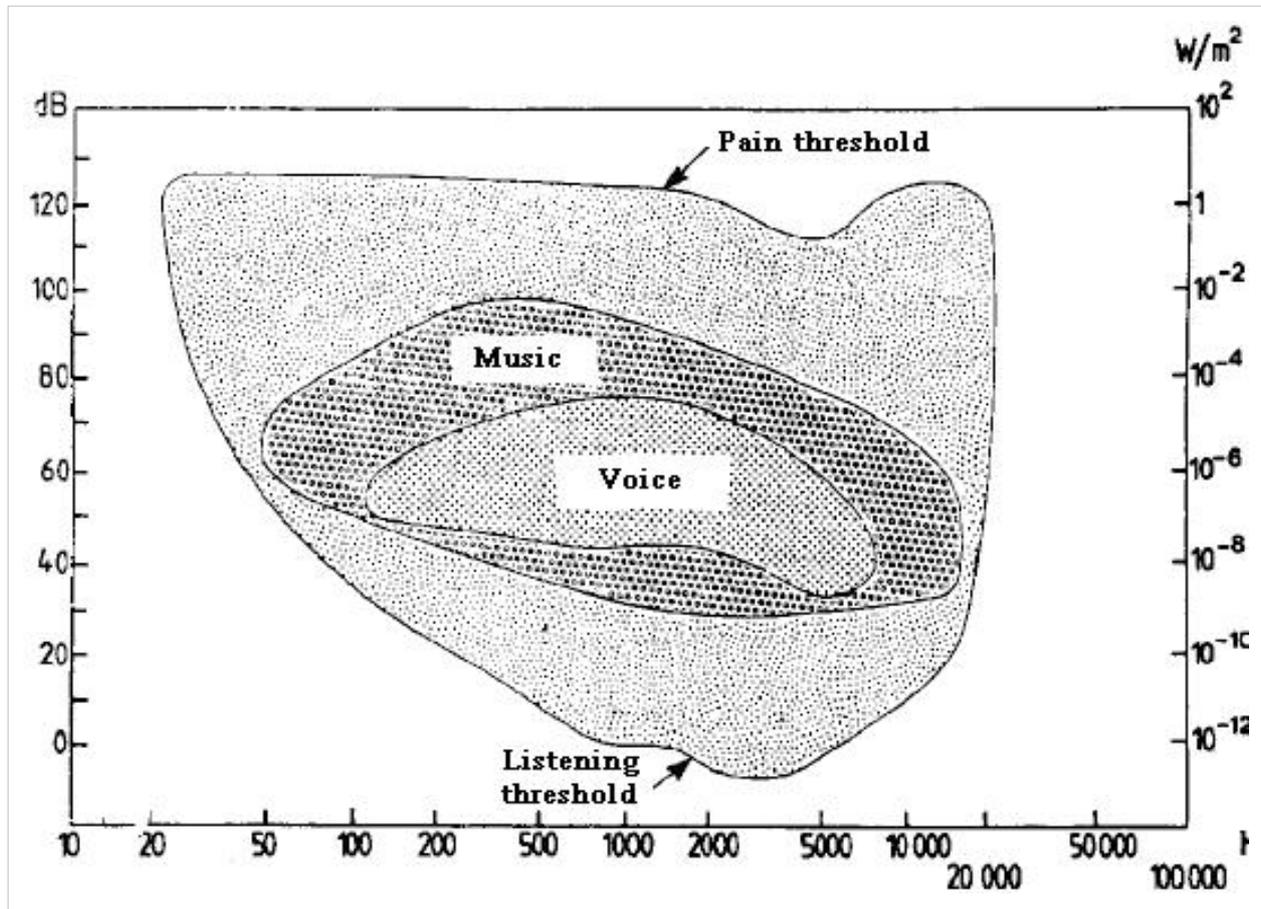
Fundamentals of Psychoacoustics



Due to the famous principle enounced by Gustav Theodor Fechner, the sensation of perception doesn't follow a linear law, but a logarithmic one. The perception of the intensity of light, or the sensation of weight, follow this law, as well. This observation legitimates the use of logarithmic scales in the field of acoustics. A 80dB (10^{-4} W/m²) sound seems to be twice as loud as a 70 dB (10^{-5} W/m²) sound, although there is a factor 10 between the two acoustic powers. This is quite a naïve law, but it led to a new way of thinking acoustics, by trying to describe the auditive sensations. That's the aim of psychoacoustics. By now, as the neurophysiologic mechanisms of human hearing haven't been successfully modelled, the only way of dealing with psychoacoustics is by finding metrics that best describe the different aspects of sound.

Perception of sound

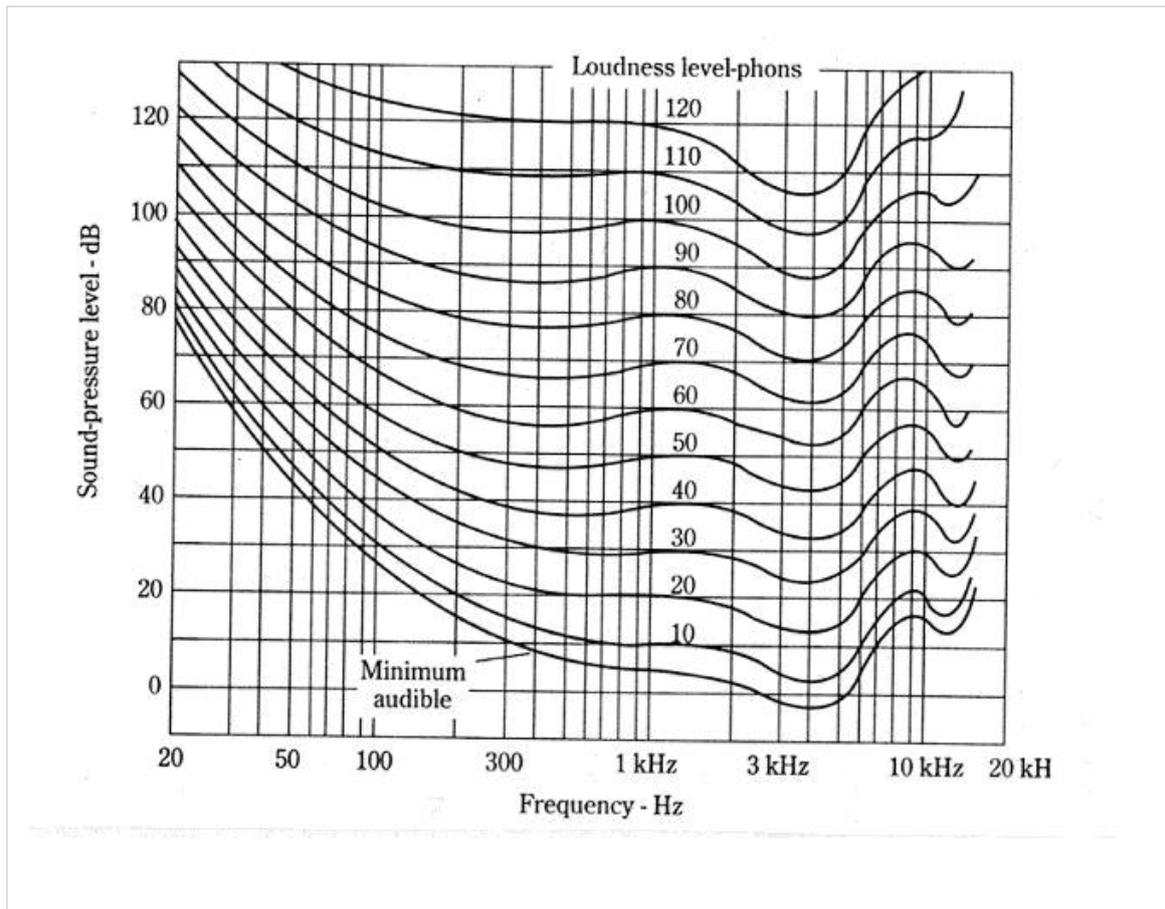
The study of sound perception is limited by the complexity of the human ear mechanisms. The figure below represents the domain of perception and the thresholds of pain and listening. The pain threshold is not frequency-dependent (around 120 dB in the audible bandwidth). At the opposite side, the listening threshold, as all the equal loudness curves, is frequency-dependent.



Phons and sones

Phons

Two sounds of equal intensity do not have the same loudness, because of the frequency sensibility of the human ear. A 80 dB sound at 100 Hz is not as loud as a 80 dB sound at 3 kHz. A new unit, the phon, is used to describe the loudness of a harmonic sound. X phons means “as loud as X dB at 1000 Hz”. Another tool is used : the equal loudness curves, a.k.a. Fletcher curves.



Sones

Another scale currently used is the sone, based upon the rule of thumb for loudness. This rule states that the sound must be increased in intensity by a factor 10 to be perceived as twice as loud. In decibel (or phon) scale, it corresponds to a 10 dB (or phons) increase. The sone scale's purpose is to translate those scales into a linear one.

$$\log(S) = 0,03(L_{ph} - 40)$$

Where S is the sone level, and L_{ph} the phon level. The conversion table is as follows:

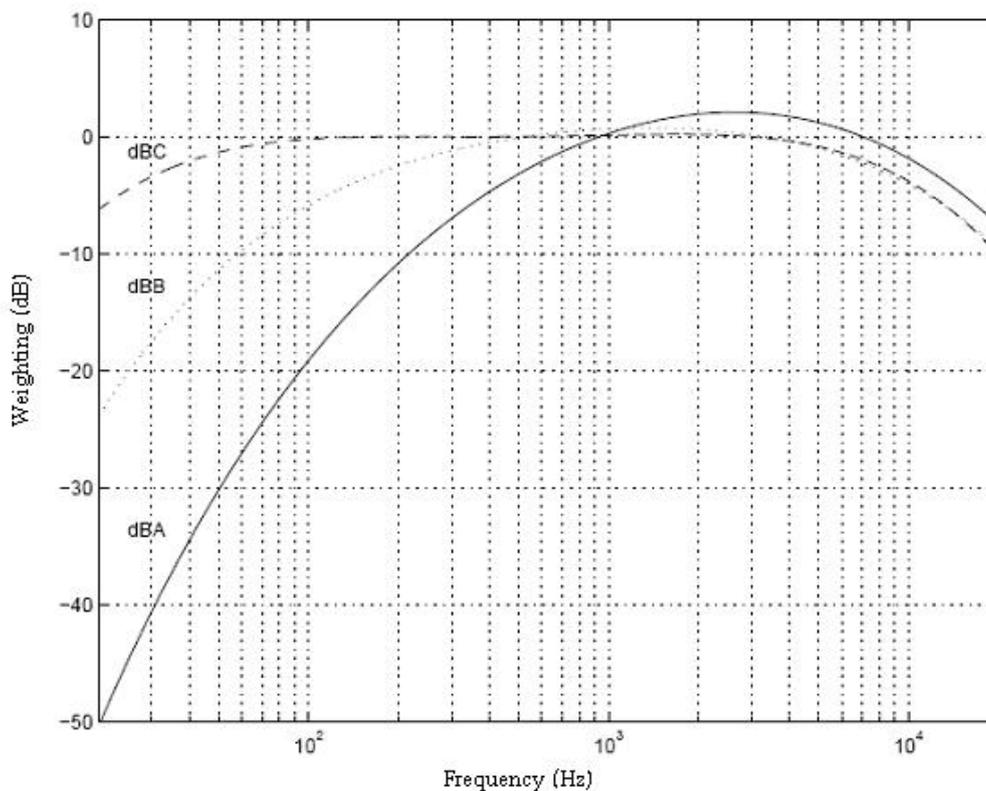
Phons	Sones
100	64
90	32
80	16
70	8
60	4
50	2
40	1

Metrics

We will now present five psychoacoustics parameters to provide a way to predict the subjective human sensation.

dB A

The measurement of noise perception with the sone or phon scale is not easy. A widely used measurement method is a weighting of the sound pressure level, according to frequency repartition. For each frequency of the density spectrum, a level correction is made. Different kinds of weightings (dB A, dB B, dB C) exist in order to approximate the human ear at different sound intensities, but the most commonly used is the dB A filter. Its curve is made to match the ear equal loudness curve for 40 phons, and as a consequence it's a good approximation of the phon scale.



Example : for a harmonic 40 dB sound, at 200 Hz, the correction is -10 dB, so this sound is 30 dB A.

Loudness

It measures the sound strength. Loudness can be measured in sone, and is a dominant metric in psychoacoustics.

Tonality

As the human ear is very sensible to the pure harmonic sounds, this metric is a very important one. It measures the number of pure tones in the noise spectrum. A broadband sound has a very low tonality, for example.

Roughness

It describes the human perception of temporal variations of sounds. This metric is measured in asper.

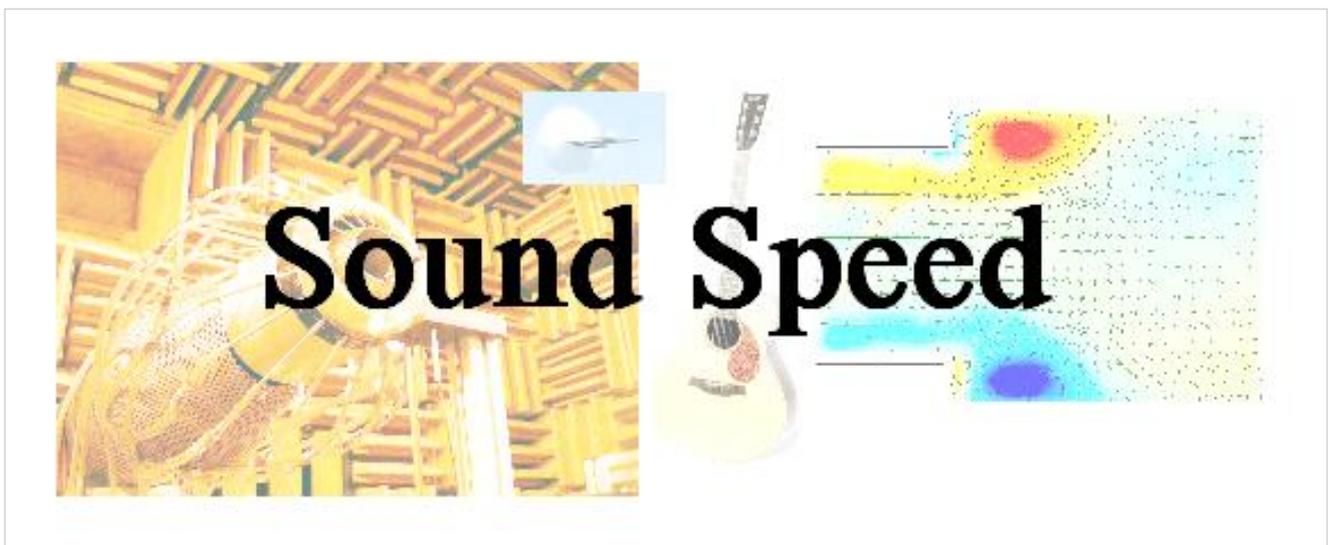
Sharpness

Sharpness is linked to the spectral characteristics of the sound. A high-frequency signal has a high value of sharpness. This metric is measured in acum.

Blocking effect

A sinusoidal sound can be masked by a white noise in a narrowing bandwidth. A white noise is a random signal with a flat power spectral density. In other words, the signal's power spectral density has equal power in any band, at any centre frequency, having a given bandwidth. If the intensity of the white noise is high enough, the sinusoidal sound will not be heard. For example, in a noisy environment (in the street, in a workshop), a great effort has to be made in order to distinguish someone's talking.

Sound Speed



The **speed of sound** c (from Latin *celeritas*, "velocity") varies depending on the medium through which the sound waves pass. It is usually quoted in describing properties of substances (e.g. see the article on sodium). In conventional use and in scientific literature sound velocity v is the same as sound speed c . Sound velocity c or velocity of sound should not be confused with sound particle velocity v , which is the velocity of the individual particles.

More commonly the term refers to the speed of sound in air. The speed varies depending on atmospheric conditions; the most important factor is the temperature. The humidity has very little effect on the speed of sound, while the

static sound pressure (air pressure) has none. Sound travels slower with an increased altitude (elevation if you are on solid earth), primarily as a result of temperature and humidity changes. An approximate speed (in metres per second) can be calculated from:

$$c_{\text{air}} = (331.5 + (0.6 \cdot \theta)) \text{ m/s}$$

where θ (theta) is the temperature in degrees Celsius.

Details

A more accurate expression for the speed of sound is

$$c = \sqrt{\kappa \cdot R \cdot T}$$

where

- R (287.05 J/(kg·K) for air) is the gas constant for air: the universal gas constant R , which units of J/(mol·K), is divided by the molar mass of air, as is common practice in aerodynamics)
- κ (kappa) is the adiabatic index (1.402 for air), sometimes noted γ
- T is the absolute temperature in kelvins.

In the standard atmosphere :

T_0 is 273.15 K (= 0 °C = 32 °F), giving a value of 331.5 m/s (= 1087.6 ft/s = 1193 km/h = 741.5 mph = 643.9 knots).

T_{20} is 293.15 K (= 20 °C = 68 °F), giving a value of 343.4 m/s (= 1126.6 ft/s = 1236 km/h = 768.2 mph = 667.1 knots).

T_{25} is 298.15 K (= 25 °C = 77 °F), giving a value of 346.3 m/s (= 1136.2 ft/s = 1246 km/h = 774.7 mph = 672.7 knots).

In fact, assuming an ideal gas, the speed of sound c depends on temperature only, **not on the pressure**. Air is almost an ideal gas. The temperature of the air varies with altitude, giving the following variations in the speed of sound using the standard atmosphere - *actual conditions may vary*. Any qualification of the speed of sound being "at sea level" is also irrelevant. Speed of sound varies with altitude (height) only because of the changing temperature!

Altitude	Temperature	m/s	km/h	mph	knots
Sea level (?)	15 °C (59 °F)	340	1225	761	661
11,000 m–20,000 m (Cruising altitude of commercial jets, and first supersonic flight)	-57 °C (-70 °F)	295	1062	660	573
29,000 m (Flight of X-43A)	-48 °C (-53 °F)	301	1083	673	585

In a **Non-Dispersive Medium** – Sound speed is independent of frequency, therefore the speed of energy transport and sound propagation are the same. For audio sound range air is a non-dispersive medium. We should also note that air contains CO2 which is a dispersive medium and it introduces dispersion to air at ultrasound frequencies (28kHz).

In a **Dispersive Medium** – Sound speed is a function of frequency. The spatial and temporal distribution of a propagating disturbance will continually change. Each frequency component propagates at each its own phase speed, while the energy of the disturbance propagates at the group velocity. Water is an example of a dispersive medium.

In general, the speed of sound c is given by

$$c = \sqrt{\frac{C}{\rho}}$$

where

C is a coefficient of stiffness

ρ is the density

Thus the speed of sound increases with the stiffness of the material, and decreases with the density.

In a fluid the only non-zero stiffness is to volumetric deformation (a fluid does not sustain shear forces).

Hence the speed of sound in a fluid is given by

$$c = \sqrt{\frac{K}{\rho}}$$

where

K is the adiabatic bulk modulus

For a gas, K is approximately given by

$$K = \kappa \cdot p$$

where

κ is the adiabatic index, sometimes called γ .

p is the pressure.

Thus, for a gas the speed of sound can be calculated using:

$$c = \sqrt{\frac{\kappa \cdot p}{\rho}}$$

which using the ideal gas law is identical to:

$$c = \sqrt{\kappa \cdot R \cdot T}$$

(Newton famously considered the speed of sound before most of the development of thermodynamics and so incorrectly used isothermal calculations instead of adiabatic. His result was missing the factor of κ but was otherwise correct.)

In a solid, there is a non-zero stiffness both for volumetric and shear deformations. Hence, in a solid it is possible to generate sound waves with different velocities dependent on the deformation mode.

In a solid rod (with thickness much smaller than the wavelength) the speed of sound is given by:

$$c = \sqrt{\frac{E}{\rho}}$$

where

E is Young's modulus

ρ (rho) is density

Thus in steel the speed of sound is approximately 5100 m/s.

In a solid with lateral dimensions much larger than the wavelength, the sound velocity is higher. It is found by replacing Young's modulus with the plane wave modulus, which can be expressed in terms of the Young's modulus and Poisson's ratio as:

$$M = E \frac{1 - \nu}{1 - \nu - 2\nu^2}$$

For air, see density of air.

The speed of sound in water is of interest to those mapping the ocean floor. In saltwater, sound travels at about 1500 m/s and in freshwater 1435 m/s. These speeds vary due to pressure, depth, temperature, salinity and other factors.

For general equations of state, if classical mechanics is used, the speed of sound c is given by

$$c^2 = \frac{\partial p}{\partial \rho}$$

where differentiation is taken with respect to adiabatic change.

If relativistic effects are important, the speed of sound S is given by:

$$S^2 = c^2 \left. \frac{\partial p}{\partial e} \right|_{\text{adiabatic}}$$

(Note that $e = \rho(c^2 + e^G)$ is the relativistic internal energy density).

This formula differs from the classical case in that ρ has been replaced by e/c^2 .

Speed of sound in air

Impact of temperature			
θ in °C	c in m/s	ρ in kg/m ³	Z in N·s/m ³
-10	325.4	1.341	436.5
-5	328.5	1.316	432.4
0	331.5	1.293	428.3
+5	334.5	1.269	424.5
+10	337.5	1.247	420.7
+15	340.5	1.225	417.0
+20	343.4	1.204	413.5
+25	346.3	1.184	410.0
+30	349.2	1.164	406.6

Mach number is the ratio of the object's speed to the speed of sound in air (medium).

Sound in solids

In solids, the velocity of sound depends on *density* of the material, not its temperature. Solid materials, such as steel, conduct sound much faster than air.

Experimental methods

In air a range of different methods exist for the measurement of sound.

Single-shot timing methods

The simplest concept is the measurement made using two microphones and a fast recording device such as a digital storage scope. This method uses the following idea.

If a sound source and two microphones are arranged in a straight line, with the sound source at one end, then the following can be measured:

1. The distance between the microphones (x)
2. The time delay between the signal reaching the different microphones (t)

Then $v = x/t$

An older method is to create a sound at one end of a field with an object that can be seen to move when it creates the sound. When the observer sees the sound-creating device act they start a stopwatch and when the observer hears the sound they stop their stopwatch. Again using $v = x/t$ you can calculate the speed of sound. A separation of at least 200 m between the two experimental parties is required for good results with this method.

Other methods

In these methods the time measurement has been replaced by a measurement of the inverse of time (frequency).

Kundt's tube is an example of an experiment which can be used to measure the speed of sound in a small volume, it has the advantage of being able to measure the speed of sound in any gas. This method uses a powder to make the nodes and antinodes visible to the human eye. This is an example of a compact experimental setup.

A tuning fork can be held near the mouth of a long pipe which is dipping into a barrel of water, in this system it is the case that the pipe can be brought to resonance if the length of the air column in the pipe is equal to $(2n+1)\lambda/4$ where n is an integer. As the antinodal point for the pipe at the open end is slightly outside the mouth of the pipe it is best to find two or more points of resonance and then measure half a wavelength between these.

Here it is the case that $v = f\lambda$

External links

- Calculation: Speed of sound in air and the temperature ^[1]
- The speed of sound, the temperature, and ... **not** the air pressure ^[2]
- Properties Of The U.S. Standard Atmosphere 1976 ^[3]

References

[1] <http://www.sengpielaudio.com/calculator-speedsound.htm>

[2] <http://www.sengpielaudio.com/SpeedOfSoundPressure.pdf>

[3] <http://www.pdas.com/atmos.htm>

Filter Design and Implementation



Introduction

Acoustic filters, or mufflers, are used in a number of applications requiring the suppression or attenuation of sound. Although the idea might not be familiar to many people, acoustic mufflers make everyday life much more pleasant. Many common appliances, such as refrigerators and air conditioners, use acoustic mufflers to produce a minimal working noise. The application of acoustic mufflers is mostly directed to machine components or areas where there is a large amount of radiated sound such as high pressure exhaust pipes, gas turbines, and rotary pumps.

Although there are a number of applications for acoustic mufflers, there are really only two main types which are used. These are absorptive and reactive mufflers. Absorptive mufflers incorporate sound absorbing materials to attenuate the radiated energy in gas flow. Reactive mufflers use a series of complex passages to maximize sound attenuation while meeting set specifications, such as pressure drop, volume flow, etc. Many of the more complex mufflers today incorporate both methods to optimize sound attenuation and provide realistic specifications.

In order to fully understand how acoustic filters attenuate radiated sound, it is first necessary to briefly cover some basic background topics. For more information on wave theory and other material necessary to study acoustic filters please refer to the references below.

Basic wave theory

Although not fundamentally difficult to understand, there are a number of alternate techniques used to analyze wave motion which could seem overwhelming to a novice at first. Therefore, only 1-D wave motion will be analyzed to keep most of the mathematics as simple as possible. This analysis is valid, with not much error, for the majority of pipes and enclosures encountered in practice.

Plane-wave pressure distribution in pipes

The most important equation used is the wave equation in 1-D form (See [1],[2], 1-D Wave Equation ^[1], for information).

Therefore, it is reasonable to suggest, if plane waves are propagating, that the pressure distribution in a pipe is given by:

$$\mathbf{p} = \mathbf{P_i}e^{j[\omega t - kx]} + \mathbf{P_r}e^{j[\omega t + kx]}$$

where P_i and P_r are incident and reflected wave amplitudes respectively. Also note that bold notation is used to indicate the possibility of complex terms. The first term represents a wave travelling in the +x direction and the second term, -x direction.

Since acoustic filters or mufflers typically attenuate the radiated sound power as much as possible, it is logical to assume that if we can find a way to maximize the ratio between reflected and incident wave amplitude then we will effectively attenuated the radiated noise at certain frequencies. This ratio is called the reflection coefficient and is given by:

$$\mathbf{R} = \left(\frac{\mathbf{P_r}}{\mathbf{P_i}} \right)$$

It is important to point out that wave reflection only occurs when the impedance of a pipe changes. It is possible to match the end impedance of a pipe with the characteristic impedance of a pipe to get no wave reflection. For more information see [1] or [2].

Although the reflection coefficient isn't very useful in its current form since we want a relation describing sound power, a more useful form can be derived by recognizing that the power intensity coefficient is simply the magnitude of reflection coefficient square [1]:

$$R_\pi = |\mathbf{R}|^2$$

As one would expect, the power reflection coefficient must be less than or equal to one. Therefore, it is useful to define the transmission coefficient as:

$$T_\pi = (1 - R_\pi)$$

which is the amount of power transmitted. This relation comes directly from conservation of energy. When talking about the performance of mufflers, typically the power transmission coefficient is specified.

Basic filter design

For simple filters, a long wavelength approximation can be made to make the analysis of the system easier. When this assumption is valid (e.g. low frequencies) the components of the system behave as lumped acoustical elements. Equations relating the various properties are easily derived under these circumstances.

The following derivations assume long wavelength. Practical applications for most conditions are given later.

Low-pass filter

These are devices that attenuate the radiated sound power at higher frequencies. This means the power transmission coefficient is approximately 1 across the band pass at low frequencies (see figure to right).

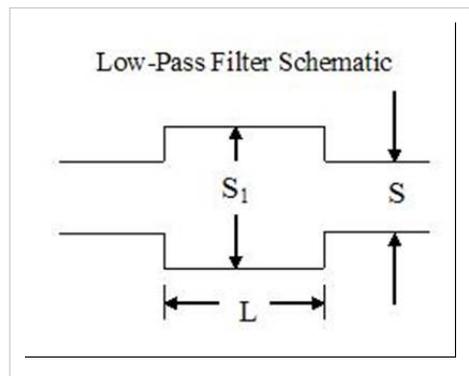
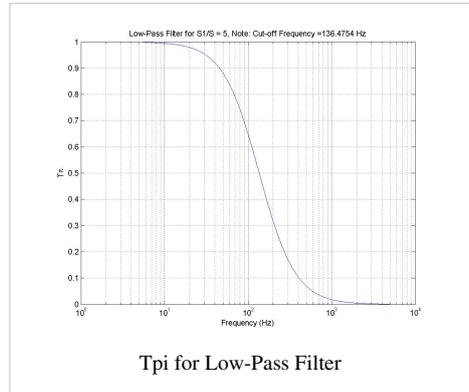
This is equivalent to an expansion in a pipe, with the volume of gas located in the expansion having an acoustic compliance (see figure to right). Continuity of acoustic impedance (see Java Applet at: Acoustic Impedance Visualization [2]) at the junction, see [1], gives a power transmission coefficient of:

$$T_{\pi} = \left(\frac{1}{1 + \left(\frac{S_1 - S}{2S} \right) kL} \right)^2$$

where k is the wavenumber (see Wave Properties), L & S_1 are length and area of expansion respectively, and S is the area of the pipe.

The cut-off frequency is given by:

$$f_c = \left(\frac{cS}{\pi L(S_1 - S)} \right)$$

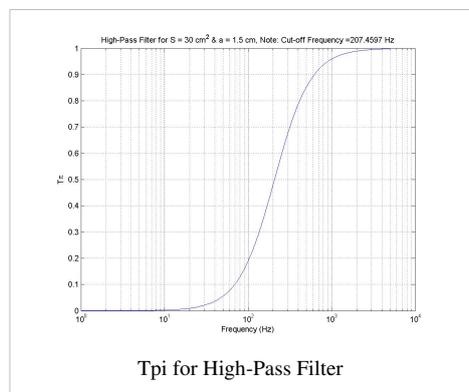


High-pass filter

These are devices that attenuate the radiated sound power at lower frequencies. Like before, this means the power transmission coefficient is approximately 1 across the band pass at high frequencies (see figure to right).

This is equivalent to a short side branch (see figure to right) with a radius and length much smaller than the wavelength (lumped element assumption). This side branch acts like an acoustic mass and applies a different acoustic impedance to the system than the low-pass filter. Again using continuity of acoustic impedance at the junction yields a power transmission coefficient of the form [1]:

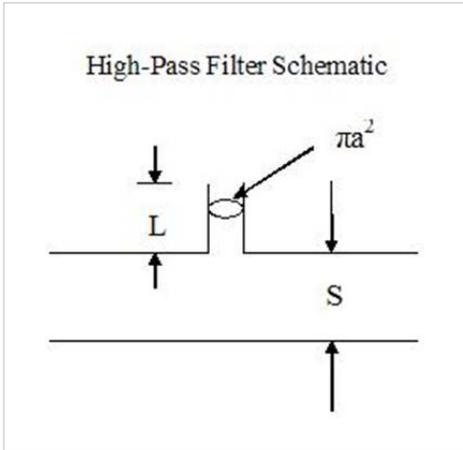
$$T_{\pi} = \left(\frac{1}{1 + \left(\frac{\pi a^2}{2SLk} \right)^2} \right)^2$$



where a and L are the area and effective length of the small tube, and S is the area of the pipe.

The cut-off frequency is given by:

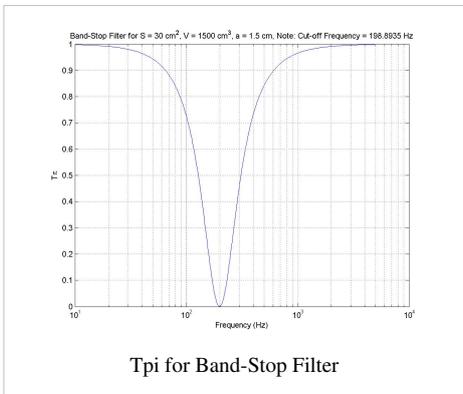
$$f_c = \left(\frac{ca^2}{2SL} \right)$$



Band-stop filter

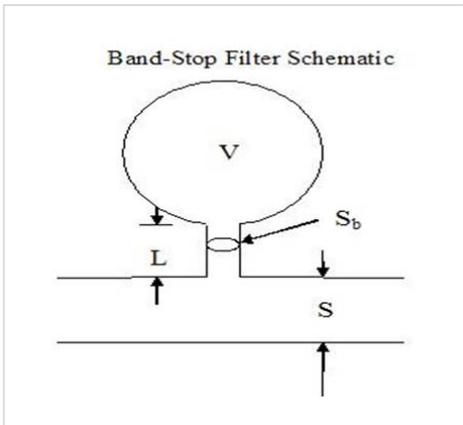
These are devices that attenuate the radiated sound power over a certain frequency range (see figure to right). Like before, the power transmission coefficient is approximately 1 in the band pass region.

Since the band-stop filter is essentially a cross between a low and high pass filter, one might expect to create one by using a combination of both techniques. This is true in that the combination of a lumped acoustic mass and compliance gives a band-stop filter. This can be realized as a helmholtz resonator (see figure to right). Again, since the impedance of the helmholtz resonator can be easily determined, continuity of acoustic impedance at the junction can give the power transmission coefficient as [1]:



$$T_\pi = \left(\frac{1}{1 + \left(\frac{c/2S}{\omega L/S_b - c^2/\omega V} \right)^2} \right)$$

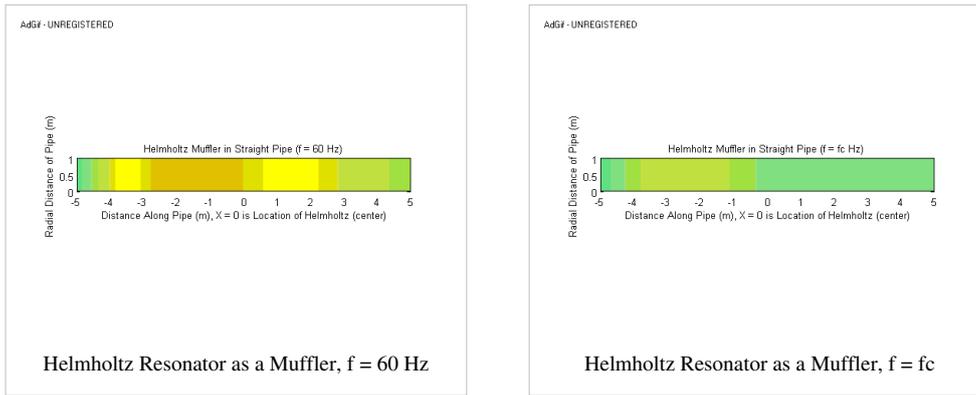
where S_b is the area of the neck, L is the effective length of the neck, V is the volume of the helmholtz resonator, and S is the area of the pipe. It is interesting to note that the power transmission coefficient is zero when the frequency is that of the resonance frequency of the helmholtz. This can be explained by the fact that at resonance the volume velocity in the neck is large with a phase such that all the incident wave is reflected back to the source [1].



The zero power transmission coefficient location is given by:

$$f_c = \left(\frac{c}{2\pi} \right) \sqrt{\left(\frac{S_b}{LV} \right)}$$

This frequency value has powerful implications. If a system has the majority of noise at one frequency component, the system can be "tuned" using the above equation, with a helmholtz resonator, to perfectly attenuate any transmitted power (see examples below).



Design

If the long wavelength assumption is valid, typically a combination of methods described above are used to design a filter. A specific design procedure is outlined for a helmholtz resonator, and other basic filters follow a similar procedure (see [1^[3]]).

Two main metrics need to be identified when designing a helmholtz resonator [3]:

1. Resonance frequency desired: $f_c = \frac{c}{2\pi} \frac{\sqrt{C_o}}{V}$ where $C_o = \frac{S}{L}$.
2. - Transmission loss: $\frac{\sqrt{C_o V}}{2S} = const$ based on TL level. This constant is found from a TL graph (see HR^[4] pp. 6).

This will result in two equations with two unknowns which can be solved for the unknown dimensions of the helmholtz resonator. It is important to note that flow velocities degrade the amount of transmission loss at resonance and tend to move the resonance location upwards [3].

In many situations, the long wavelength approximation is not valid and alternative methods must be examined. These are much more mathematically rigorous and require a complete understanding acoustics involved. Although the mathematics involved are not shown, common filters used are given in the section that follows.

Actual filter design

As explained previously, there are two main types of filters used in practice: absorptive and reactive. The benefits and drawback of each will be briefly explained, along with their relative applications (see Absorptive Mufflers).

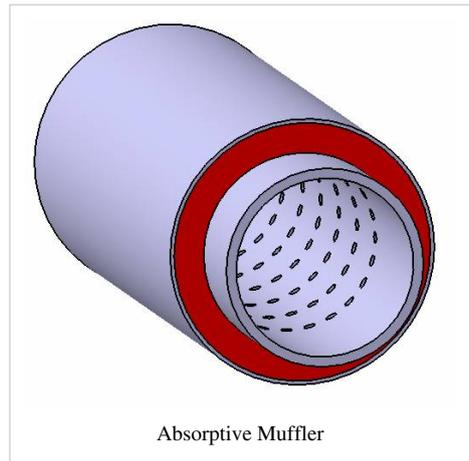
Absorptive

These are mufflers which incorporate sound absorbing materials to transform acoustic energy into heat. Unlike reactive mufflers which use destructive interference to minimize radiated sound power, absorptive mufflers are typically straight through pipes lined with multiple layers of absorptive materials to reduce radiated sound power. The most important property of absorptive mufflers is the attenuation constant. Higher attenuation constants lead to more energy dissipation and lower radiated sound power.

Advantages of Absorptive Mufflers [3]:
(1) - High amount of absorption at larger frequencies. (2) - Good for applications involving broadband (constant across the spectrum) and narrowband noise. (3) - Reduced amount of back pressure compared to reactive mufflers.
Disadvantages of Absorptive Mufflers [3]:
(1) - Poor performance at low frequencies. (2) - Material can degrade under certain circumstances (high heat, etc.).

Examples

There are a number of applications for absorptive mufflers. The most well known application is in race cars, where engine performance is desired. Absorptive mufflers don't create a large amount of back pressure (as in reactive mufflers) to attenuate the sound, which leads to higher muffler performance. It should be noted however, that the radiate sound is much higher. Other applications include plenum chambers (large chambers lined with absorptive materials, see picture below), lined ducts, and ventilation systems.



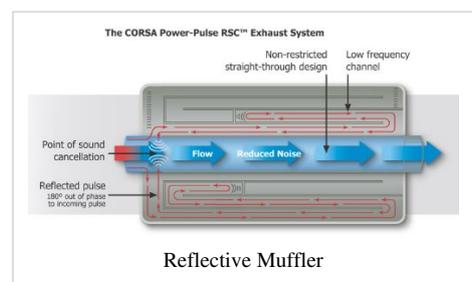
Reactive

Reactive mufflers use a number of complex passages (or lumped elements) to reduce the amount of acoustic energy transmitted. This is accomplished by a change in impedance at the intersections, which gives rise to reflected waves (and effectively reduces the amount of transmitted acoustic energy). Since the amount of energy transmitted is minimized, the reflected energy back to the source is quite high. This can actually degrade the performance of engines and other sources. Opposite to absorptive mufflers, which dissipate the acoustic energy, reactive mufflers keep the energy contained within the system. See #The_reflector_muffler Reactive Mufflers for more information.

Advantages of Reactive Mufflers [3]:
(1) - High performance at low frequencies. (2) - Typically give high insertion loss, IL, for stationary tones. (3) - Useful in harsh conditions.
Disadvantages of Reactive Mufflers [3]:
(1) - Poor performance at high frequencies. (2) - Not desirable characteristics for broadband noise.

Examples

Reactive mufflers are the most widely used mufflers in combustion engines[1^[5]]. Reactive mufflers are very efficient in low frequency applications (especially since simple lumped element analysis can be applied). Other application areas include: harsh environments (high temperature/velocity engines, turbines, etc.), specific frequency attenuation (using a helmholtz like device, a specific frequency can be toned to give total attenuation of radiated sound power), and a need for low radiated sound power (car mufflers, air conditioners, etc.).



Performance

There are 3 main metrics used to describe the performance of mufflers; Noise Reduction, Insertion Loss, and Transmission Loss. Typically when designing a muffler, 1 or 2 of these metrics is given as a desired value.

Noise Reduction (NR)

Defined as the difference between sound pressure levels on the source and receiver side. It is essentially the amount of sound power reduced between the location of the source and termination of the muffler system (it doesn't have to be the termination, but it is the most common location) [3].

$$NR = (L_{p1} - L_{p2})$$

where L_{p1} and L_{p2} is sound pressure levels at source and receiver respectively. Although NR is easy to measure, pressure typically varies at source side due to standing waves [3].

Insertion Loss (IL)

Defined as difference of sound pressure level at the receiver with and without sound attenuating barriers. This can be realized, in a car muffler, as the difference in radiated sound power with just a straight pipe to that with an expansion chamber located in the pipe. Since the expansion chamber will attenuate some of the radiate sound power, the pressure at the receiver with sound attenuating barriers will be less. Therefore, a higher insertion loss is desired [3].

$$IL = (L_{p,without} - L_{p,with})$$

where $L_{p,without}$ and $L_{p,with}$ are pressure levels at receiver without and with a muffler system respectively. Main problem with measuring IL is that the barrier or sound attenuating system needs to be removed without changing the source [3].

Transmission Loss (TL)

Defined as the difference between the sound power level of the incident wave to the muffler system and the transmitted sound power. For further information see [Transmission Loss^[6]] [3].

$$TL = 10 \log \left(\frac{1}{\tau} \right) \text{ with } \tau = \left(\frac{I_t}{I_i} \right)$$

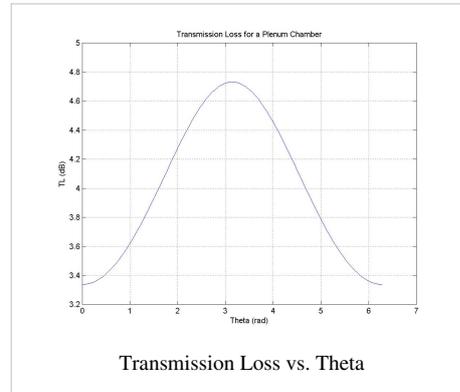
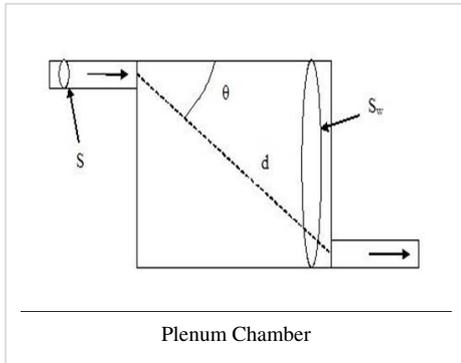
where I_t and I_i are the transmitted and incident wave power respectively. From this expression, it is obvious the problem with measure TL is decomposing the sound field into incident and transmitted waves which can be difficult to do for complex systems (analytically).

Examples

(1) - For a plenum chamber (see figure below):

$$TL = -10\log \left(S \left(\frac{\cos\theta}{2\pi d^2} + \frac{1 - \alpha}{\alpha S_w} \right) \right) \text{ in dB}$$

where α is average absorption coefficient.



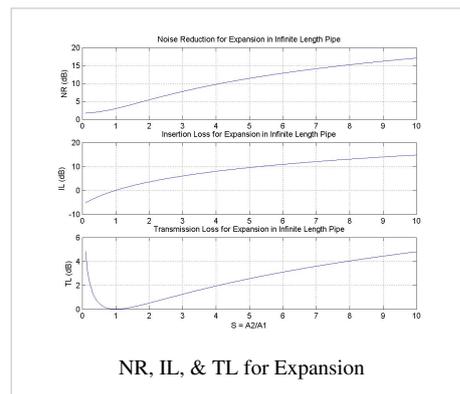
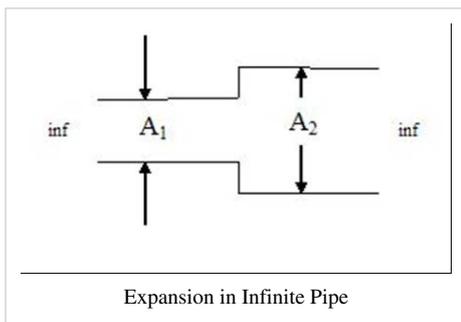
(2) - For an expansion (see figure below):

$$NR = 10\log \left[\frac{1}{2} \left| e^{-ikx_s} + \left(\frac{1 - S}{1 + S} \right) e^{ikx_s} \right|^2 (1 + S)^2 \right]$$

$$IL = 10\log \left[\frac{(1 + S)^2}{4} \right]$$

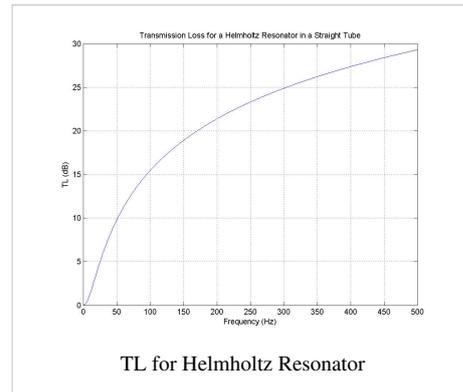
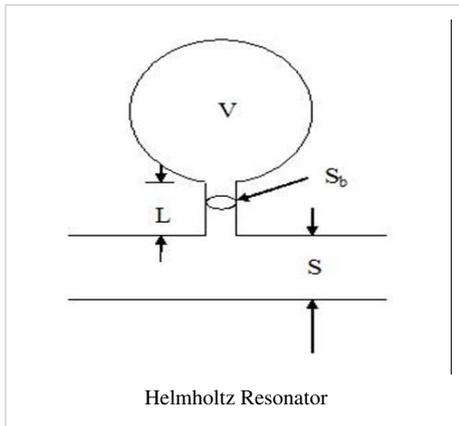
$$TL = 10\log \left[\frac{(1 + S)^2}{4S} \right]$$

where $S = \left(\frac{A_2}{A_1} \right)$



(3) - For a helmholtz resonator (see figure below):

$$TL = 10\log \left[1 + \left(\frac{\left(\frac{c}{2S_b} \right)}{\omega LS - \left(\frac{c^2}{\omega V} \right)} \right)^2 \right] \text{ in dB}$$



gdnr

Links

1. Muffler/silencer applications and descriptions of performance criteria [Exhaust Silencers ^[3]]
2. Engineering Acoustics, Purdue University - ME 513 ^[7].
3. Sound Propagation Animations ^[8]
4. Exhaust Muffler Design ^[9]
5. Project Proposal & Outline

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2. Acoustics; Pierce, Acoustical Society of America, 1989
3. - ME 413 Noise Control, Dr. Mongeau, Purdue University

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- [2] <http://www.ndt-ed.org/EducationResources/CommunityCollege/Ultrasonics/Physics/acousticimpedance.htm>
- [3] <http://www.silex.com/pdfs/Exhaust%20Silencers.pdf>
- [4] [http://mecheng.osu.edu/~selamet/docs/2003_JASA_113\(4\)_1975-1985_helmholtz_ext_neck.pdf](http://mecheng.osu.edu/~selamet/docs/2003_JASA_113(4)_1975-1985_helmholtz_ext_neck.pdf)
- [5] <http://www.eiwilliams.com/steel/index.php?p=EngineSilencers#All>
- [6] <http://freespace.virgin.net/mark.davidson3/TL/TL.html>
- [7] <http://widget.ecn.purdue.edu/~me513/>
- [8] <http://widget.ecn.purdue.edu/~me513/animate.html>
- [9] <http://myfwc.com/boating/airboat/Section3.pdf>

Flow-induced Oscillations of a Helmholtz Resonator



Introduction

The importance of flow excited acoustic resonance lies in the large number of applications in which it occurs. Sound production in organ pipes, compressors, transonic wind tunnels, and open sunroofs are only a few examples of the many applications in which flow excited resonance of Helmholtz resonators can be found.[4] An instability of the fluid motion coupled with an acoustic resonance of the cavity produce large pressure fluctuations that are felt as increased sound pressure levels.

Passengers of road vehicles with open sunroofs often experience discomfort, fatigue, and dizziness from self-sustained oscillations inside the car cabin. This phenomenon is caused by the coupling of acoustic and hydrodynamic flow inside a cavity which creates strong pressure oscillations in the passenger compartment in the 10 to 50 Hz frequency range. Some effects experienced by vehicles with open sunroofs when buffeting include: dizziness, temporary hearing reduction, discomfort, driver fatigue, and in extreme cases nausea. The importance of reducing interior noise levels inside the car cabin relies primarily in reducing driver fatigue and improving sound transmission from entertainment and communication devices.

This Wikibook page aims to theoretically and graphically explain the mechanisms involved in the flow-excited acoustic resonance of Helmholtz resonators. The interaction between fluid motion and acoustic resonance will be explained to provide a thorough explanation of the behavior of self-oscillatory Helmholtz resonator systems. As an application example, a description of the mechanisms involved in sunroof buffeting phenomena will be developed at the end of the page.

Feedback loop analysis

As mentioned before, the self-sustained oscillations of a Helmholtz resonator in many cases is a continuous interaction of hydrodynamic and acoustic mechanisms. In the frequency domain, the flow excitation and the acoustic behavior can be represented as transfer functions. The flow can be decomposed into two volume velocities.

qr: flow associated with acoustic response of cavity

qo: flow associated with excitation

Acoustical characteristics of the resonator

Lumped parameter model

The lumped parameter model of a Helmholtz resonator consists of a rigid-walled volume open to the environment through a small opening at one end. The dimensions of the resonator in this model are much less than the acoustic wavelength, in this way allowing us to model the system as a lumped system.

Figure 2 shows a sketch of a Helmholtz resonator on the left, the mechanical analog on the middle section, and the electric-circuit analog on the right hand side. As shown in the Helmholtz resonator drawing, the air mass flowing through an inflow of volume velocity includes the mass inside the neck (M_o) and an end-correction mass (M_{end}). Viscous losses at the edges of the neck length are included as well as the radiation resistance of the tube. The electric-circuit analog shows the resonator modeled as a forced harmonic oscillator. [1] [2][3]

Figure 2

V: cavity volume

ρ : ambient density

c: speed of sound

S: cross-section area of orifice

K: stiffness

M_a : acoustic mass

C_a : acoustic compliance

The equivalent stiffness K is related to the potential energy of the flow compressed inside the cavity. For a rigid wall cavity it is approximately:

$$K = \left(\frac{\rho c^2}{V}\right) S^2$$

The equation that describes the Helmholtz resonator is the following:

$$S \hat{P}_e = \frac{\hat{q}_e}{j\omega S} (-\omega^2 M + j\omega R + K)$$

\hat{P}_e : excitation pressure

M: total mass (mass inside neck M_o plus end correction, M_{end})

R: total resistance (radiation loss plus viscous loss)

From the electrical-circuit we know the following:

$$M_a = \frac{L\rho}{S}$$

$$C_a = \frac{\pi V}{\rho c^2}$$

$$L' = L + 1.7 re$$

The main cavity resonance parameters are resonance frequency and quality factor which can be estimated using the parameters explained above (assuming free field radiation, no viscous losses and leaks, and negligible wall compliance effects)

$$\omega_r^2 = \frac{1}{M_a C_a}$$

$$f_r = c2\pi\sqrt{\frac{S}{L'V}}$$

The sharpness of the resonance peak is measured by the quality factor Q of the Helmholtz resonator as follows:

$$Q = 2\pi\sqrt{V\left(\frac{L'}{S}\right)^3}$$

f_r : resonance frequency in Hz

ω_r : resonance frequency in radians

L : length of neck

L' : corrected length of neck

From the equations above, the following can be deduced:

- The greater the volume of the resonator, the lower the resonance frequencies.
- If the length of the neck is increased, the resonance frequency decreases.

Production of self-sustained oscillations

The acoustic field interacts with the unstable hydrodynamic flow above the open section of the cavity, where the grazing flow is continuous. The flow in this section separates from the wall at a point where the acoustic and hydrodynamic flows are strongly coupled. [5]

The separation of the boundary layer at the leading edge of the cavity (front part of opening from incoming flow) produces strong vortices in the main stream. As observed in Figure 3, a shear layer crosses the cavity orifice and vortices start to form due to instabilities in the layer at the leading edge.

Figure 3

From Figure 3, L is the length of the inner cavity region, d denotes the diameter or length of the cavity length, D represents the height of the cavity, and δ describes the gradient length in the grazing velocity profile (boundary layer thickness).

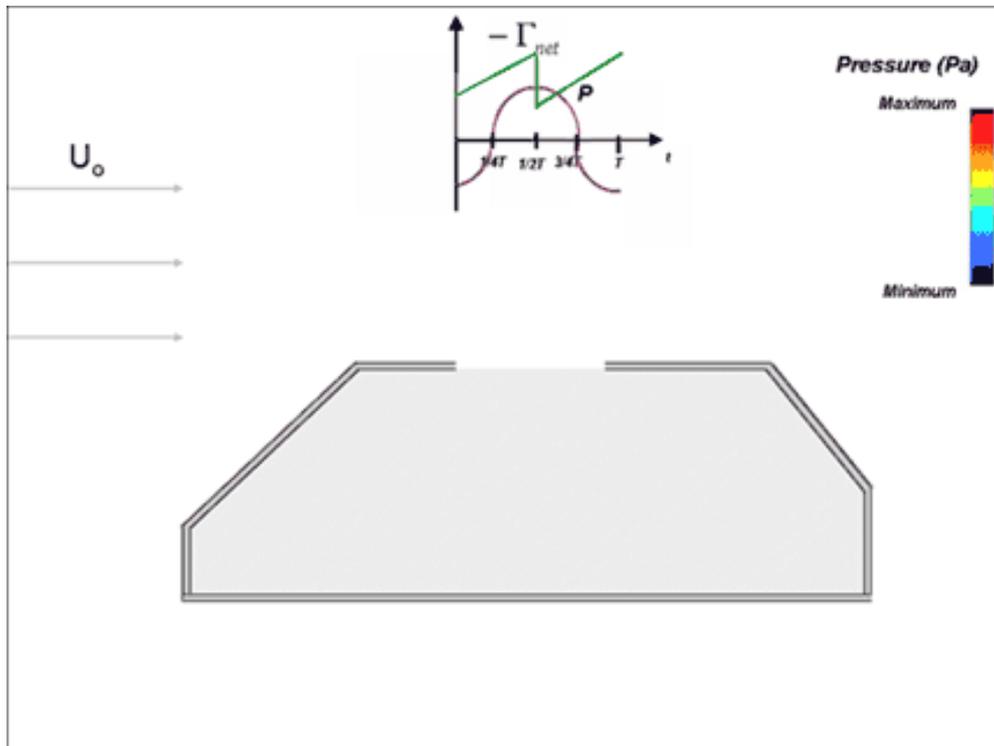
The velocity in this region is characterized to be unsteady and the perturbations in this region will lead to self-sustained oscillations inside the cavity. Vortices will continually form in the opening region due to the instability of the shear layer at the leading edge of the opening.

Applications to Sunroof Buffeting

How are vortices formed during buffeting?

In order to understand the generation and convection of vortices from the shear layer along the sunroof opening, the animation below has been developed. At a certain range of flow velocities, self-sustained oscillations inside the open cavity (sunroof) will be predominant. During this period of time, vortices are shed at the trailing edge of the opening and continue to be convected along the length of the cavity opening as pressure inside the cabin decreases and increases. Flow visualization experimentation is one method that helps obtain a qualitative understanding of vortex formation and conduction.

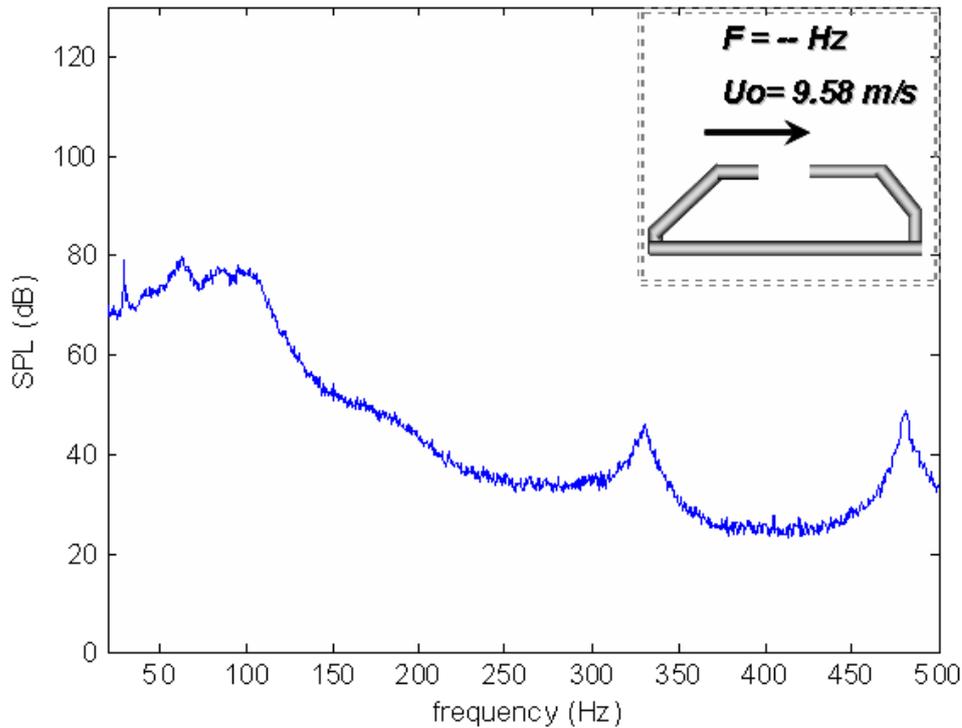
The animation below, shows in the middle, a side view of a car cabin with the sunroof open. As the air starts to flow at a certain mean velocity U_0 , air mass will enter and leave the cabin as the pressure decreases and increases again. At the right hand side of the animation, a legend shows a range of colors to determine the pressure magnitude inside the car cabin. At the top of the animation, a plot of circulation and acoustic cavity pressure versus time for one period of oscillation is shown. The symbol x moving along the acoustic cavity pressure plot is synchronized with pressure fluctuations inside the car cabin and with the legend on the right. For example, whenever the x symbol is located at the point where $t=0$ (when the acoustic cavity pressure is minimum) the color of the car cabin will match that of the minimum pressure in the legend (blue).



The perturbations in the shear layer propagate with a velocity of the order of $1/2U_0$ which is half the mean inflow velocity. [5] After the pressure inside the cavity reaches a minimum (blue color) the air mass position in the neck of the cavity reaches its maximum outward position. At this point, a vortex is shed at the leading edge of the sunroof opening (front part of sunroof in the direction of inflow velocity). As the pressure inside the cavity increases (progressively to red color) and the air mass at the cavity entrance is moved inwards, the vortex is displaced into the neck of the cavity. The maximum downward displacement of the vortex is achieved when the pressure inside the cabin is also maximum and the air mass in the neck of the Helmholtz resonator (sunroof opening) reaches its maximum downward displacement. For the rest of the remaining half cycle, the pressure cavity falls and the air below the neck of the resonator is moved upwards. The vortex continues displacing towards the downstream edge of the sunroof where it is convected upwards and outside the neck of the resonator. At this point the air below the neck reaches its maximum upwards displacement.[4] And the process starts once again.

How to identify buffeting

Flow induced tests performed over a range of flow velocities are helpful to determine the change in sound pressure levels (SPL) inside the car cabin as inflow velocity is increased. The following animation shows typical auto spectra results from a car cabin with the sunroof open at various inflow velocities. At the top right hand corner of the animation, it is possible to see the inflow velocity and resonance frequency corresponding to the plot shown at that instant of time.



It is observed in the animation that the SPL increases gradually with increasing inflow velocity. Initially, the levels are below 80 dB and no major peaks are observed. As velocity is increased, the SPL increases throughout the frequency range until a definite peak is observed around a 100 Hz and 120 dB of amplitude. This is the resonance frequency of the cavity at which buffeting occurs. As it is observed in the animation, as velocity is further increased, the peak decreases and disappears.

In this way, sound pressure level plots versus frequency are helpful in determining increased sound pressure levels inside the car cabin to find ways to minimize them. Some of the methods used to minimize the increased SPL levels achieved by buffeting include: notched deflectors, mass injection, and spoilers.

Useful websites

This link: [1] takes you to the website of EXA Corporation, a developer of PowerFlow for Computational Fluid Dynamics (CFD) analysis.

This link: [2] is a small news article about the current use of (CFD) software to model sunroof buffeting.

This link: [3] is a small industry brochure that shows the current use of CFD for sunroof buffeting.

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Active Control



Introduction

The principle of active control of noise, is to create destructive interferences using a secondary source of noise. Thus, any noise can theoretically disappear. But as we will see in the following sections, only low frequencies noises can be reduced for usual applications, since the amount of secondary sources required increases very quickly with frequency. Moreover, predictable noises are much easier to control than unpredictable ones. The reduction can reach up to 20dB for the best cases. But since good reduction can only be reached for low frequencies, the perception we have of the resulting sound is not necessarily as good as the theoretical reduction. This is due to psychoacoustics considerations, which will be discussed later on.

Fundamentals of active control of noise

Control of a monopole by another monopole

Even for the free space propagation of an acoustic wave created by a punctual source it is difficult to reduce noise in a large area, using active noise control, as we will see in the section.

In the case of an acoustic wave created by a monopolar source, the Helmholtz equation becomes:

$$\Delta p + k^2 p = -j\omega\rho_0 q$$

where q is the flow of the noise sources.

The solution for this equation at any M point is:

$$p_p(M) = \frac{j\omega\rho_0 q_p}{4\pi} \frac{e^{-jk r_p}}{r_p}$$

where the p mark refers to the primary source.

Let us introduce a secondary source in order to perform active control of noise. The acoustic pressure at that same M point is now:

$$p(M) = \frac{j\omega\rho_0 q_p}{4\pi} \frac{e^{-jk r_p}}{r_p} + \frac{j\omega\rho_0 q_s}{4\pi} \frac{e^{-jk r_s}}{r_s}$$

It is now obvious that if we chose $q_s = -q_p \frac{r_s}{r_p} e^{-jk(r_p - r_s)}$ there is no more noise at the M point. This is the most

simple example of active control of noise. But it is also obvious that if the pressure is zero in M , there is no reason why it should also be zero at any other N point. This solution only allows to reduce noise in one very small area.

However, it is possible to reduce noise in a larger area far from the source, as we will see in this section. In fact the expression for acoustic pressure far from the primary source can be approximated by:

$$p(M) = \frac{j\omega\rho_0}{4\pi} \frac{e^{-jk r_p}}{r_p} (q_p + q_s e^{-jk D \cos \theta})$$

As shown in the previous section we can adjust the secondary source in order to get no noise in M. In that case, the acoustic pressure in any other N point of the space remains low if the primary and secondary sources are close enough. More precisely, it is possible to have a pressure close to zero in the whole space if the M point is equally distant from the two sources and if: $D < \lambda/6$ where D is the distance between the primary and secondary sources. As we will see later on, it is easier to perform active control of noise with more than one source controlling the primary source, but it is of course much more expensive.

A commonly admitted estimation of the number of secondary sources which are necessary to reduce noise in an R radius sphere, at a frequency f is:

$$N = \frac{36\pi R^2 f^2}{c^2}$$

This means that if you want no noise in a one meter diameter sphere at a frequency below 340Hz, you will need 30 secondary sources. This is the reason why active control of noise works better at low frequencies.

Active control for waves propagation in ducts and enclosures

This section requires from the reader to know the basis of modal propagation theory, which will not be explained in this article.

Ducts

For an infinite and straight duct with a constant section, the pressure in areas without sources can be written as an infinite sum of propagation modes:

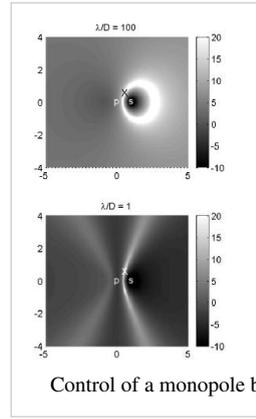
$$p(x, y, z, \omega) = \sum_{n=1}^N a_n(\omega) \phi_n(x, y) e^{-jk_n z}$$

where ϕ are the eigen functions of the Helmholtz equation and a_n represent the amplitudes of the modes.

The eigen functions can either be obtained analytically, for some specific shapes of the duct, or numerically. By putting pressure sensors in the duct and using the previous equation, we get a relation between the pressure matrix P (pressure for the various frequencies) and the A matrix of the amplitudes of the modes. Furthermore, for linear sources, there is a relation between the A matrix and the U matrix of the signal sent to the secondary sources: $A_s = KU$ and hence: $A = A_p + A_s = A_p + KU$.

Our purpose is to get: $A=0$, which means: $A_p + KU = 0$. This is possible every time the rank of the K matrix is bigger than the number of the propagation modes in the duct.

Thus, it is theoretically possible to have no noise in the duct in a very large area not too close from the primary sources if there are more secondary sources than propagation modes in the duct. Therefore, it is obvious that active noise control is more appropriate for low frequencies. In fact the more the frequency is low, the less propagation modes there will be in the duct. Experiences show that it is in fact possible to reduce the noise from over 60dB.



Enclosures

The principle is rather similar to the one described above, except the resonance phenomenon has a major influence on acoustic pressure in the cavity. In fact, every mode that is not resonant in the considered frequency range can be neglected. In a cavity or enclosure, the number of these modes rise very quickly as frequency rises, so once again, low frequencies are more appropriate. Above a critical frequency, the acoustic field can be considered as diffuse. In that case, active control of noise is still possible, but it is theoretically much more complicated to set up.

Active control and psychoacoustics

As we have seen, it is possible to reduce noise with a finite number of secondary sources. Unfortunately, the perception of sound of our ears does not only depend on the acoustic pressure (or the decibels). In fact, it sometimes happen that even though the number of decibels has been reduced, the perception that we have is not really better than without active control.

Active control systems

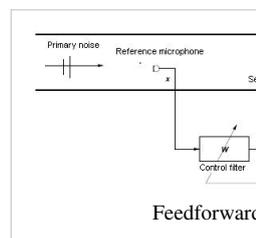
Since the noise that has to be reduced can never be predicted exactly, a system for active control of noise requires an auto adaptable algorithm. We have to consider two different ways of setting up the system for active control of noise depending on whether it is possible or not to detect the noise from the primary source before it reaches the secondary sources. If this is possible, a feed forward technique will be used (aircraft engine for example). If not a feed back technique will be preferred.

Feedforward

In the case of a feed forward, two sensors and one secondary source are required. The sensors measure the sound pressure at the primary source (detector) and at the place we want noise to be reduced (control sensor). Furthermore, we should have an idea of what the noise from the primary source will become as he reaches the control sensor. Thus we approximately know what correction should be made, before the sound wave reaches the control sensor (forward). The control sensor will only correct an eventual or residual error. The feedforward technique allows to reduce one specific noise (aircraft engine for example) without reducing every other sound (conversations, ...). The main issue for this technique is that the location of the primary source has to be known, and we have to be sure that this sound will be detected beforehand. Therefore portative systems based on feed forward are impossible since it would require having sensors all around the head.

Feedback

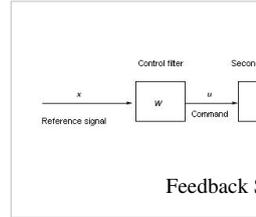
In that case, we do not exactly know where the sound comes from; hence there is only one sensor. The sensor and the secondary source are very close from each other and the correction is done in real time: as soon as the sensor gets the information the signal is treated by a filter which sends the corrected signal to the secondary source. The main issue with feedback is that every noise is reduced and it is even theoretically impossible to have a standard conversation.



Applications

Noise cancelling headphone

Usual headphones become useless when the frequency gets too low. As we have just seen active noise cancelling headphones require the feedback technique since the primary sources can be located all around the head. This active control of noise is not really efficient at high frequencies since it is limited by the Larsen effect. Noise can be reduced up to 30dB at a frequency range between 30Hz and 500Hz.



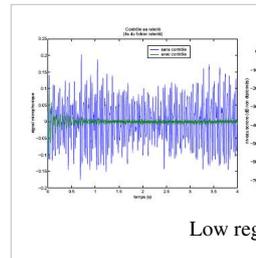
Active control for cars

Noise reduction inside cars can have a significant impact on the comfort of the driver. There are three major sources of noise in a car: the motor, the contact of tires on the road, and the aerodynamic noise created by the air flow around the car. In this section, active control for each of those sources will be briefly discussed.

Motor noise

This noise is rather predictable since it a consequence of the rotation of the pistons in the motor. Its frequency is not exactly the motor's rotational speed though. However, the frequency of this noise is in between 20Hz and 200Hz, which means that an active control is theoretically possible. The following pictures show the result of an active control, both for low and high regime.

Even though these results show a significant reduction of the acoustic pressure, the perception inside the car is not really better with this active control system, mainly for psychoacoustics reasons which were mentioned above. Moreover such a system is rather expensive and thus are not used in commercial cars.



Tires noise

This noise is created by the contact between the tires and the road. It is a broadband noise which is rather unpredictable since the mechanisms are very complex. For example, the different types of roads can have a significant impact on the resulting noise. Furthermore, there is a cavity around the tires, which generate a resonance phenomenon. The first frequency is usually around 200Hz. Considering the multiple causes for that noise and its unpredictability, even low frequencies become hard to reduce. But since this noise is broadband, reducing low frequencies is not enough to reduce the overall noise. In fact an active control system would mainly be useful in the case of an unfortunate amplification of a specific mode.

Aerodynamic noise

This noise is a consequence of the interaction between the air flow around the car and the different appendices such as the rear views for example. Once again, it is an unpredictable broadband noise, which makes it difficult to reduce with an active control system. However, this solution can become interesting in the case an annoying predictable resonance would appear.

Active control for aeronautics

The noise of aircraft propellers is highly predictable since the frequency is quite exactly the rotational frequency multiplied by the number of blades. Usually this frequency is around some hundreds of Hz. Hence, an active control system using the feedforward technique provides very satisfying noise reductions. The main issues are the cost and the weigh of such a system. The fan noise on aircraft engines can be reduced in the same manner.

Further reading

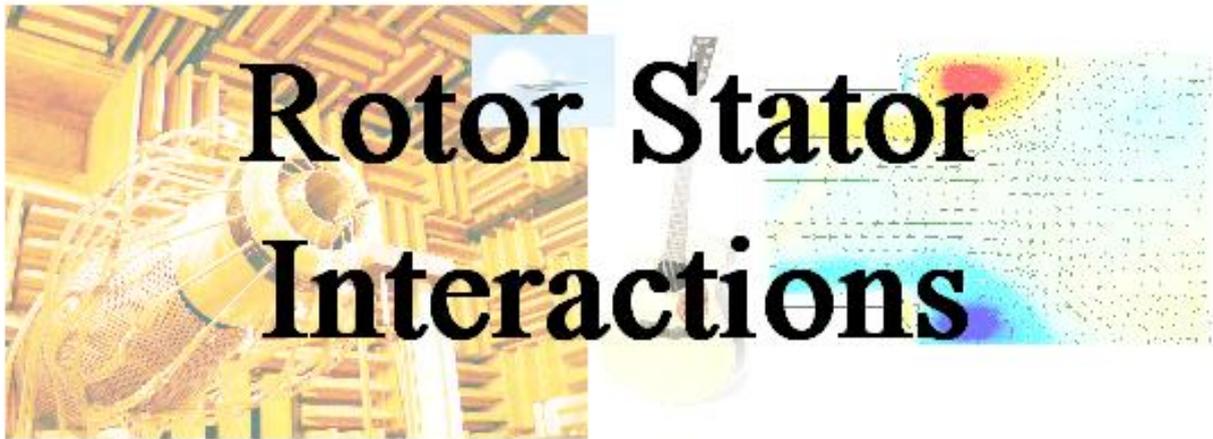
- "Active Noise Control" ^[1] at Dirac delta.

References

- [1] <http://www.diracdelta.co.uk/science/source/a/c/active%20noise%20control/source.html>

Applications in Transport Industry

Rotor Stator Interactions



An important issue for the aeronautical industry is the reduction of aircraft noise. The characteristics of the turbomachinery noise are to be studied. The rotor/stator interaction is a predominant part of the noise emission. We will present an introduction to these interaction theory, whose applications are numerous. For example, the conception of air-conditioning ventilators requires a full understanding of this interaction.

Noise emission of a Rotor-Stator mechanism

A Rotor wake induces on the downstream Stator blades a fluctuating vane loading, which is directly linked to the noise emission.

We consider a B blades Rotor (at a rotation speed of Ω) and a V blades stator, in a unique Rotor/Stator configuration. The source frequencies are multiples of $B\Omega$, that is to say $mB\Omega$. For the moment we don't have access to the source levels F_m . The noise frequencies are also $mB\Omega$, not depending on the number of blades of the stator. Nevertheless, this number V has a predominant role in the noise levels (P_m) and directivity, as it will be discussed later.

Example

For an airplane air-conditioning ventilator, reasonable data are :

$$B = 13 \text{ and } \Omega = 12000 \text{ rnd/min}$$

The blade passing frequency is 2600 Hz, so we only have to include the first two multiples (2600 Hz and 5200 Hz), because of the human ear high-sensibility limit. We have to study the frequencies $m=1$ and $m=2$.

Optimization of the number of blades

As the source levels can't be easily modified, we have to focus on the interaction between those levels and the noise levels.

The transfer function $\frac{F_m}{P_m}$ contains the following part :

$$\sum_{s=-\infty}^{s=+\infty} e^{-\frac{i(mB-sV)\pi}{2}} J_{mB-sV}(mBM)$$

Where m is the Mach number and J_{mB-sV} the Bessel function of $mB-sV$ order. In order to minimize the influence of the transfer function, the goal is to reduce the value of this Bessel function. To do so, the argument must be smaller than the order of the Bessel function.

Back to the example :

For $m=1$, with a Mach number $M=0.3$, the argument of the Bessel function is about 4. We have to avoid having $mB-sV$ inferior than 4. If $V=10$, we have $13-1 \times 10=3$, so there will be a noisy mode. If $V=19$, the minimum of $mB-sV$ is 6, and the noise emission will be limited.

Remark :

The case that is to be strictly avoided is when $mB-sV$ can be nul, which causes the order of the Bessel function to be 0. As a consequence, we have to take care having B and V prime numbers.

Determination of source levels

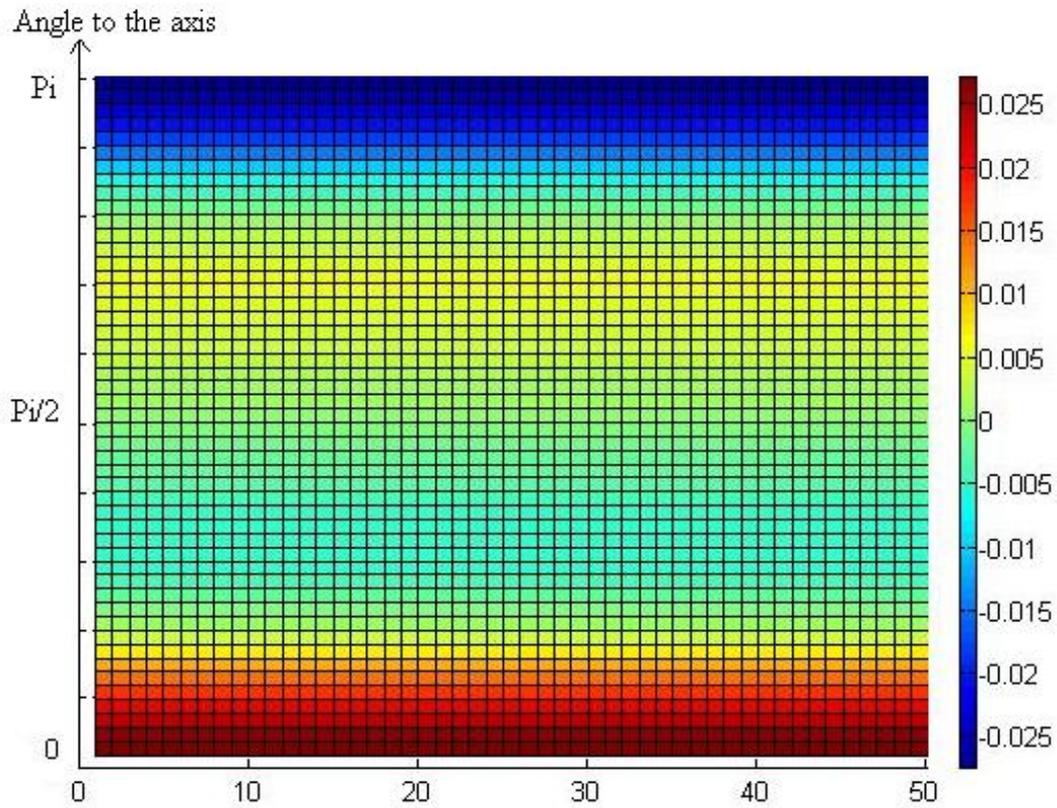
The minimization of the transfer function $\frac{F_m}{P_m}$ is a great step in the process of reducing the noise emission.

Nevertheless, to be highly efficient, we also have to predict the source levels F_m . This will lead us to choose to minimize the Bessel functions for the most significant values of m . For example, if the source level for $m=1$ is very higher than for $m=2$, we will not consider the Bessel functions of order $2B-sV$. The determination of the source levels is given by the Sears theory, which will not be discussed here.

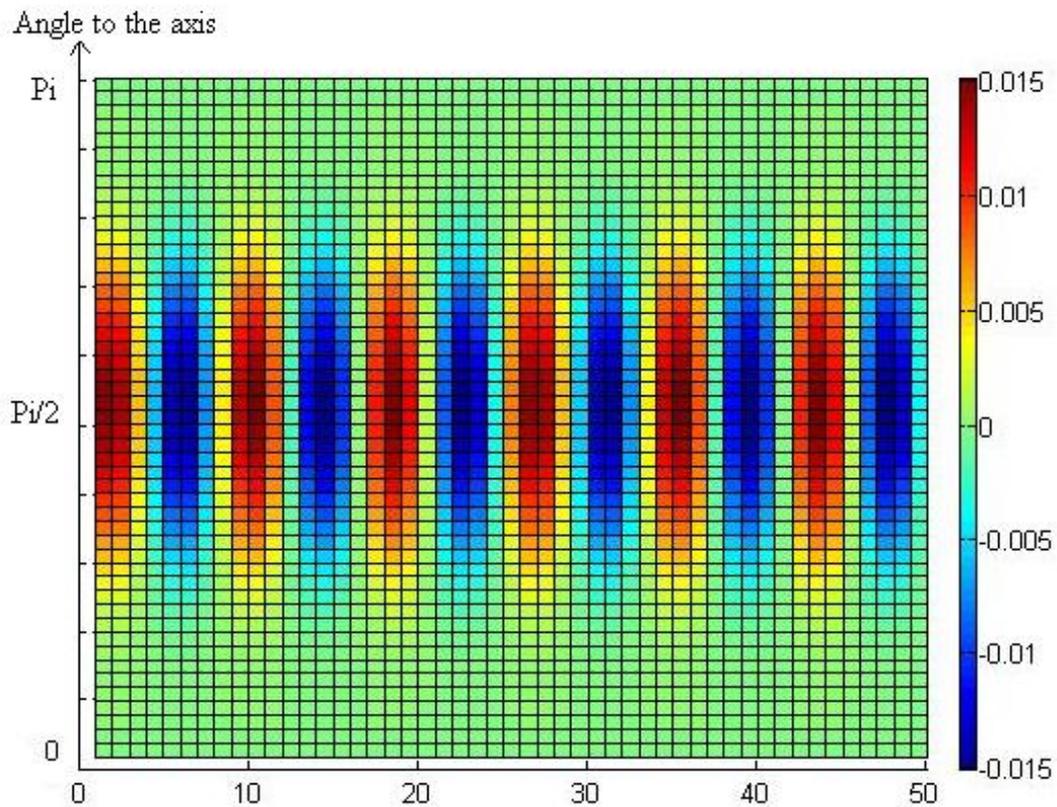
Directivity

All this study was made for a specific direction : the axis of the Rotor/Stator. All the results are acceptable when the noise reduction is ought to be in this direction. In the case where the noise to reduce is perpendicular to the axis, the results are very different, as those figures shown :

For B=13 and V=13, which is the worst case, we see that the sound level is very high on the axis (for $\theta = 0$)



For B=13 and V=19, the sound level is very low on the axis but high perpendicularly to the axis (for $\theta = \pi/2$)



Further reading

This module discusses rotor/stator interaction, the predominant part of the noise emission of turbomachinery. See Acoustics/Noise from Cooling Fans for a discussion of other noise sources.

External references

- Prediction of rotor wake-stator interaction noise by P. Sijtsma and J.B.H.M. Schulten ^[1]

References

- [1] <http://www.nlr.nl/documents/publications/2003/2003-124-tp.pdf>

Car Mufflers



Introduction

A car muffler is a component of the exhaust system of a car. The exhaust system has mainly 3 functions:

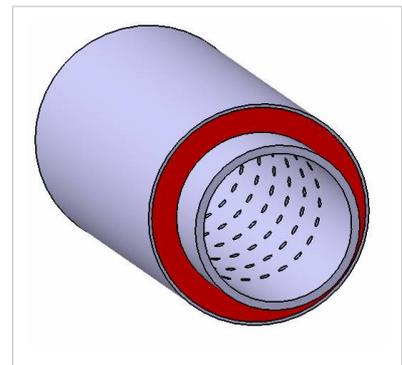
1. Getting the hot and noxious gas from the engine away from the vehicle
2. Reduce exhaust emission
3. Attenuating the noise output from the engine

The last specified function is the function of the car muffler. It is necessary because the gas coming from the combustion in the pistons of the engine would generate an extremely loud noise if it were sent directly in the ambient air surrounding engine through the exhaust valves. There are 2 techniques used to dampen the noise: absorption and reflection. Each technique has its advantages and disadvantages.

The absorber muffler

The muffler is composed of a tube covered by sound absorbing material. The tube is perforated so that some part of the sound wave goes through the perforation to the absorbing material. The absorbing material is usually made of fiberglass or steel wool. The dampening material is protected from the surrounding by a supplementary coat made of a bend metal sheet.

The advantage of this method is low back pressure with a relatively simple design. The inconvenience of this method is low sound damping ability compared to the other techniques, especially at low frequency.



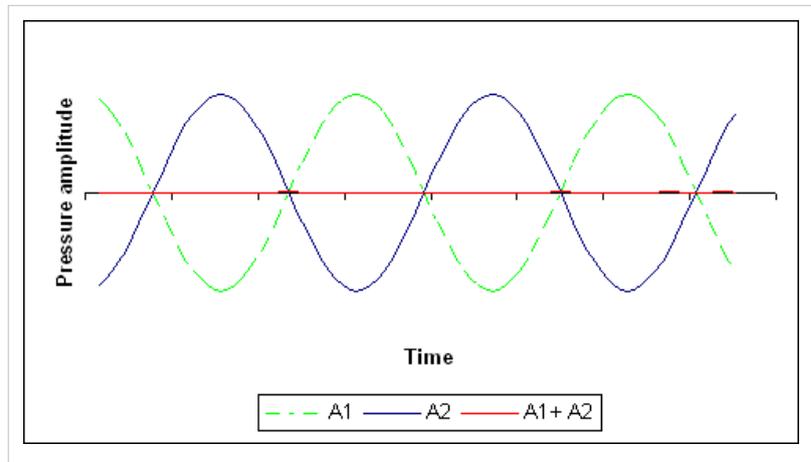
The mufflers using the absorption technique are usually sports vehicle because they increase the performances of the engine because of their low back pressure. A trick to improve their muffling ability consists of lining up several "straight" mufflers.

The reflector muffler

Principle: Sound wave reflection is used to create a maximum amount of destructive interferences

Definition of destructive interferences

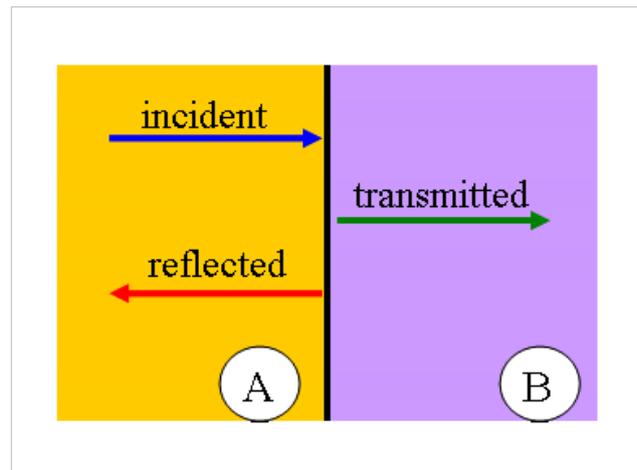
Let's consider the noise a person would hear when a car drives past. This sound would physically correspond to the pressure variation of the air which would make his ear-drum vibrate. The curve A1 of the graph 1 could represent this sound. The pressure amplitude is a function of the time at a certain fixed place. If another sound wave A2 is produced at the same time,



the pressure of the two waves will add. If the amplitude of A1 is exactly the opposite of the amplitude A2, then the sum will be zero, which corresponds physically to the atmospheric pressure. The listener would thus hear nothing although there are two radiating sound sources. A2 is called the destructive interference.

Definition of the reflection

The sound is a traveling wave i.e. its position changes in function of the time. As long as the wave travels in the same medium, there is no change of speed and amplitude. When the wave reaches a frontier between two mediums which have different impedances, the speed, and the pressure amplitude change (and so does the angle if the wave does not propagate perpendicularly to the frontier). The figure 1 shows two medium A and B and the 3 waves: incident transmitted and reflected.

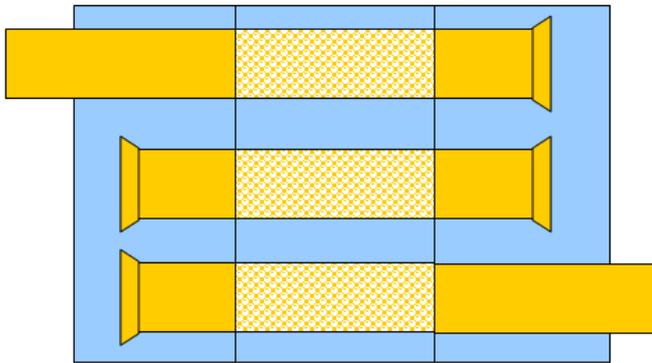


Example

If plane sound waves are propagating across a tube and the section of the tube changes, the impedance of the tube will change. Part of the incident waves will be transmitted through the discontinuity and the other part will be reflected.

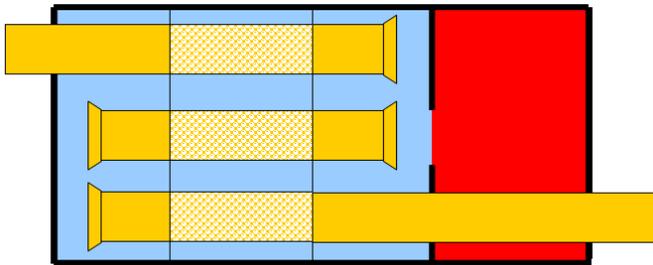
Animation

Mufflers using the reflection technique are the most commonly used because they dampen the noise much better than the absorber muffler. However they induce a higher back pressure, lowering the performance of the engine (an engine would be most efficient or powerful without the use of a muffler).



The upper right image represents a car muffler's typical architecture. It is composed of 3 tubes. There are 3 areas separated by plates, the part of the tubes located in the middle area are perforated. A small quantity of pressure "escapes" from the tubes through the perforation and cancel one another.

Some mufflers using the reflection principle may incorporate cavities which dampen noise. These cavities are called Helmholtz Resonators in acoustics. This feature is usually only available for up market class mufflers.



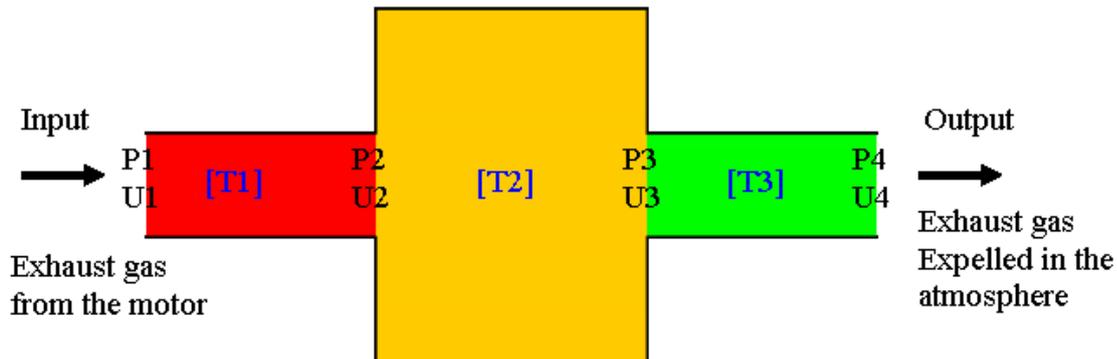
Back pressure

Car engines are 4 stroke cycle engines. Out of these 4 strokes, only one produces the power, this is when the explosion occurs and pushes the pistons back. The other 3 strokes are necessary evil that don't produce energy. They on the contrary consume energy. During the exhaust stroke, the remaining gas from the explosion is expelled from the cylinder. The higher the pressure behind the exhaust valves (i.e. back pressure), and the higher effort necessary to expel the gas out of the cylinder. So, a low back pressure is preferable in order to have a higher engine horsepower.

Muffler modeling by transfer matrix method

This method is easy to use on computer to obtain theoretical values for the transmission loss of a muffler. The transmission loss gives a value in dB that correspond to the ability of the muffler to dampen the noise.

Example



P stands for Pressure [Pa] and U stand for volume velocity [m³/s]

$$\begin{bmatrix} P1 \\ U1 \end{bmatrix} = [T1] \begin{bmatrix} P2 \\ U2 \end{bmatrix} \text{ and } \begin{bmatrix} P2 \\ U2 \end{bmatrix} = [T2] \begin{bmatrix} P3 \\ U3 \end{bmatrix} \text{ and } \begin{bmatrix} P3 \\ U3 \end{bmatrix} = [T3] \begin{bmatrix} P4 \\ U4 \end{bmatrix}$$

So, finally:
$$\begin{bmatrix} P1 \\ U1 \end{bmatrix} = [T1] [T2] [T3] \begin{bmatrix} P4 \\ U4 \end{bmatrix}$$

with

$$[T_i] = \begin{bmatrix} \cos(kL_i) & j \sin(kL_i) \frac{\rho c}{S_i} \\ j \sin(kL_i) \frac{S_i}{\rho c} & \cos(kL_i) \end{bmatrix}$$

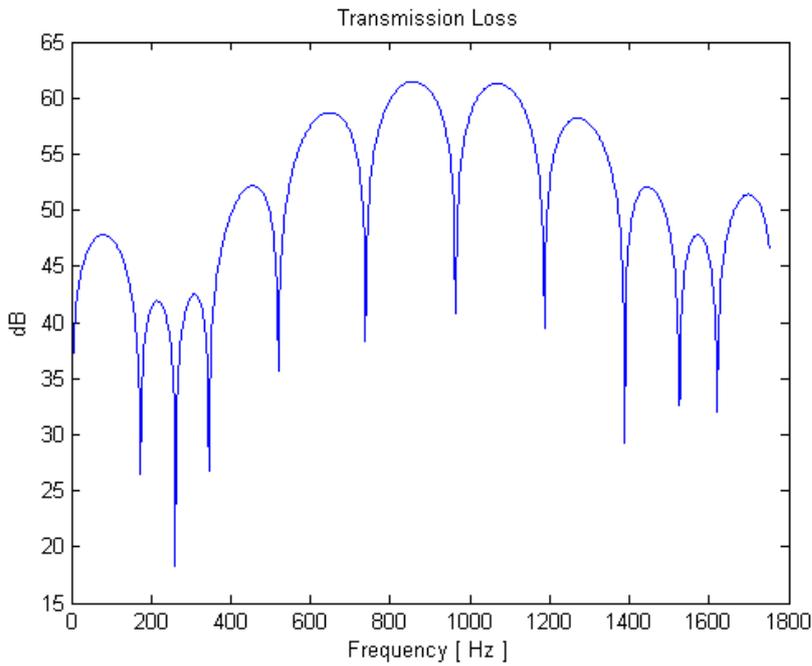
S_i stands for the cross section area

k is the angular velocity

ρ is the medium density

c is the speed of sound of the medium

Results



Matlab code of the graph above.

Comments

The higher the value of the transmission loss and the better the muffler.

The transmission loss depends on the frequency. The sound frequency of a car engine is approximately between 50 and 3000Hz. At resonance frequencies, the transmission loss is zero. These frequencies correspond to the lower peaks on the graph.

The transmission loss is independent of the applied pressure or velocity at the input.

The temperature (about 600 Fahrenheit) has an impact on the air properties : the speed of sound is higher and the mass density is lower.

The elementary transfer matrix depends on the element which is modelled. For instance the transfer matrix of a Helmutz Resonator is $\begin{bmatrix} 1 & 0 \\ \frac{1}{Z} & 1 \end{bmatrix}$ with $Z = j\rho\left(\frac{\omega L_i}{S_i} - \frac{c^2}{\omega V}\right)$

The transmission loss and the insertion loss are different terms. The transmission loss is 10 times the logarithm of the ratio output/input. The insertion loss is 10 times the logarithm of the ratio of the radiated sound power with and without muffler.

Links

- General Information about Filter Design & Implementation
- More information about the Transfer Matrix Method ^[1]
- General information about car mufflers ^[2]

References

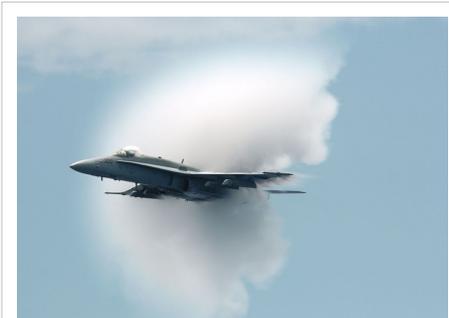
[1] <http://www.scielo.br/pdf/jbsmse/v27n2/25381.pdf>

[2] <http://auto.howstuffworks.com/muffler.htm>

Sonic Boom



A **sonic boom** is the audible component of a shock wave in air. The term is commonly used to refer to the air shocks caused by the supersonic flight of military aircraft or passenger transports such as Concorde (Mach 2.2, no longer flying) and the Space Shuttle (Mach 27). Sonic booms generate enormous amounts of sound energy, sounding much like an explosion; typically the shock front may approach 100 megawatts per square meter, and may exceed 200 decibels.



Warplane passing the sound barrier.

Cause of sonic booms

As an object moves through the air it creates a series of pressure waves in front and behind it, similar to the bow and stern waves created by a boat. These waves travel at the speed of sound, and as the speed of the aircraft increases the waves are forced together or 'compressed' because they cannot "get out of the way" of each other, eventually merging into a single shock wave at the speed of sound. This critical speed is known as Mach 1 and is approximately 1,225 km/h (761 mph) at sea level.

In smooth flight, the shock wave starts at the nose of the aircraft and ends at the tail. There is a sudden drop in pressure at the nose, decreasing steadily to a negative pressure at the tail, where it suddenly returns to normal. This "overpressure profile" is known as the N-wave due to its shape. We experience the "boom" when there is a sudden drop in pressure, so the N-wave causes two booms, one when the initial pressure rise from the nose hits, and another when the tail passes and the pressure suddenly returns to normal. This leads to a distinctive "double boom" from supersonic aircraft. When maneuvering the pressure; distribution changes into different forms, with a characteristic U-wave shape. Since the boom is being generated continually as long as the aircraft is supersonic, it traces out a path on the ground following the aircraft's flight path, known as the **boom carpet**.

A sonic boom or "tunnel boom" can also be caused by high-speed trains in tunnels (e.g. the Japanese Shinkansen). In order to reduce the sonic boom effect, a special shape of the train car and a widened opening of the tunnel entrance is necessary. When a high speed train enters a tunnel, the sonic boom effect occurs at the tunnel exit. In contrast to the (super)sonic boom of an aircraft, this "tunnel boom" is caused by a rapid change of subsonic flow (due to the sudden narrowing of the surrounding space) rather than by a shock wave. In close range to the tunnel exit this phenomenon can cause disturbances to residents.

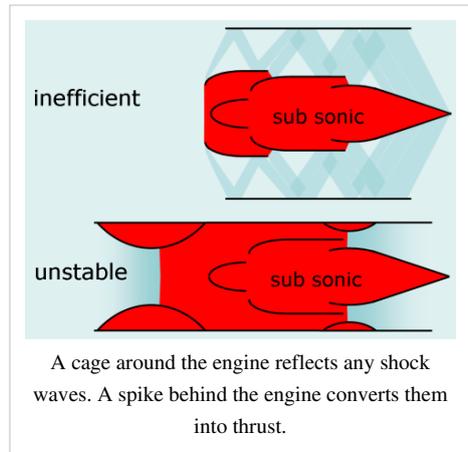
Characteristics

The power, or volume, of the shock wave is dependent on the quantity of air that is being accelerated, and thus the size and weight of the aircraft. As the aircraft increases speed the shocks grow "tighter" around the craft, and do not become much "louder". At very high speeds and altitudes the cone does not intersect the ground, and no boom will be heard. The "length" of the boom from front to back is dependent on the length of the aircraft, although to a factor of 3:2 not 1:1. Longer aircraft therefore "spread out" their booms more than smaller ones, which leads to a less powerful boom.

The nose shockwave compresses and pulls the air along with the aircraft so that the aircraft behind its shockwave sees subsonic airflow.

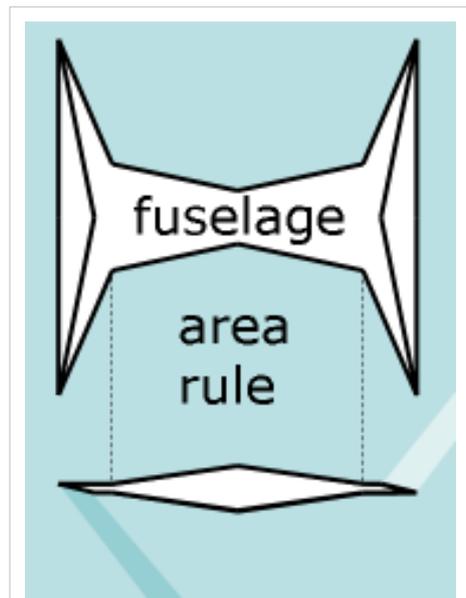


When an aircraft is near the sound barrier, an unusual cloud sometimes forms in its wake. A Prandtl-Glauert Singularity results from a drop in pressure, due to shock wave formation. This pressure change causes a sharp drop in temperature, which in humid conditions leads the water vapor in the air to condense into droplets and form the cloud.



However, this means that several smaller shock waves can, and usually do, form at other points on the aircraft, primarily any convex points or curves, the leading wing edge and especially the inlet to engines. These secondary shockwaves are caused by the subsonic air behind the main shockwave being forced to go supersonic again by the shape of the aircraft (for example due to the air's acceleration over the top of a curved wing).

The later shock waves are somehow faster than the first one, travel faster and add to the main shockwave at some distance away from the aircraft to create a much more defined N-wave shape. This maximizes both the magnitude and the "rise time" of the shock, which makes it seem louder. On most designs the characteristic distance is about 40,000ft, meaning that below this altitude the sonic boom will be "softer". However the drag at this altitude or below makes supersonic travel particularly inefficient, which poses a serious problem.



To generate lift a supersonic airplane has to produce at least two shock waves: One over-pressure downwards wave, and one under-pressure upwards wave. Whitcomb area rule states, we can reuse air displacement without generating additional shock waves. In this case the fuselage reuses some displacement of the wings.

Abatement

In the late 1950s when SST designs were being actively pursued, it was thought that although the boom would be very large, they could avoid problems by flying higher. This premise was proven false when the North American B-70 *Valkyrie* started flying and it was found that the boom was a very real problem even at 70,000ft (21,000m). It was during these tests that the N-wave was first characterized.

Richard Seebass and his colleague Albert George at Cornell University studied the problem extensively, and eventually defined a "figure of merit", **FM**, to characterize the sonic boom levels of different aircraft. FM is proportional to the aircraft weight divided by the three-halves of the aircraft length, $FM = W/(3/2 \cdot L) = 2W/3L$. The lower this value, the less boom the aircraft generates, with figures of about 1 or lower being considered acceptable. Using this calculation they found FM's of about 1.4 for Concorde, and 1.9 for the Boeing 2707. This eventually doomed most SST projects as public resentment, somewhat blown out of proportion, mixed with politics eventually resulted in laws that made any such aircraft impractical (flying only over water for instance).

Seebass-George also worked on the problem from another angle, examining ways to reduce the "peaks" of the N-wave and therefore smooth out the shock into something less annoying. Their theory suggested that body shaping might be able to use the secondary shocks to either "spread out" the N-wave, or interfere with each other to the same end. Ideally this would raise the characteristic altitude from 40,000ft to 60,000, which is where most SST designs fly. The design required some fairly sophisticated shaping in order to achieve the dual needs of reducing the shock and still leaving an aerodynamically efficient shape, and therefore had to wait for the advent of computer-aided design before being able to be built.

This remained untested for decades, until DARPA started the **Quiet Supersonic Platform** project and funded the **Shaped Sonic Boom Demonstration** aircraft to test it. SSBDD used a F-5 Freedom Fighter modified with a new body shape, and was tested over a two year period in what has become the most extensive study on the sonic boom to date.

After measuring the 1,300 recordings, some taken inside the shock wave by a chase plane, the SSBD demonstrated a reduction in boom by about one-third. Although one-third is not a huge reduction, it could reduce Concorde below the $FM = 1$ limit for instance.

There are theoretical designs that do not appear to create sonic booms at all, such as the Busemann's Biplane. Nobody has been able to suggest a practical implementation of this concept, as yet.

Perception and noise

The sound of a sonic boom depends largely on the distance between the observer and the aircraft producing the sonic boom. A sonic boom is usually heard as a deep double "boom" as the aircraft is usually some distance away. However, as those who have witnessed landings of space shuttles have heard, when the aircraft is nearby the sonic boom is a sharper "bang" or "crack". The sound is much like the "aerial bombs" used at firework displays.

In 1964, NASA and the FAA began the Oklahoma City sonic boom tests, which caused eight sonic booms per day over a period of six months. Valuable data was gathered from the experiment, but 15,000 complaints were generated and ultimately entangled the government in a class action lawsuit, which it lost on appeal in 1969.

In late October 2005, Israel began using nighttime sonic boom raids against civilian populations in the Gaza Strip [1] as a method of psychological warfare. The practice was condemned by the United Nations. A senior Israeli army intelligence source said the tactic was intended to break civilian support for armed Palestinian groups.

Media

These videos include jets achieving supersonic speeds.

First supersonic flight (info)

Chuck Yeager broke the sound barrier on October 14, 1947 in the Bell X-1.

F-14 Tomcat sonic boom flyby (with audio) (info)

F-14 Tomcat flies at Mach 1 over the water, creating a sonic boom as it passes.

F-14A Tomcat supersonic flyby (info)

Supersonic F-14A Tomcat flying by the USS Theodore Roosevelt CVN-71 in 1986 for the tiger cruise.

Shuttle passes sound barrier (info)

Space shuttle Columbia crosses the sound barrier at 45 seconds after liftoff.

External links

- NASA opens new chapter in supersonic flight ^[2]
 - "Sonic Boom ^[3]," a tutorial from the "Sonic Boom, Sound Barrier, and Condensation Clouds ^[4]" (or "Sonic Boom, Sound Barrier, and Prandtl-Glauert Condensation Clouds") collection of tutorials by Mark S. Cramer, Ph.D. at <http://FluidMech.net> (Tutorials, Sound Barrier).
 - decibel chart including sonic booms ^[5]
-

References

- [1] <http://www.guardian.co.uk/israel/Story/0,2763,1607450,00.html>
- [2] <http://spaceflightnow.com/news/n0309/06supersonic/>
- [3] <http://fluidmech.net/tutorials/sonic/sonicboom.htm>
- [4] <http://fluidmech.net/tutorials/sonic/sonicboom-intro.htm>
- [5] <http://www.makeitlouder.com/Decibel%20Level%20Chart.txt>

Sonar

SONAR (**s**ound **n**avigation and **r**anging) is a technique that uses sound propagation under water to navigate or to detect other vessels. There are two kinds of sonar: active and passive.

History

The French physicist Paul Langevin, working with a Russian émigré electrical engineer, Constantin Chilowski, invented the first active sonar-type device for detecting submarines in 1915. Although piezoelectric transducers later superseded the electrostatic transducers they used, their work influenced the future of sonar designs. In 1916, under the British Board of Inventions and Research, Canadian physicist Robert Boyle took on the project, which subsequently passed to the **Anti- (or Allied) Submarine Detection Investigation Committee**, producing a prototype for testing in mid-1917, hence the British acronym **ASDIC**.

By 1918, both the U.S. and Britain had built active systems. The UK tested what they still called ASDIC on *HMS Antrim* in 1920, and started production of units in 1922. The 6th Destroyer Flotilla had ASDIC-equipped vessels in 1923. An anti-submarine school, *HMS Osprey*, and a training flotilla of four vessels were established on Portland in 1924.

The U.S. Sonar QB set arrived in 1931. By the outbreak of World War II, the Royal Navy had five sets for different surface ship classes, and others for submarines. The greatest advantage came when it was linked to the Squid anti-submarine weapon.

Active sonar

Active sonar creates a pulse of sound, often called a "ping", and then listens for reflections of the pulse. To measure the distance to an object, one measures the time from emission of a pulse to reception. To measure the bearing, one uses several hydrophones, and measures the relative arrival time to each in a process called beamforming.

The pulse may be at constant frequency or a chirp of changing frequency. For a chirp, the receiver correlates the frequency of the reflections to the known chirp. The resultant processing gain allows the receiver to derive the same information as if a much shorter pulse of the same total energy were emitted. In practice, the chirp signal is sent over a longer time interval; therefore the instantaneous emitted power will be reduced, which simplifies the design of the transmitter. In general, long-distance active sonars use lower frequencies. The lowest have a bass "BAH-WONG" sound.

The most useful small sonar looks roughly like a waterproof flashlight. One points the head into the water, presses a button, and reads a distance. Another variant is a "fishfinder" that shows a small display with shoals of fish. Some civilian sonars approach active military sonars in capability, with quite exotic three-dimensional displays of the area near the boat. However, these sonars are not designed for stealth.

When active sonar is used to measure the distance to the bottom, it is known as echo sounding.

Active sonar is also used to measure distance through water between two sonar transponders. A transponder is a device that can transmit and receive signals but when it receives a specific interrogation signal it responds by transmitting a specific reply signal. To measure distance, one transponder transmits an interrogation signal and

measures the time between this transmission and the receipt of the other transponder's reply. The time difference, scaled by the speed of sound through water and divided by two, is the distance between the two transponders. This technique, when used with multiple transponders, can calculate the relative positions of static and moving objects in water.

Analysis of active sonar data

Active sonar data is obtained by measuring detected sound for a short period of time after the issuing of a ping; this time period is selected so as to ensure that the ping's reflection will be detected. The distance to the seabed (or other acoustically reflective object) can be calculated from the elapsed time between the ping and the detection of its reflection. Other properties can also be detected from the shape of the ping's reflection:

- When collecting data on the seabed, some of the reflected sound will typically reflect off the air-water interface, and then reflect off the seabed a second time. The size of this second echo provides information about the acoustic hardness of the seabed.
- The roughness of a seabed affects the variance in reflection time. For a smooth seabed, all of the reflected sound will take much the same path, resulting in a sharp spike in the data. For a rougher seabed, sound will be reflected back over a larger area of seabed, and some sound may bounce between seabed features before reflecting to the surface. A less sharp spike in the data therefore indicates a rougher seabed.

Sonar and marine animals

Some marine animals, such as whales and dolphins, use echolocation systems similar to active sonar to locate predators and prey. It is feared that sonar transmitters could confuse these animals and cause them to lose their way, perhaps preventing them from feeding and mating. A recent article on the BBC Web site (see below) reports findings published in the journal *Nature* to the effect that military sonar may be inducing some whales to experience decompression sickness (and resultant beachings).

High-powered sonar transmitters may indirectly harm marine animals, although scientific evidence suggests that a confluence of factors must first be present. In the Bahamas in 2000, a trial by the United States Navy of a 230 decibel transmitter in the frequency range 3 – 7 kHz resulted in the beaching of sixteen whales, seven of which were found dead. The Navy accepted blame in a report published in the *Boston Globe* on 1 January, 2002. However, at low powers, sonar can protect marine mammals against collisions with ships.

A kind of sonar called mid-frequency sonar has been correlated with mass cetacean strandings throughout the world's oceans, and has therefore been singled out by environmentalists as causing the death of marine mammals. International press coverage of these events can be found at this active sonar news clipping ^[1] Web site. A lawsuit was filed in Santa Monica, California on 19 October, 2005 contending that the U.S. Navy has conducted sonar exercises in violation of several environmental laws, including the National Environmental Policy Act, the Marine Mammal Protection Act, and the Endangered Species Act.

Passive sonar

Passive sonar listens without transmitting. It is usually employed in military settings, although a few are used in science applications.

Speed of sound

Sonar operation is affected by sound speed. Sound speed is slower in fresh water than in sea water. In all water sound velocity is affected by density (or the mass per unit of volume). Density is affected by temperature, dissolved molecules (usually salinity), and pressure. The speed of sound (in feet per second) is approximately equal to $4388 + (11.25 \times \text{temperature (in } ^\circ\text{F)}) + (0.0182 \times \text{depth (in feet)} + \text{salinity (in parts-per-thousand)})$. This is an empirically

derived approximation equation that is reasonably accurate for normal temperatures, concentrations of salinity and the range of most ocean depths. Ocean temperature varies with depth, but at between 30 and 100 metres there is often a marked change, called the thermocline, dividing the warmer surface water from the cold, still waters that make up the rest of the ocean. This can frustrate sonar, for a sound originating on one side of the thermocline tends to be bent, or refracted, off the thermocline. The thermocline may be present in shallower coastal waters, however, wave action will often mix the water column and eliminate the thermocline. Water pressure also affects sound propagation. Increased pressure increases the density of the water and raises the sound velocity. Increases in sound velocity cause the sound waves to refract away from the area of higher velocity. The mathematical model of refraction is called Snell's law.

Sound waves that are radiated down into the ocean bend back up to the surface in great arcs due to the effect of pressure on sound. The ocean must be at least 6000 feet (1850 meters) deep, or the sound waves will echo off the bottom instead of refracting back upwards. Under the right conditions these waves will then be focused near the surface and refracted back down and repeat another arc. Each arc is called a convergence zone. Where an arc intersects the surface a CZ annulus is formed. The diameter of the CZ depends on the temperature and salinity of the water. In the North Atlantic, for example, CZs are found approximately every 33 nautical miles (61 km), depending on the season, forming a pattern of concentric circles around the sound source. Sounds that can be detected for only a few miles in a direct line can therefore also be detected hundreds of miles away. Typically the first, second and third CZ are fairly useful; further out than that the signal is too weak, and thermal conditions are too unstable, reducing the reliability of the signals. The signal is naturally attenuated by distance, but modern sonar systems are very sensitive.

Identifying sound sources

Military sonar has a wide variety of techniques for identifying a detected sound. For example, U.S. vessels usually operate 60 Hz alternating current power systems. If transformers are mounted without proper vibration insulation from the hull, or flooded, the 60 Hz sound from the windings and generators can be emitted from the submarine or ship, helping to identify its nationality. In contrast, most European submarines have 50 Hz power systems. Intermittent noises (such as a wrench being dropped) may also be detectable to sonar.

Passive sonar systems may have large sonic databases, however most classification is performed manually by the sonar operator. A computer system frequently uses these databases to identify classes of ships, actions (i.e., the speed of a ship, or the type of weapon released), and even particular ships. Publications for classification of sounds are provided by and continually updated by the U.S. Office of Naval Intelligence.

Sonar in warfare

Modern naval warfare makes extensive use of sonar. The two types described before are both used, but from different platforms, i.e., types of water-borne vessels.

Active sonar is extremely useful, since it gives the exact position of an object. Active sonar works the same way as radar: a signal is emitted. The sound wave then travels in many directions from the emitting object. When it hits an object, the sound wave is then reflected in many other directions. Some of the energy will travel back to the emitting source. The echo will enable the sonar system or technician to calculate, with many factors such as the frequency, the energy of the received signal, the depth, the water temperature, etc., the position of the reflecting object. Using active sonar is somewhat hazardous however, since it does not allow the sonar to identify the target, and any vessel around the emitting sonar will detect the emission. Having heard the signal, it is easy to identify the type of sonar (usually with its frequency) and its position (with the sound wave's energy). Moreover, active sonar, similar to radar, allows the user to detect objects at a certain range but also enables other platforms to detect the active sonar at a far greater range.

Since active sonar does not allow an exact identification and is very noisy, this type of detection is used by fast platforms (planes, helicopters) and by noisy platforms (most surface ships) but rarely by submarines. When active sonar is used by surface ships or submarines, it is typically activated very briefly at intermittent periods, to reduce the risk of detection by an enemy's passive sonar. As such, active sonar is normally considered a backup to passive sonar. In aircraft, active sonar is used in the form of disposable sonobuoys that are dropped in the aircraft's patrol area or in the vicinity of possible enemy sonar contacts.

Passive sonar has fewer drawbacks. Most importantly, it is silent. Generally, it has a much greater range than active sonar, and allows an identification of the target. Since any motorized object makes some noise, it may be detected eventually. It simply depends on the amount of noise emitted and the amount of noise in the area, as well as the technology used. To simplify, passive sonar "sees" around the ship using it. On a submarine, the nose mounted passive sonar detects in directions of about 270° , centered on the ship's alignment, the hull-mounted array of about 160° on each side, and the towed array of a full 360° . The no-see areas are due to the ship's own interference. Once a signal is detected in a certain direction (which means that something makes sound in that direction, this is called broadband detection) it is possible to zoom in and analyze the signal received (narrowband analysis). This is generally done using a Fourier transform to show the different frequencies making up the sound. Since every engine makes a specific noise, it is easy to identify the object.

Another use of the passive sonar is to determine the target's trajectory. This process is called Target Motion Analysis (TMA), and the resultant "solution" is the target's range, course, and speed. TMA is done by marking from which direction the sound comes at different times, and comparing the motion with that of the operator's own ship. Changes in relative motion are analyzed using standard geometrical techniques along with some assumptions about limiting cases.

Passive sonar is stealthy and very useful. However, it requires high-tech components (band pass filters, receivers) and is costly. It is generally deployed on expensive ships in the form of arrays to enhance the detection. Surface ships use it to good effect; it is even better used by submarines, and it is also used by airplanes and helicopters, mostly to a "surprise effect", since submarines can hide under thermal layers. If a submarine captain believes he is alone, he may bring his boat closer to the surface and be easier to detect, or go deeper and faster, and thus make more sound.

In the United States Navy, a special badge known as the Integrated Undersea Surveillance System Badge is awarded to those who have been trained and qualified in sonar operation and warfare.

In World War II, the Americans used the term **SONAR** for their system. The British still called their system **ASDIC**. In 1948, with the formation of NATO, standardization of signals led to the dropping of ASDIC in favor of sonar.

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- [1] http://www.anon.org/lfas_news.jsp

Interior Sound Transmission

Introduction to NVH

Noise is characterized by frequency (20-20kHz), level (dB) and quality. Noise may be undesirable in some cases, i.e. road NVH yet may be desirable in other cases, i.e. powerful sounding engine.

Vibration is defined as the motion sensed by the body, mainly in 0.5Hz - 50Hz range. It is characterized by frequency, level and direction.

Harshness is defined as rough, grating or discordant sensation.

Sound quality is defined according to Oxford English Dictionary, "That distinctive Quality of a Sound other than it's Pitch or Loudness"

In generally, ground vehicle NVH and sound quality design are attributed by the following categories:

- 1.) Powertrain NVH and SQ: Interior Idle NVH, Acceleration NVH, Deceleration NVH, Cruising NVH, Sound Quality Character, Diesel Combustion Noise, Engine Start-Up/Shut-Down
- 2.) Wind Noise: Motorway Speed Wind Noise (80-130kph), High Speed Wind Noise (>130kph), Open Glazing Wind Noise
- 3.) Road NVH: Road Noise, Road Vibration, Impact Noise
- 4.) Operational Sound Quality: Closure Open/Shut Sound Quality, Customer Operated Feature Sound Quality, Audible Warning Sounds
- 5.) Squeaks and Rattles

== Noise generation ==

Applications in Room Acoustics

Anechoic and Reverberation Rooms



Introduction

Acoustic experiments often require to realise measurements in rooms with special characteristics. Two types of rooms can be distinguished: anechoic rooms and reverberation rooms.

Anechoic room

The principle of this room is to simulate a free field. In a free space, the acoustic waves are propagated from the source to infinity. In a room, the reflections of the sound on the walls produce a wave which is propagated in the opposite direction and comes back to the source. In anechoic rooms, the walls are very absorbent in order to eliminate these reflections. The sound seems to die down rapidly. The materials used on the walls are rockwool, glasswool or foams, which are materials that absorb sound in relatively wide frequency bands. Cavities are dug in the wool so that the large wavelength corresponding to bass frequencies are absorbed too. Ideally the sound pressure level of a punctual sound source decreases about 6 dB per a distance doubling.

Anechoic rooms are used in the following experiments:

Intensimetry: measurement of the acoustic power of a source.

Study of the source directivity.

Reverberation room

The walls of a reverberation room mostly consist of concrete and are covered with reflecting paint. Alternative design consist of sandwich panels with metal surface. The sound reflects a lot of time on the walls before dying down. It gives a similar impression of a sound in a cathedral. Ideally all sound energy is absorbed by air. Because of all these reflections, a lot of plane waves with different directions of propagation interfere in each point of the room. Considering all the waves is very complicated so the acoustic field is simplified by the diffuse field hypothesis: the

field is homogeneous and isotropic. Then the pressure level is uniform in the room. The truth of this thesis increases with ascending frequency, resulting in a lower limiting frequency for each reverberation room, where the density of standing waves is sufficient.

Several conditions are required for this approximation: The absorption coefficient of the walls must be very low ($\alpha < 0.2$) The room must have geometrical irregularities (non-parallel walls, diffusor objects) to avoid nodes of pressure of the resonance modes.

With this hypothesis, the theory of Sabine can be applied. It deals with the reverberation time which is the time required to the sound level to decrease of 60dB. T depends on the volume of the room V, the absorption coefficient α_i and the area S_i of the different materials in the room :

Reverberation rooms are used in the following experiments:

- measurement of the ability of a material to absorb a sound
- measurement of the ability of a partition to transmit a sound
- Intensimetry
- measurement of sound power

Basic Room Acoustic Treatments

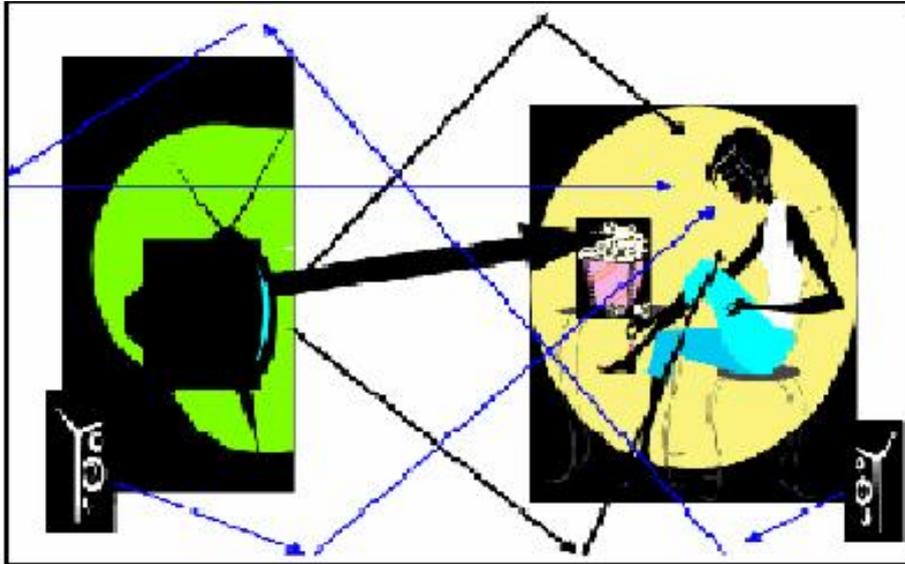


Introduction

Many people use one or two rooms in their living space as "theatrical" rooms where theater or music room activities commence. It is a common misconception that adding speakers to the room will enhance the quality of the room acoustics. There are other simple things that can be done to increase the room's acoustics to produce sound that is similar to "theater" sound. This site will take you through some simple background knowledge on acoustics and then explain some solutions that will help improve sound quality in a room.

Room sound combinations

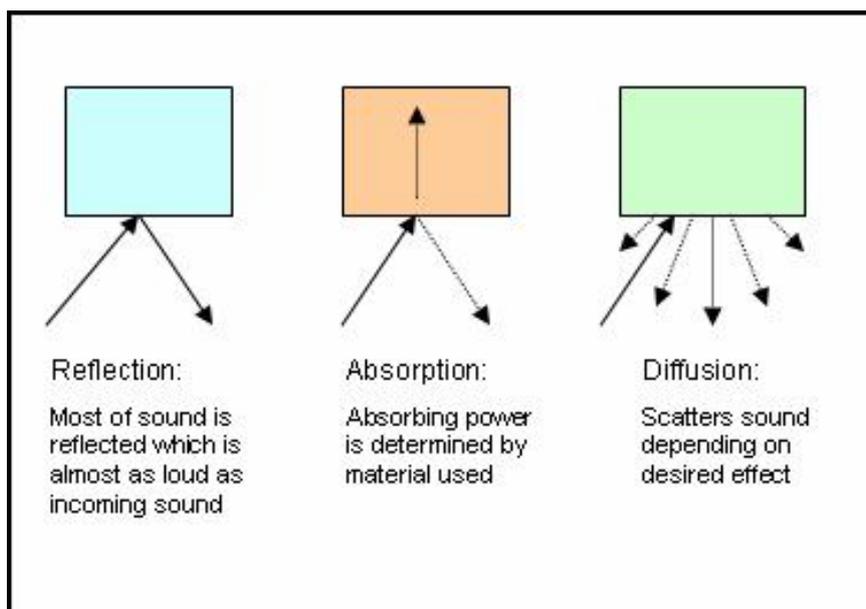
The sound you hear in a room is a combination of direct sound and indirect sound. Direct sound will come directly from your speakers while the other sound you hear is reflected off of various objects in the room.



The Direct sound is coming right out of the TV to the listener, as you can see with the heavy black arrow. All of the other sound is reflected off surfaces before they reach the listener.

Good and bad reflected sound

Have you ever listened to speakers outside? You might have noticed that the sound is thin and dull. This occurs because when sound is reflected, it is fuller and louder than it would be if it were in an open space. So when sound is reflected, it can add a fullness, or spaciousness. The bad part of reflected sound occurs when the reflections amplify some notes, while cancelling out others, making the sound distorted. It can also affect tonal quality and create an echo-like effect. There are three types of reflected sound, pure reflection, absorption, and diffusion. Each reflection type is important in creating a "theater" type acoustic room.



Reflected sound

Reflected sound waves, good and bad, affect the sound you hear, where it comes from, and the quality of the sound when it gets to you. The bad news when it comes to reflected sound is standing waves.

These waves are created when sound is reflected back and forth between any two parallel surfaces in your room, ceiling and floor or wall to wall.

Standing waves can distort noises 300Hz and down. These noises include the lower mid frequency and bass ranges. Standing waves tend to collect near the walls and in corners of a room, these collecting standing waves are called room resonance modes.

Finding your room resonance modes

First, specify room dimensions (length, width, and height). Then follow this example:

To Find the Main Resonance Mode :

$$F = \frac{c}{2 * D}$$

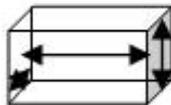
c = speed of sound (air = 1130 ft/s)

F = Frequency of Resonance Mode (Hz)

D = Distance between parallel walls (ft)

EXAMPLE

Room (23'*11'*8')



$$F = \frac{1130}{2 * 23} = 25\text{Hz}$$

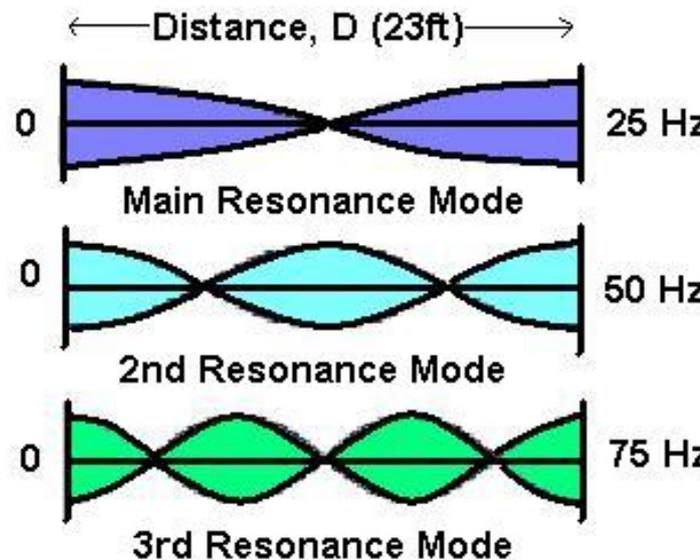
Main Mode = F = 25Hz

Weaker Modes at multiples of Main Mode

$$2 * 25, 3 * 25 \dots$$

Do this for each room dimension (l, w, h)

Resonance Mode Frequencies appearing in more than one dimension will be trouble frequencies.



Working with room resonance modes to increase sound quality

There are some room dimensions that produce the largest amount of standing waves.

1. Cube
2. Room with 2 out of the three dimensions equal
3. Rooms with dimensions that are multiples of each other

Absorbed

The sound that humans hear is actually a form of acoustic energy. Different materials absorb different amounts of this energy at different frequencies. When considering room acoustics, there should be a good mix of high frequency absorbing materials and low frequency absorbing materials. A table including information on how different common household absorb sound can be found here ^[1].

Diffused sound

Using devices that diffuse sound is a fairly new way of increasing acoustic performance in a room. It is a means to create sound that appears to be "live". They can replace echo-like reflections without absorbing too much sound.

Some ways of determining where diffusive items should be placed were found on this website ^[2].

1. If you have carpet or drapes already in your room, use diffusion to control side wall reflections.
2. A bookcase filled with odd-sized books makes an effective diffuser.
3. Use absorptive material on room surfaces between your listening position and your front speakers, and treat the back wall with diffusive material to re-distribute the reflections.

How to find overall trouble spots in a room

Every surface in a room does not have to be treated in order to have good room acoustics. Here is a simple method of finding trouble spots in a room.

1. Grab a friend to hold a mirror along the wall near a certain speaker at speaker height.
2. The listener sits in a spot of normal viewing.
3. The friend then moves slowly toward the listening position (stay along the wall).
4. Mark each spot on the wall where the listener can see any of the room speakers in the mirror.
5. Congratulations! These are the trouble spots in the room that need an absorptive material in place. Don't forget that diffusive material can also be placed in those positions.

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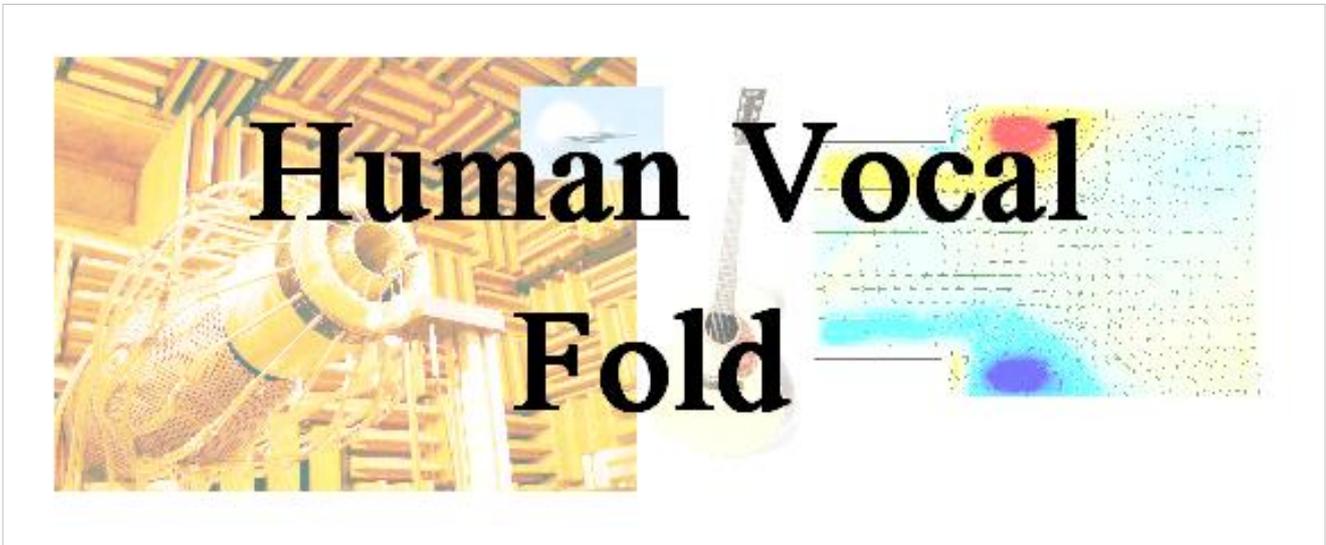
- Acoustic Room Treatment Articles ^[3]
- Room Acoustics: Acoustic Treatments ^[4]
- Home Improvement: Acoustic Treatments ^[5]
- Crutchfield Advisor ^[6]

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- [5] http://www.diynetwork.com/diy/hi_family_room/article/0,2037,DIY_13912_3471072,00.html
- [6] http://www.crutchfieldadvisor.com/S-hpU9sw2hgbG/learningcenter/home/speakers_roomacoustics.html?page=1

Applications in Psychoacoustics

Human Vocal Fold



Physiology of vocal fold

Human vocal fold is a set of lip-like tissues located inside the larynx ^[1], and is the source of sound for a human and many animals.

The Larynx is located at the top of trachea. It is mainly composed of cartilages and muscles, and the largest cartilage, thyroid, is well known as the "Adam's Apple."

The organ has two main functions; to act as the last protector of the airway, and to act as a sound source for voice. This page focuses on the latter function.

Links on Physiology:

- [Discover The Larynx](#) ^[2]

Voice production

Although the science behind sound production for a vocal fold is complex, it can be thought of as similar to a brass player's lips, or a whistle made out of grass. Basically, vocal folds (or lips or a pair of grass) make a constriction to the airflow, and as the air is forced through the narrow opening, the vocal folds oscillate. This causes a periodical change in the air pressure, which is perceived as sound.

Vocal Folds Video ^[3]

When the airflow is introduced to the vocal folds, it forces open the two vocal folds which are nearly closed initially. Due to the stiffness of the folds, they will then try to close the opening again. And now the airflow will try to force the folds open etc... This creates an oscillation of the vocal folds, which in turn, as I stated above, creates sound. However, this is a damped oscillation, meaning it will eventually achieve an equilibrium position and stop oscillating. So how are we able to "sustain" sound?

As it will be shown later, the answer seems to be in the changing shape of vocal folds. In the opening and the closing stages of the oscillation, the vocal folds have different shapes. This affects the pressure in the opening, and creates the extra pressure needed to push the vocal folds open and sustain oscillation. This part is explained in more detail in the "Model" section.

This flow-induced oscillation, as with many fluid mechanics problems, is not an easy problem to model. Numerous attempts to model the oscillation of vocal folds have been made, ranging from a single mass-spring-damper system to finite element models. In this page I would like to use my single-mass model to explain the basic physics behind the oscillation of a vocal fold.

Information on vocal fold models: National Center for Voice and Speech ^[4]

Model

The most simple way of simulating the motion of vocal folds is to use a single mass-spring-damper system as shown above. The mass represents one vocal fold, and the second vocal fold is assumed to be symmetry about the axis of symmetry. Position 3 represents a location immediately past the exit (end of the mass), and position 2 represents the glottis (the region between the two vocal folds).

The pressure force

The major driving force behind the oscillation of vocal folds is the pressure in the glottis. The Bernoulli's equation from fluid mechanics states that:

$$P_1 + \frac{1}{2}\rho U^2 + \rho gh = \text{Constant} \text{ -----EQN 1}$$

Neglecting potential difference and applying EQN 1 to positions 2 and 3 of Figure 1,

$$P_2 + \frac{1}{2}\rho U_2^2 = P_3 + \frac{1}{2}\rho U_3^2 \text{ -----EQN 2}$$

Note that the pressure and the velocity at position 3 cannot change. This makes the right hand side of EQN 2 constant. Observation of EQN 2 reveals that in order to have oscillating pressure at 2, we must have oscillation velocity at 2. The flow velocity inside the glottis can be studied through the theories of the orifice flow.

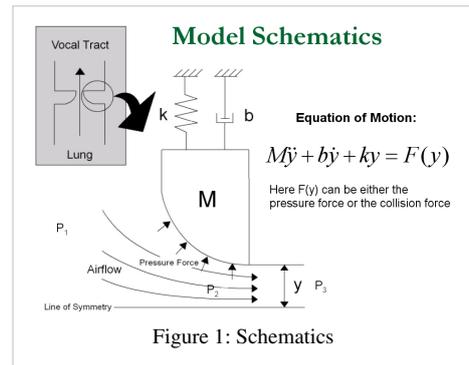
The constriction of airflow at the vocal folds is much like an orifice flow with one major difference: with vocal folds, the orifice profile is continuously changing. The orifice profile for the vocal folds can open or close, as well as change the shape of the opening. In Figure 1, the profile is converging, but in another stage of oscillation it takes a diverging shape.

The orifice flow is described by Blevins as:

$$U = C \frac{2(P_1 - P_3)}{\rho} \text{ -----EQN 3}$$

Where the constant C is the orifice coefficient, governed by the shape and the opening size of the orifice. This number is determined experimentally, and it changes throughout the different stages of oscillation.

Solving equations 2 and 3, the pressure force throughout the glottal region can be determined.



The Collision Force

As the video of the vocal folds shows, vocal folds can completely close during oscillation. When this happens, the Bernoulli equation fails. Instead, the collision force becomes the dominating force. For this analysis, Hertz collision model was applied.

$$F_H = k_H \delta^{3/2} (1 + b_H \delta') \text{-----EQN 4}$$

where

$$k_H = \frac{4}{3} \frac{E}{1 - \mu_H^2} \sqrt{r}$$

Here delta is the penetration distance of the vocal fold past the line of symmetry.

Simulation of the model

The pressure and the collision forces were inserted into the equation of motion, and the result was simulated.

Figure 2 shows that an oscillating volumetric flow rate was achieved by passing a constant airflow through the vocal folds. When simulating the oscillation, it was found that the collision force limits the amplitude of oscillation rather than drive the oscillation. Which tells us that the pressure force is what allows the sustained oscillation to occur.

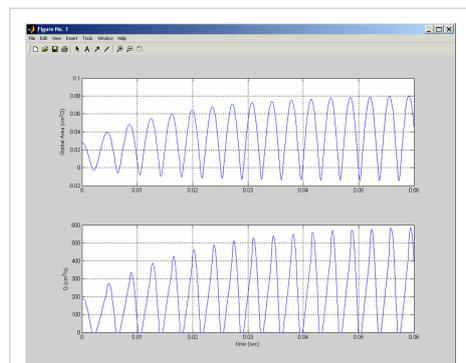


Figure 2: Area Opening and Volumetric Flow Rate

The acoustic output

This model showed that the changing profile of glottal opening causes an oscillating volumetric flow rate through the vocal folds. This will in turn cause an oscillating pressure past the vocal folds. This method of producing sound is unusual, because in most other means of sound production, air is compressed periodically by a solid such as a speaker cone.

Past the vocal folds, the produced sound enters the vocal tract. Basically this is the cavity in the mouth as well as the nasal cavity. These cavities act as acoustic filters, modifying the character of the sound. These are the characters that define the unique voice each person produces.

Related links

- FEA Model ^[5]
- Two Mass Model ^[6]

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Threshold of Hearing & Pain

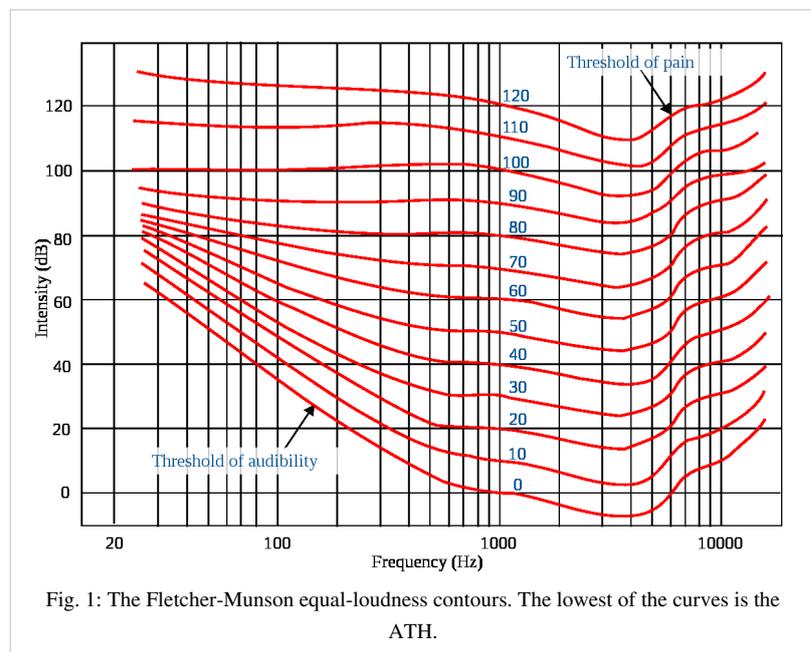


The **threshold of hearing** is the Sound pressure level SPL of 20 μPa (micropascals) = 2×10^{-5} pascal (Pa). This low threshold of amplitude (strength or sound pressure level) is frequency dependent. See the frequency curve in Fig. 2 below

The **absolute threshold of hearing** (ATH) is the minimum amplitude (level or strength) of a pure tone that the average ear with normal hearing can hear in a noiseless environment.

The **threshold of pain** is the SPL beyond which sound becomes unbearable for a human listener. This threshold varies only slightly with frequency. Prolonged exposure to sound pressure levels in excess of the threshold of pain can cause physical damage, potentially leading to hearing impairment.

Different values for the threshold of pain:



Threshold of pain	
SPL	sound pressure
120 dBSPL	20 Pa
130 dBSPL	63 Pa
134 dBSPL	100 Pa
137.5 dBSPL	150 Pa
140 dBSPL	200 Pa

The **Threshold of hearing** is frequency dependent, and typically shows a minimum (indicating the ear's maximum sensitivity) at frequencies between 1 kHz and 5 kHz. A typical ATH curve is pictured in Fig. 1. The absolute threshold of hearing represents the lowest curve amongst the set of equal-loudness contours, with the highest curve representing the threshold of pain.

In psychoacoustic audio compression, the ATH is used, often in combination with masking curves, to calculate which spectral components are inaudible and may thus be ignored in the coding process; any part of an audio spectrum which has an amplitude (level or strength) below the ATH may be removed from an audio signal without any audible change to the signal.

The ATH curve rises with age as the human ear becomes more insensitive to sound, with the greatest changes occurring at frequencies higher than 2 kHz. Curves for subjects of various age groups are illustrated in Fig. 2. The data is from the United States *Occupational Health and Environment Control, Standard Number: 1910.95 App F*

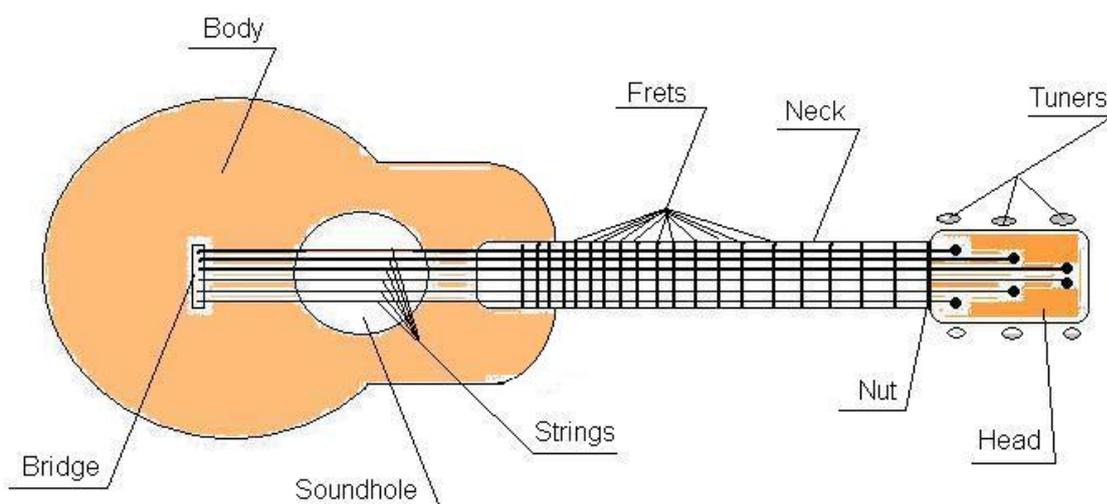
../Human Vocal Fold/ - **Acoustics** - ../How an Acoustic Guitar Works/

Musical Acoustics Applications

How an Acoustic Guitar Works



What are sound vibrations that contribute to sound production? First of all, there are the strings. Any string that is under tension will vibrate at a certain frequency. The weight and length of the string, the tension in the string, and the compliance of the string determine the frequency at which it vibrates. The guitar controls the length and tension of six differently weighted strings to cover a very wide range of frequencies. Second, there is the body of the guitar. The guitar body is connected is very important for the lower frequencies of the guitar. The air mass just inside the sound hole oscillates, compressing and decompressing the compliant air inside the body. In practice this concept is called a Helmholtz resonator. Without this, it would difficult to produce the wonderful timbre of the guitar.



The strings

The strings of the guitar vary in linear density, length, and tension. This gives the guitar a wide range of attainable frequencies. The larger the linear density is, the slower the string vibrates. The same goes for the length; the longer the string is the slower it vibrates. This causes a low frequency. Inversely, if the strings are less dense and/or shorter they create a higher frequency. The resonance frequencies of the strings can be calculated by

$$f_1 = \frac{\sqrt{T}}{2L\rho_1}$$

T = String tension
 ρ_1 = Linear density
 L = String length

The string length, L , in the equation is what changes when a player presses on a string at a certain fret. This will shorten the string which in turn increases the frequency it produces when plucked. The spacing of these frets is important. The length from the nut to bridge determines how much space goes between each fret. If the length is 25 inches, then the position of the first fret should be located $(25/17.817)$ inches from the nut. Then the second fret should be located $(25-(25/17.817))/17.817$ inches from the first fret. This results in the equation

$$d = \frac{L}{17.817}$$

d = Spacing between frets
 L = Length from previous fret to bridge

When a string is plucked, a disturbance is formed and travels in both directions away from point where the string was plucked. These "waves" travel at a speed that is related to the tension and linear density and can be calculated by

$$c = \sqrt{\frac{T}{\rho_1}}$$

c = Wave speed
 T = String tension
 ρ_1 = Linear density

The waves travel until they reach the boundaries on each end where they are reflected back. The link below displays how the waves propagate in a string.

Plucked String @ www.phys.unsw.edu ^[1]

The strings themselves do not produce very much sound because they are so thin. They can't "push" the air that surrounds them very effectively. This is why they are connected to the top plate of the guitar body. They need to transfer the frequencies they are producing to a large surface area which can create more intense pressure disturbances.

The body

The body of the guitar transfers the vibrations of the bridge to the air that surrounds it. The top plate contributes to most of the pressure disturbances, because the player dampens the back plate and the sides are relatively stiff. This is why it is important to make the top plate out of a light springy wood, like spruce. The more the top plate can vibrate, the louder the sound it produces will be. It is also important to keep the top plate flat, so a series of braces are located on the inside to strengthen it. Without these braces the top plate would bend and crack under the large stress created by the tension in the strings. This would also affect the magnitude of the sound being transmitted. The warped plate would not be able to "push" air very efficiently. A good experiment to try, in order to see how important this part of the guitar is in the amplification process, is as follows:

1. Start with an ordinary rubber band, a large bowl, adhesive tape, and plastic wrap.
2. Stretch the rubber band and pluck it a few times to get a good sense for how loud it is.
3. Stretch the plastic wrap over the bowl to form a sort of drum.
4. Tape down one end of the rubber band to the plastic wrap.
5. Stretch the rubber band and pluck it a few times.
6. The sound should be much louder than before.

The air

The final part of the guitar is the air inside the body. This is very important for the lower range of the instrument. The air just inside the soundhole oscillates compressing and expanding the air inside the body. This is just like blowing across the top of a bottle and listening to the tone it produces. This forms what is called a Helmholtz resonator. For more information on Helmholtz resonators go to Helmholtz Resonance ^[2]. This link also shows the correlation to acoustic guitars in great detail. The acoustic guitar makers often tune these resonators to have a resonance frequency between F#2 and A2 (92.5 to 110.0 Hz). Having such a low resonance frequency is what aids the amplification of the lower frequency strings. To demonstrate the importance of the air in the cavity, simply play an open A on the guitar (the fifth string - second lowest note). Now, as the string is vibrating, place a piece of cardboard over the soundhole. The sound level is reduced dramatically. This is because you've stopped the vibration of the air mass just inside the soundhole, causing only the top plate to vibrate. Although the top plate still vibrates a transmits sound, it isn't as effective at transmitting lower frequency waves, thus the need for the Helmholtz resonator.

References

- [1] <http://www.phys.unsw.edu.au/music/objects/waves&strings2.swf>
- [2] <http://www.phys.unsw.edu.au/~jw/Helmholtz.html>

Basic Acoustics of the Marimba



Introduction

One of my favorite instruments is the marimba. Like a xylophone, a marimba has octaves of wooden bars that are struck with mallets to produce tones. Unlike the harsh sound of a xylophone, a marimba produces a deep, rich tone. Marimbas are not uncommon and are played in most high school bands.

Now, while all the trumpet and flute and clarinet players are busy tuning up their instruments, the marimba player is back in the percussion section with her feet up just relaxing. This is a bit surprising, however, since the marimba is a melodic instrument that needs to be in tune to sound good. So what gives? Why is the marimba never tuned? How would you even go about tuning a marimba? To answer these questions, the acoustics behind (or within) a marimba must be understood.

Components of sound

What gives the marimba its unique sound? It can be boiled down to two components: the bars and the resonators. Typically, the bars are made of rosewood (or some synthetic version of wood). They are cut to size depending on what note is desired, then the tuning is refined by shaving wood from the underside of the bar.

Example

Rosewood bar, middle C, 1 cm thick

The equation that relates the length of the bar with the desired frequency comes from the theory of modeling a bar that is free at both ends. This theory yields the following equation:

$$Length = \sqrt{\frac{3.011^2 \cdot \pi \cdot t \cdot c}{8 \cdot \sqrt{12} \cdot f}}$$

where t is the thickness of the bar, c is the speed of sound in the bar, and f is the frequency of the note.

For rosewood, $c = 5217$ m/s. For middle C, $f = 262$ Hz.

Therefore, to make a middle C key for a rosewood marimba, cut the bar to be:

$$Length = \sqrt{\frac{3.011^2 \cdot \pi \cdot .01 \cdot 5217}{8 \cdot \sqrt{12} \cdot 262}} = .45m = 45cm^{***}$$

The resonators are made from metal (usually aluminum) and their lengths also differ depending on the desired note. It is important to know that each resonator is open at the top but closed by a stopper at the bottom end.

Example

Aluminum resonator, middle C

The equation that relates the length of the resonator with the desired frequency comes from modeling the resonator as a pipe that is driven at one end and closed at the other end. A "driven" pipe is one that has a source of excitation (in this case, the vibrating key) at one end. This model yields the following:

$$Length = \frac{c}{4 \cdot f}$$

where c is the speed of sound in air and f is the frequency of the note.

For air, $c = 343$ m/s. For middle C, $f = 262$ Hz.

Therefore, to make a resonator for the middle C key, the resonator length should be:

$$Length = \frac{343}{4 \cdot 262} = .327m = 32.7cm$$

Resonator shape

The shape of the resonator is an important factor in determining the quality of sound that can be produced. The ideal shape is a sphere. This is modeled by the Helmholtz resonator. However, mounting big, round, beach ball-like resonators under the keys is typically impractical. The worst choices for resonators are square or oval tubes. These shapes amplify the non-harmonic pitches sometimes referred to as "junk pitches". The round tube is typically chosen because it does the best job (aside from the sphere) at amplifying the desired harmonic and not much else.

As mentioned in the second example above, the resonator on a marimba can be modeled by a closed pipe. This model can be used to predict what type of sound (full and rich vs dull) the marimba will produce. Each pipe is a "quarter wave resonator" that amplifies the sound waves produced by of the bar. This means that in order to produce a full, rich sound, the length of the resonator must exactly match one-quarter of the wavelength. If the length is off, the marimba will produce a dull or off-key sound for that note.

Why would the marimba need tuning?

In the theoretical world where it is always 72 degrees with low humidity, a marimba would not need tuning. But, since weather can be a factor (especially for the marching band) marimbas do not always perform the same way. Hot and cold weather can wreak havoc on all kinds of percussion instruments, and the marimba is no exception. On hot days, the marimba tends to be sharp and for cold days it tends to be flat. This is the exact opposite of what happens to string instruments. Why? The tone of a string instrument depends mainly on the tension in the string, which decreases as the string expands with heat. The decrease in tension leads to a flat note. Marimbas on the other hand produce sound by moving air through the resonators. The speed at which this air is moved is the speed of sound, which varies proportionately with temperature! So, as the temperature increases, so does the speed of sound. From the equation given in example 2 from above, you can see that an increase in the speed of sound (c) means a longer pipe is needed to resonate the same note. If the length of the resonator is not increased, the note will sound sharp. Now, the heat can also cause the wooden bars to expand, but the effect of this expansion is insignificant compared to the effect of the change in the speed of sound.

Tuning myths

It is a common myth among percussionists that the marimba can be tuned by simply moving the resonators up or down (while the bars remain in the same position.) The thought behind this is that by moving the resonators down, for example, you are in effect lengthening them. While this may sound like sound reasoning, it actually does not hold true in practice. Judging by how the marimba is constructed (cutting bars and resonators to specific lengths), it seems that there are really two options to consider when looking to tune a marimba: shave some wood off the underside of the bars, or change the length of the resonator. For obvious reasons, shaving wood off the keys every time the weather changes is not a practical solution. Therefore, the only option left is to change the length of the resonator.

As mentioned above, each resonator is plugged by a stopper at the bottom end. So, by simply shoving the stopper farther up the pipe, you can shorten the resonator and sharpen the note. Conversely, pushing the stopper down the pipe can flatten the note. Most marimbas do not come with tunable resonators, so this process can be a little challenging. (Broomsticks and hammers are common tools of the trade.)

Example

Middle C Resonator lengthened by 1 cm

For ideal conditions, the length of the middle C (262 Hz) resonator should be 32.7 cm as shown in example 2. Therefore, the change in frequency for this resonator due to a change in length is given by:

$$\Delta Frequency = 262Hz - \frac{c}{4 \cdot (.327 + \Delta L)}$$

If the length is increased by 1 cm, the change in frequency will be:

$$\Delta Frequency = \frac{343}{4 \cdot (.327 + .01)} - 262Hz = 7.5Hz$$

The acoustics behind the tuning a marimba go back to the design that each resonator is to be $\frac{1}{4}$ of the total wavelength of the desired note. When marimbas get out of tune, this length is no longer exactly equal to $\frac{1}{4}$ the wavelength due to the lengthening or shortening of the resonator as described above. Because the length has changed, resonance is no longer achieved, and the tone can become muffled or off-key.

Conclusions

Some marimba builders are now changing their designs to include tunable resonators. Since any leak in the end-seal will cause major loss of volume and richness of the tone, this is proving to be a very difficult task. At least now, though, armed with the acoustic background of their instruments, percussionists everywhere will now have something to do when the conductor says, "tune up!"

Links and References

1. <http://www.gppercussion.com/html/resonators.html>
2. <http://www.mostlymarimba.com/>
3. <http://www.craftymusicteachers.com/bassmarimba/>

Bessel Functions and the Kettledrum



Introduction

In class, we have begun to discuss the solutions of multidimensional wave equations. A particularly interesting aspect of these multidimensional solutions are those of bessel functions for circular boundary conditions. The practical application of these solutions is the kettledrum. This page will explore in qualitative and quantitative terms how the kettledrum works. More specifically, the kettledrum will be introduced as a circular membrane and its solution will be discussed with visuals (e.g. visualization of bessel functions, video of kettledrums and audio forms (wav files of kettledrums playing. In addition, links to more information about this material, including references will be included.

The math behind the kettledrum: the brief version

When one looks at how a kettledrum produces sound, one should look no farther than the drum head. The vibration of this circular membrane (and the air in the drum enclosure) is what produces the sound in this instrument. The mathematics behind this vibrating drum are relatively simple. If one looks at a small element of the drum head, it looks exactly like the mathematical model for a vibrating string (see:). The only difference is that there are two dimensions where there are forces on the element, the two dimensions that are planar to the drum. As this is the same situation, we have the same equation, except with another spatial term in the other planar dimension. This allows us

to model the drum head using a helmholtz equation. The next step (solved in detail below) is to assume that the displacement of the drum head (in polar coordinates) is a product of two separate functions for theta and r. This allows us to turn the PDE into two ODES which are readily solved and applied to the situation of the kettledrum head. For more info, see below.

The math behind the kettledrum: the derivation

So starting with the trusty general Helmholtz equation:

$$\nabla^2 \Psi + k^2 \Psi = 0$$

Where k is the wave number, the frequency of the forced oscillations divided by the speed of sound in the membrane.

Since we are dealing with a circular object, it make sense to work in polar coordinates (in terms of radius and angle) instead of rectangular coordinates. For polar coordinates the Laplacian term of the helmholtz relation (∇^2) becomes $\partial^2 \Psi / \partial r^2 + 1/r \partial \Psi / \partial r + 1/r^2 \partial^2 \Psi / \partial \theta^2$

Now lets assume that: $\Psi(r, \theta) = R(r)\Theta(\theta)$

This assumption follows the method of separation of variables. (see Reference 3 for more info) Substituting this result back into our trusty Helmholtz equation gives the following:

$$r^2 / R(d^2 R / dr^2 + 1/r dR / dr) + k^2 r^2 = -1 / \Theta d^2 \Theta / d\theta^2$$

Since we separated the variables of the solution into two one-dimensional functions, the partial derivatives become ordinary derivatives. Both sides of this result must equal the same constant. For simplicity, I will use λ^2 as this constant. This results in the following two equations:

$$d^2 \Theta / d\theta^2 = -\lambda^2 \Theta$$

$$d^2 R / dr^2 + 1/r dR / dr + (k^2 - \lambda^2 / r^2) R = 0$$

The first of these equations readily seen as the standard second order ordinary differential equation which has a harmonic solution of sines and cosines with the frequency based on λ . The second equation is what is known as Bessel's Equation. The solution to this equation is cryptically called Bessel functions of order λ of the first and second kind. These functions, while sounding very intimidating, are simply oscillatory functions of the radius times the wave number that are unbounded when kr (for the function of the second kind) approaches zero and diminish as kr get larger. (For more information on what these functions look like see References 1,2, and 3)

Now that we have the general solution to this equation, we can now model a infinite radius kettledrum head. However, since i have yet to see an infinite kettle drum, we need to constrain this solution of a vibrating membrane to a finite radius. We can do this by applying what we know about our circular membrane: along the edges of the kettledrum, the drum head is attached to the drum. This means that there can be no displacement of the membrane at the termination at the radius of the kettle drum. This boundary condition can be mathematically described as the following:

$$R(a) = 0$$

Where a is the arbitrary radius of the kettledrum. In addition to this boundary condition, the displacement of the drum head at the center must be finite. This second boundary condition removes the bessel function of the second kind from the solution. This reduces the R part of our solution to:

$$R(r) = AJ_\lambda(kr)$$

Where J_λ is a bessel function of the first kind of order λ . Apply our other boundary condition at the radius of the drum requires that the wave number k must have discrete values, (j_{mn}/a) which can be looked up. Combining all of these gives us our solution to how a drum head behaves (which is the real part of the following):

$$y_{\lambda n}(r, \theta, t) = A_{\lambda n} J_{\lambda n}(k_{\lambda n} r) e^{j\lambda\theta + j\omega_{\lambda n} t}$$

The math behind the kettledrum: the entire drum

The above derivation is just for the drum head. An actual kettledrum has one side of this circular membrane surrounded by an enclosed cavity. This means that air is compressed in the cavity when the membrane is vibrating, adding more complications to the solution. In mathematical terms, this makes the partial differential equation non-homogeneous or in simpler terms, the right side of the Helmholtz equation does not equal zero. This result requires significantly more derivation, and will not be done here. If the reader cares to know more, these results are discussed in the two books under references 6 and 7.

Sites of interest

As one can see from the derivation above, the kettledrum is very interesting mathematically. However, it also has a rich historical music tradition in various places of the world. As this page's emphasis is on math, there are few links provided below that reference this rich history.

- A discussion of Persian kettledrums: Kettle drums of Iran and other countries ^[1]
- A discussion of kettledrums in classical music: Kettle drum Lit. ^[2]
- A massive resource for kettledrum history, construction and technique" Vienna Symphonic Library ^[3]

References

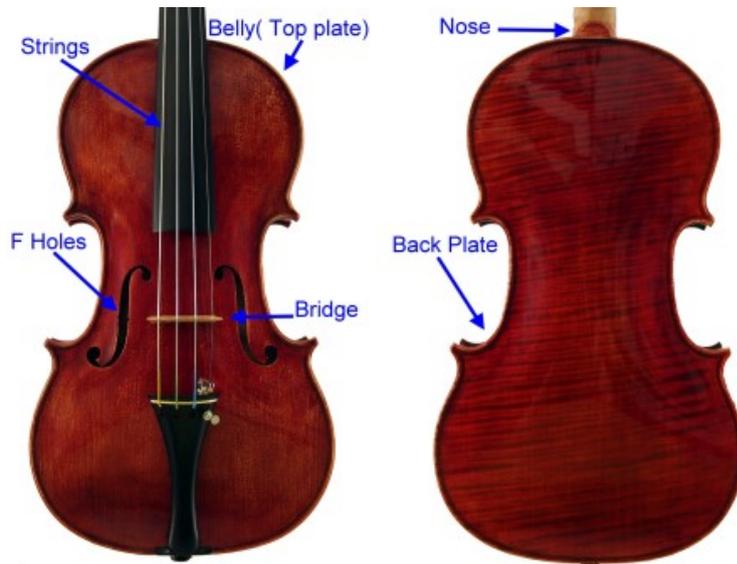
1. Eric W. Weisstein. "Bessel Function of the First Kind." From MathWorld--A Wolfram Web Resource. <http://mathworld.wolfram.com/BesselFunctionoftheFirstKind.html>
2. Eric W. Weisstein. "Bessel Function of the Second Kind." From MathWorld--A Wolfram Web Resource. <http://mathworld.wolfram.com/BesselFunctionoftheSecondKind.html>
3. Eric W. Weisstein. "Bessel Function." From MathWorld--A Wolfram Web Resource. <http://mathworld.wolfram.com/BesselFunction.html>
4. Eric W. Weisstein et al. "Separation of Variables." From MathWorld--A Wolfram Web Resource. <http://mathworld.wolfram.com/SeparationofVariables.html>
5. Eric W. Weisstein. "Bessel Differential Equation." From MathWorld--A Wolfram Web Resource. <http://mathworld.wolfram.com/BesselDifferentialEquation.html>
6. Kinsler and Frey, "Fundamentals of Acoustics", fourth edition, Wiley & Sons
7. Haberman, "Applied Partial Differential Equations", fourth edition, Prentice Hall Press

References

- [1] <http://www.drumdojo.com/world/persia/kettledrums.htm>
[2] <http://ctr.umkc.edu/user/mgarlitos/timp.html>
[3] <http://www.vsl.co.at/en-us/70/3196/3198/5675.vsl>

Acoustics in Violins

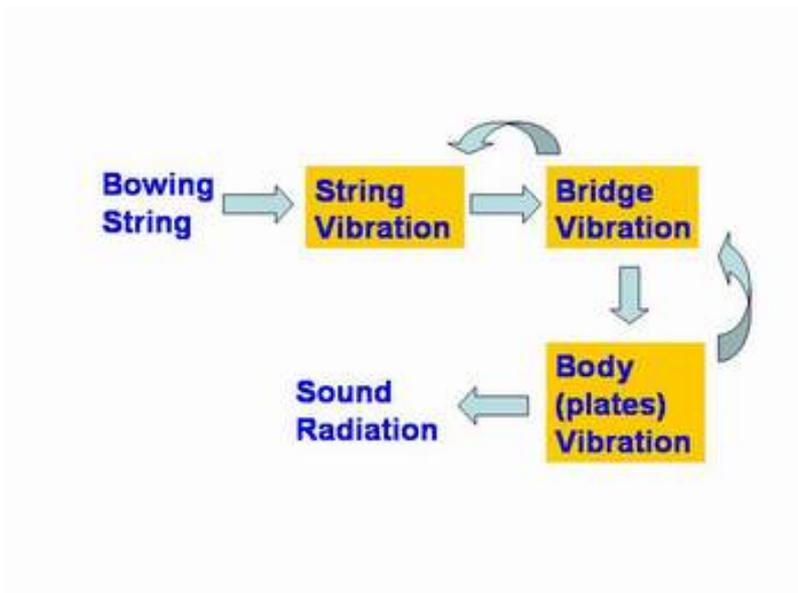
For a detailed anatomy of the violin, please refer to Atelierla Bussiere ^[1].



How does a violin make sound?

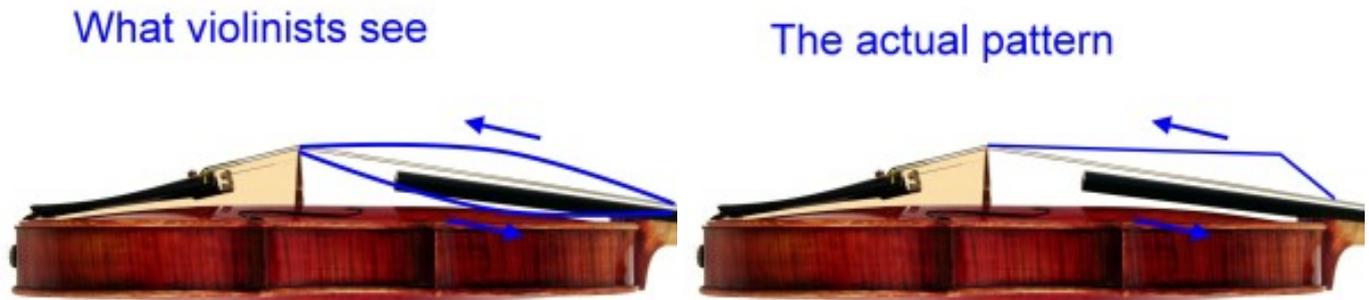
General concept

When a violinist bows a string, which can produce vibrations with abundant harmonics, the vibrations of the strings are structurally transmitted to the bridge and the body of the instrument through the bridge. The bridge transmits the vibrational energy produced by the strings to the body through its feet, further triggering the vibration of body. The vibration of the body determines sound radiation and sound quality, along with the resonance of the cavity.



String

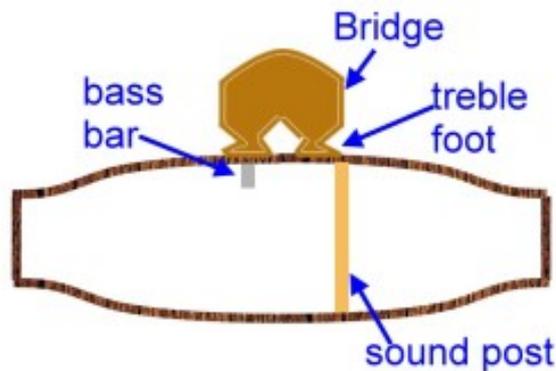
The vibration pattern of the strings can be easily be observed. To the naked eye, the string appears to move back and forth in a parabolic shape (see figure), which resembles the first mode of free vibration of a stretched string. The vibration of strings was first investigated by Hermann Von Helmholtz ^[2], the famous mathematician and physicist in 19th century. A surprising scenario was discovered that the string actually moves in an inverse “V” shape rather than parabolas (see figure). What we see is just an envelope of the motion of the string. To honor his findings, the motion of bowed strings had been called “Helmholtz motion.”



Bridge

The primary role of the bridge is to transform the motion of vibrating strings into periodic driving forces by its feet to the top plate of the violin body. The configuration of the bridge can be referred to the figure. The bridge stands on the belly between f holes, which have two primary functions. One is to connect the air inside the body with outside air, and the other one is to make the belly between f holes move more easily than other parts of the body. The fundamental frequency of a violin bridge was found to be around 3000 Hz when it is on a rigid support, and it is an effective energy-transmitting medium to transmit the energy from the string to body at frequencies from 1 kHz to 4kHz, which is in the range of keen sensitivity of human hearing. In order to darken the sound of violin, the player attaches a mute on the bridge. The mute is actually an additional mass which reduces the fundamental frequency of the bridge. As a result, the sound at higher frequencies is diminished since the force transferred to the body has been decreased. On the other hand, the fundamental frequency of the bridge can be raised by attaching an additional stiffness in the form of tiny wedges, and the sound at higher frequencies will be amplified accordingly.

The sound post connects the flexible belly to the much stiffer back plate. The sound post can prevent the collapse of the belly due to high tension force in the string, and, at the same time, couples the vibration of the plate. The bass bar under the belly extends beyond the f holes and transmits the force of the bridge to a larger area of the belly. As can be seen in the figure, the motion of the treble foot is restricted by the sound post, while, conversely, the foot over bass bar can move up and down more easily. As a result, the bridge tends to move up and down, pivoting about the treble foot. The forces appearing at the two feet remain equal and opposite up to 1 kHz. At higher frequencies, the forces become uneven. The force on the soundpost foot predominates at some frequencies, while it is the bass bar foot at some.



Body

The body includes top plate, back plate, the sides, and the air inside, all of which serve to transmit the vibration of the bridge into the vibration of air surrounding the violin. For this reason, the violin needs a relatively large surface area to push enough amount of air back and forth. Thus, the top and back plates play important roles in the mechanism. Violin makers have traditionally paid much attention to the vibration of the top and back plates of the violin by listening to the tap tones, or, recently, by observing the vibration mode shapes of the body plates. The vibration modes of an assembled violin are, however, much more complicated.

The vibration modes of top and back plates can be easily observed in a similar technique first performed by Ernest Florens Friedrich Chaldni (1756 – 1827), who is often respectfully referred “the father of acoustics.” First, the fine sand is uniformly sprinkled on the plate. Then, the plate can be resonated, either by a powerful sound wave tuned to the desired frequencies, by being bowed by a violin bow, or by being excited mechanically or electromechanically at desired frequencies. Consequently, the sand disperses randomly due to the vibration of plate. Some of the sand falls outside the region of plate, while some of the sand is collected by the nodal regions, which have relatively small movement, of the plate. Hence, the mode shapes of the plate can be visualized in this manner, which can be referred to the figures in the reference site, Violin Acoustics^[3]. The first seven modes of the top and back plates of violin are presented, with nodal lines depicted by using black sands.

The air inside the body is also important, especially in the range of lower frequencies. It is like the air inside a bottle when you blow into the neck, or, as known as Helmholtz resonance, which has its own modes of vibration. The air inside the body can communicate with air outside through the f holes, and the outside air serves as medium carrying waves from the violin.

See www.violinbridges.co.uk^[4] for more articles on bridges and acoustics.

Sound radiation

A complete description of sound radiation of a violin should include the information about radiation intensity as functions both of frequency and location. The sound radiation can be measured by a microphone connected to a pressure level meter which is rotatably supported on a stand arm around the violin, while the violin is fastened at the neck by a clip. The force is introduced into the violin by using a miniature impact hammer at the upper edge of the bridge in the direction of bowing. The detail can be referred to Martin Schleske, master studio for violinmaking^[5]. The radiation intensity of different frequencies at different locations can be represented by directional characteristics, or acoustic maps. The directional characteristics of a violin can be shown in the figure in the website of Martin Schleske^[6], where the radial distance from the center point represents the absolute value of the sound level (re 1Pa/N) in dB, and the angular coordinate of the full circle indicates the measurement point around the instrument. According to the directional characteristics of violins, the principal radiation directions for the violin in the horizontal plane can be established. For more detail about the principal radiation direction for violins at different frequencies, please refer to reference (Meyer 1972).

References and other links

- Violin Acoustics ^[7]
- Paul Galluzzo's Homepage ^[8]
- Martin Schleske, master studio for violinmaking ^[5]
- Atelierla Bussiere ^[9]
- Fletcher, N. H., and Rossing, T. D., *The physics of musical instrument*, Springer-Verlag, 1991
- Meyer, J., "Directivity of bowed stringed instruments and its effect on orchestral sound in concert halls", *J. Acoustic. Soc. Am.*, 51, 1972, pp. 1994-2009

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- [1] <http://www.atelierlabussiere.com/pcom4englis.htm>
- [2] <http://www-groups.dcs.st-and.ac.uk/~history/Mathematicians/Helmholtz.html>
- [3] <http://www.phys.unsw.edu.au/~jw/patterns2.html>
- [4] <http://www.violinbridges.co.uk>
- [5] http://www.schleske.de/index.php?lang=en&http://www.schleske.de/06geigenbauer/en_akustik3schall3messmeth.shtml
- [6] http://www.schleske.de/picshoch/a0044_abstr_netz.gif
- [7] <http://www.phys.unsw.edu.au/music/violin/>
- [8] http://www2.eng.cam.ac.uk/~pmg26/home_frame.html
- [9] <http://www.atelierlabussiere.com/>

Microphone Technique



General technique

1. A microphone should be used whose frequency response will suit the frequency range of the voice or instrument being recorded.
2. Vary microphone positions and distances until you achieve the monitored sound that you desire.
3. In the case of poor room acoustics, place the microphone very close to the loudest part of the instrument being recorded or isolate the instrument.
4. Personal taste is the most important component of microphone technique. Whatever sounds right to you, *is* right.

Types of microphones

Dynamic microphones

These are the most common general-purpose microphones. They do not require power to operate. If you have a microphone that is used for live performance, it is probably a dynamic mic.

They have the advantage that they can withstand very high sound pressure levels (high volume) without damage or distortion, and tend to provide a richer, more intense sound than other types. Traditionally, these mics did not provide as good a response on the highest frequencies (particularly above 10 kHz), but some recent models have come out that attempt to overcome this limitation.

In the studio, dynamic mics are often used for high sound pressure level instruments such as drums, guitar amps and brass instruments. Models that are often used in recording include the Shure SM57 and the Sennheiser MD421.

Condenser microphones

These microphones are often the most expensive microphones a studio owns. They require power to operate, either from a battery or phantom power, provided using the mic cable from an external mixer or pre-amp. These mics have a built-in pre-amplifier that uses the power. Some vintage microphones have a tube amplifier, and are referred to as tube condensers.

While they cannot withstand the very high sound pressure levels that dynamic mics can, they provide a flatter frequency response, and often the best response at the highest frequencies. Not as good at conveying intensity, they are much better at providing a balanced accurate sound.

Condenser mics come with a variety of sizes of transducers. They are usually grouped into smaller format condensers, which often are long cylinders about the size of a nickel coin in diameter, and larger format condensers, the transducers of which are often about an inch in diameter or slightly larger.

In the studio, condenser mics are often used for instruments with a wide frequency range, such as an acoustic piano, acoustic guitar, voice, violin, cymbals, or an entire band or chorus. On louder instruments they do not use close miking with condensers. Models that are often used in recording include the Shure SM81 (small format), AKG C414 (large format) and Neumann U87 (large format).

Ribbon microphones

Ribbon microphones are often used as an alternative to condenser microphones. Some modern ribbon microphones do not require power, and some do. The first ribbon microphones, developed at RCA in the 1930s, required no power, were quite fragile and could be destroyed by just blowing air through them. Modern ribbon mics are much more resilient, and can be used with the same level of caution as condenser mics.

Ribbon microphones provide a warmer sound than a condenser mic, with a less brittle top end. Some vocalists (including Paul McCartney) prefer them condenser mics. In the studio they are used on vocals, violins, and even drums. Popular models for recording include the Royer R121 and the AEA R84.

Working distance

Close miking

When miking at a distance of 1 inch to about 1 foot from the sound source, it is considered close miking. This technique generally provides a tight, present sound quality and does an effective job of isolating the signal and excluding other sounds in the acoustic environment.

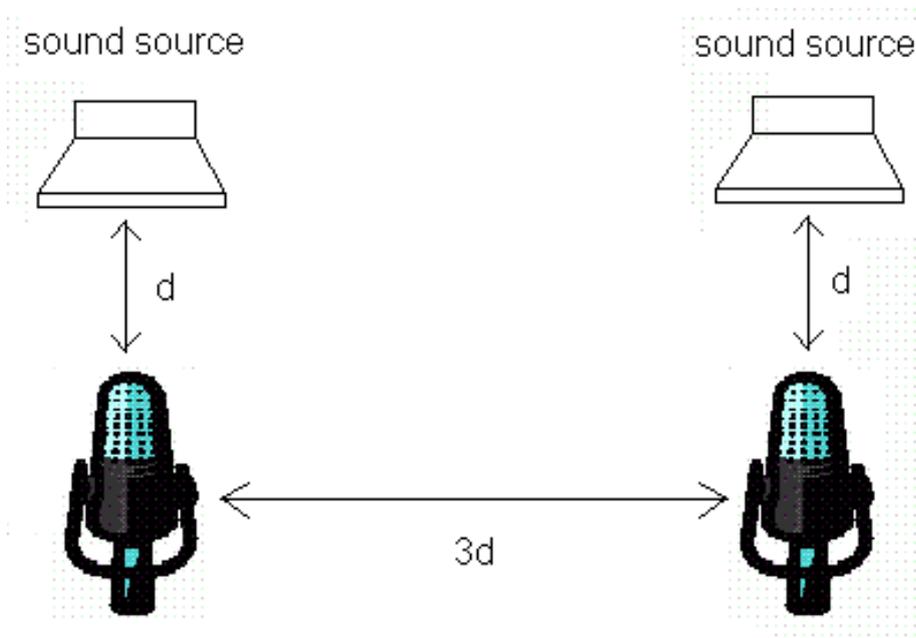
Bleed

Bleeding occurs when the signal is not properly isolated and the microphone picks up another nearby instrument. This can make the mixdown process difficult if there are multiple voices on one track. Use the following methods to prevent leakage:

- Place the microphones closer to the instruments.
- Move the instruments farther apart.
- Put some sort of acoustic barrier between the instruments.
- Use directional microphones.

A B miking

The A B miking distance rule (ratio 3 - 1) is a general rule of thumb for close miking. To prevent phase anomalies and bleed, the microphones should be placed at least three times as far apart as the distance between the instrument and the microphone.



Distant miking

Distant miking refers to the placement of microphones at a distance of 3 feet or more from the sound source. This technique allows the full range and balance of the instrument to develop and it captures the room sound. This tends to add a live, open feeling to the recorded sound, but careful consideration needs to be given to the acoustic environment.

Accent miking

Accent miking is a technique used for solo passages when miking an ensemble. A soloist needs to stand out from an ensemble, but placing a microphone too close will sound unnaturally present compared the distant miking technique used with the rest of the ensemble. Therefore, the microphone should be placed just close enough to the soloist so that the signal can be mixed effectively without sounding completely excluded from the ensemble.

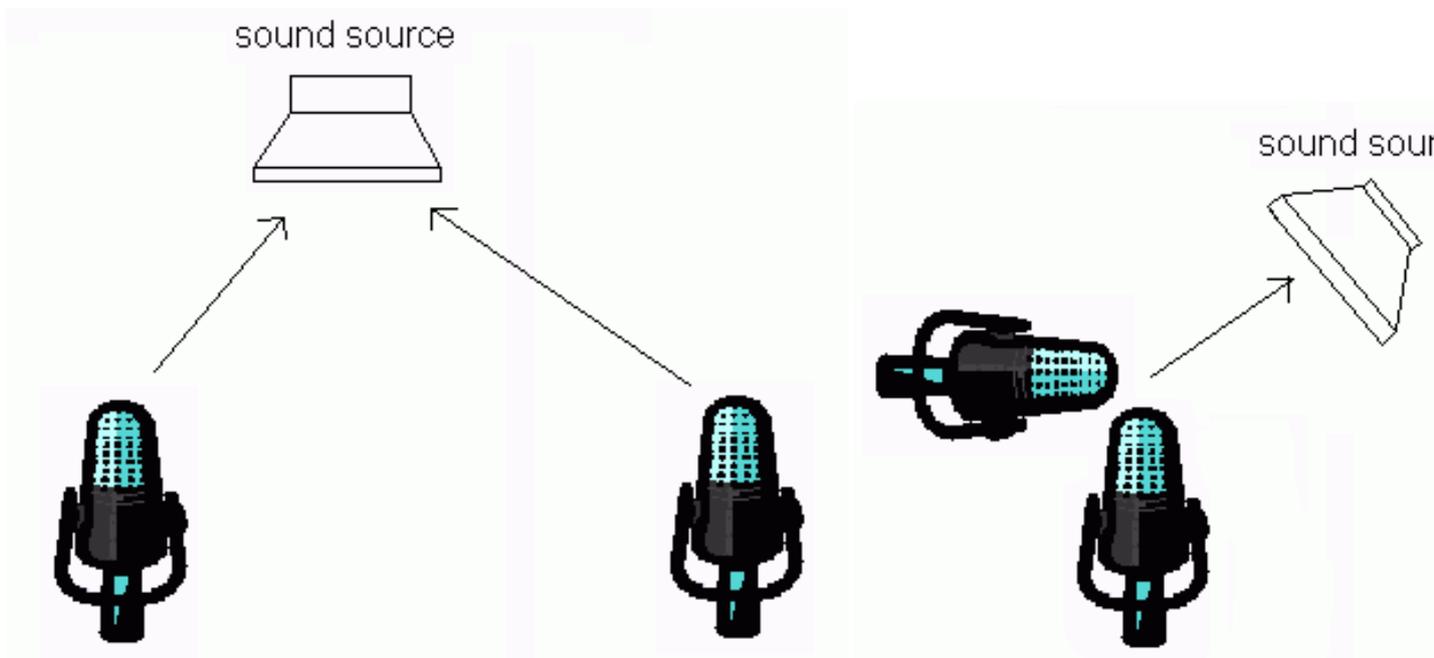
Ambient miking

Ambient miking is placing the microphones at such a distance that the room sound is more prominent than the direct signal. This technique is used to capture audience sound or the natural reverberation of a room or concert hall.

Stereo and surround technique

Stereo

Stereo miking is simply using two microphones to obtain a stereo left-right image of the sound. A simple method is the use of a spaced pair, which is placing two identical microphones several feet apart and using the difference in time and amplitude to create the image. Great care should be taken in the method as phase anomalies can occur due to the signal delay. This risk of phase anomaly can be reduced by using the X/Y method, where the two microphones are placed with the grills as close together as possible without touching. There should be an angle of 90 to 135 degrees between the mics. This technique uses only amplitude, not time, to create the image, so the chance of phase discrepancies is unlikely.



Surround

To take advantage of 5.1 sound or some other surround setup, microphones may be placed to capture the surround sound of a room. This technique essentially stems from stereo technique with the addition of more microphones. Because every acoustic environment is different, it is difficult to define a general rule for surround miking, so placement becomes dependent on experimentation. Careful attention must be paid to the distance between microphones and potential phase anomalies.

Placement for varying instruments

Amplifiers

When miking an amplified speaker, such as for electric guitars, the mic should be placed 2 to 12 inches from the speaker. Exact placement becomes more critical at a distance of less than 4 inches. A brighter sound is achieved when the mic faces directly into the center of the speaker cone and a more mellow sound is produced when placed slightly off-center. Placing off-center also reduces amplifier noise.

A bigger sound can often be achieved by using two mics. The first mic should be a dynamic mic, placed as described in the previous paragraph. Add to this a condenser mic placed at least 3 times further back (remember the 3:1 rule), which will pick up the blended sound of all speakers, as well as some room ambience. Run the mics into separate channels and combine them to your taste.

Brass instruments

High sound-pressure levels are produced by brass instruments due to the directional characteristics of mid to mid-high frequencies. Therefore, for brass instruments such as trumpets, trombones, and tubas, microphones should face slightly off of the bell's center at a distance of one foot or more to prevent overloading from wind blasts.

Guitars

Technique for acoustic guitars is dependent on the desired sound. Placing a microphone close to the sound hole will achieve the highest output possible, but the sound may be bottom-heavy because of how the sound hole resonates at low frequencies. Placing the mic slightly off-center at 6 to 12 inches from the hole will provide a more balanced pickup. Placing the mic closer to the bridge with the same working distance will ensure that the full range of the instrument is captured.

A technique that some engineers use places a large-format condenser mic 12-18 inches away from the 12th fret of the guitar, and a small-format condenser very close to the strings nearby. Combining the two signals can produce a rich tone.

Pianos

Ideally, microphones would be placed 4 to 6 feet from the piano to allow the full range of the instrument to develop before it is captured. This isn't always possible due to room noise, so the next best option is to place the microphone just inside the open lid. This applies to both grand and upright pianos.

Percussion

One overhead microphone can be used for a drum set, although two are preferable. If possible, each component of the drum set should be miked individually at a distance of 1 to 2 inches as if they were their own instrument. This also applies to other drums such as congas and bongos. For large, tuned instruments such as xylophones, multiple mics can be used as long as they are spaced according to the 3:1 rule. Typically, dynamic mics are used for individual drum miking, while small-format condensers are used for the overheads.

Voice

Standard technique is to put the microphone directly in front of the vocalist's mouth, although placing slightly off-center can alleviate harsh consonant sounds (such as "p") and prevent overloading due to excessive dynamic range.

Woodwinds

A general rule for woodwinds is to place the microphone around the middle of the instrument at a distance of 6 inches to 2 feet. The microphone should be tilted slightly towards the bell or sound hole, but not directly in front of it.

Sound Propagation

It is important to understand how sound propagates due to the nature of the acoustic environment so that microphone technique can be adjusted accordingly. There are four basic ways that this occurs:

Reflection

Sound waves are reflected by surfaces if the object is as large as the wavelength of the sound. It is the cause of echo (simple delay), reverberation (many reflections cause the sound to continue after the source has stopped), and standing waves (the distance between two parallel walls is such that the original and reflected waves in phase reinforce one another).

Absorption

Sound waves are absorbed by materials rather than reflected. This can have both positive and negative effects depending on whether you desire to reduce reverberation or retain a live sound.

Diffraction

Objects that may be between sound sources and microphones must be considered due to diffraction. Sound will be stopped by obstacles that are larger than its wavelength. Therefore, higher frequencies will be blocked more easily than lower frequencies.

Refraction

Sound waves bend as they pass through mediums with varying density. Wind or temperature changes can cause sound to seem like it is literally moving in a different direction than expected.

Sources

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- Shure, Inc. (2003). *Shure Product Literature*. Retrieved November 28, 2005, from <http://www.shure.com/scripts/literature/literature.aspx>.

Microphone Design and Operation



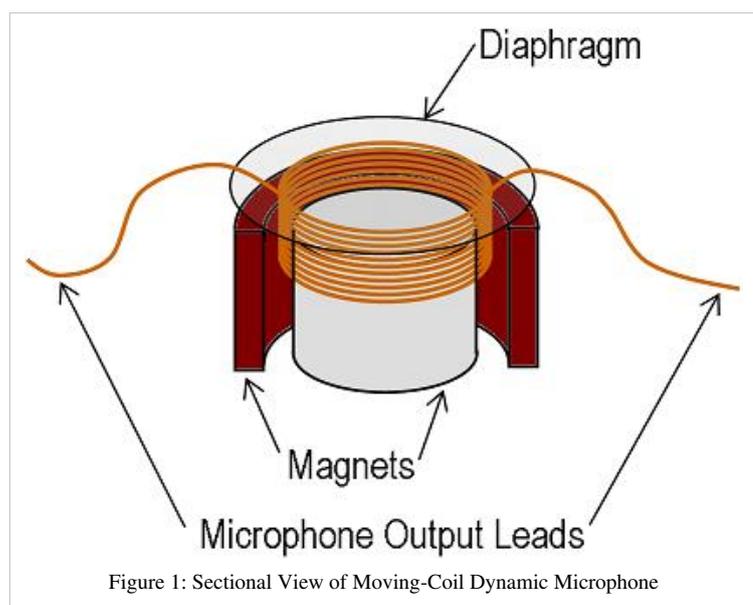
Introduction

Microphones are devices which convert pressure fluctuations into electrical signals. There are two main methods of accomplishing this task that are used in the mainstream entertainment industry. They are known as dynamic microphones and condenser microphones. Piezoelectric crystals can also be used as microphones but are not commonly used in the entertainment industry. For further information on piezoelectric transducers [Click Here](#) ^[1].

Dynamic microphones

This type of microphone converts pressure fluctuations into electrical current. These microphones work by means of the principle known as Faraday's Law. The principle states that when an electrical conductor is moved through a magnetic field, an electrical current is induced within the conductor. The magnetic field within the microphone is created using permanent magnets and the conductor is produced in two common arrangements.

The first conductor arrangement is made of a coil of wire. The wire is typically copper and is attached to a circular membrane or piston usually made from lightweight plastic or occasionally aluminum. The impinging pressure fluctuation on the piston causes it to move in the magnetic field and thus creates the desired electrical current. Figure 1 provides a sectional view of a moving-coil microphone.



The second conductor arrangement is a ribbon of metallic foil suspended between magnets. The metallic ribbon is what moves in response to a pressure fluctuation and in the same manner, an electrical current is produced. Figure 2 provides a sectional view of a ribbon microphone. In both configurations, dynamic microphones follow the same principals as acoustical transducers. For further information about transducers Click Here ^[2].

Condenser microphones

This type of microphone converts pressure fluctuations into electrical potentials through the use of changing an electrical capacitor. This is why condenser microphones are also

known as capacitor microphones. An electrical capacitor is created when two charged electrical conductors are placed at a finite distance from each other. The basic relation that describes capacitors is:

$$Q=C*V$$

where Q is the electrical charge of the capacitor's conductors, C is the capacitance, and V is the electric potential between the capacitor's conductors. If the electrical charge of the conductors is held at a constant value, then the voltage between the conductors will be inversely proportional to the capacitance. Also, the capacitance is inversely proportional to the distance between the conductors. Condenser microphones utilize these two concepts.

The capacitor in a condenser microphone is made of two parts: the diaphragm and the back plate. Figure 3 shows a section view of a condenser microphone. The diaphragm is what moves due to impinging pressure fluctuations and the back plate is held in a stationary position. When the

diaphragm moves closer to the back plate, the capacitance increases and therefore a change in electric potential is produced. The diaphragm is typically made of metallic coated Mylar. The assembly that houses both the back plate and the diaphragm is commonly referred to as a capsule.

To keep the diaphragm and back plate at a constant charge, an electric potential must be presented to the capsule. There are various ways of performing this operation. The first of which is by simply using a battery to supply the needed DC potential to the capsule. A simplified schematic of this technique is displayed in figure 4. The resistor across the leads of the capsule is very high, in the range of 10 mega ohms, to keep the charge on the capsule close to constant.

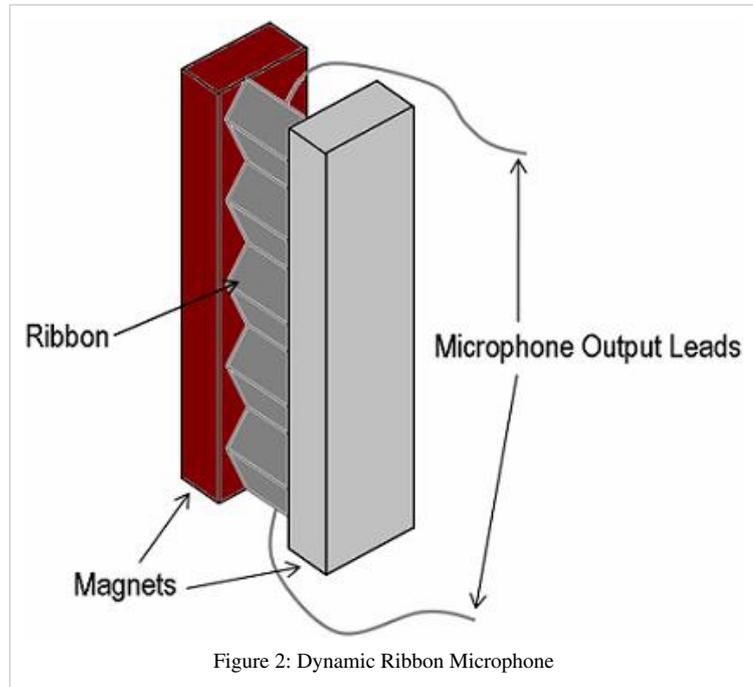


Figure 2: Dynamic Ribbon Microphone

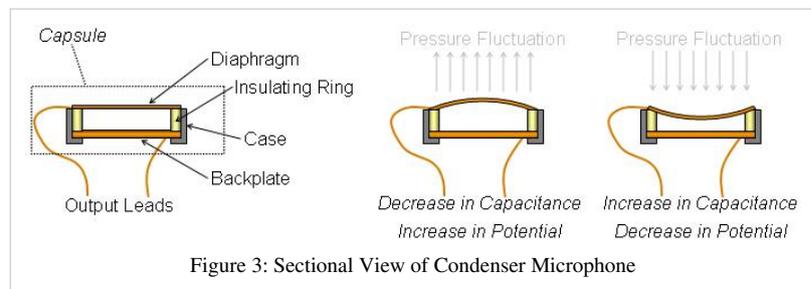


Figure 3: Sectional View of Condenser Microphone

Another technique of providing a constant charge on the capacitor is to supply a DC electric potential through the microphone cable that carries the microphone's output signal. Standard microphone cable is known as XLR cable and is terminated by three pin connectors. Pin one connects to the shield around the cable. The microphone signal is transmitted between pins two and three. Figure 5 displays the layout of dynamic microphone attached to a mixing console via XLR cable.

Phantom Supply/Powering (Audio Engineering Society, DIN 45596): The first and most popular method of providing a DC potential through a microphone cable is to supply +48 V to both of the microphone output leads, pins 2 and 3, and use the shield of the cable, pin 1, as the ground to the circuit. Because pins 2 and 3 see the same potential, any fluctuation of the microphone powering potential will not affect the microphone signal seen by the attached audio equipment. This configuration can be seen in figure 6. The +48 V will be stepped down at the microphone using a transformer and provide the potential to the back plate and diaphragm in a similar fashion as the battery solution. In fact, 9, 12, 24, 48 or 52 V can be supplied, but 48 V is the most frequent.

The second method of running the potential through the cable is to supply 12 V between pins 2 and 3. This method is referred to as **T-powering** (also known as Tonaderspeisung, AB powering; DIN 45595). The main problem with T-powering is that potential fluctuation in the powering of the capsule will be transmitted into an audio signal because the audio equipment analyzing the microphone signal will not see a difference between a potential change across pins 2 and 3 due to a pressure fluctuation and one due to the power source electric potential fluctuation.

Finally, the diaphragm and back plate can be manufactured from a material that maintains a fixed charge. These microphones are termed electrets. In early electret designs, the charge on the material tended to become unstable over time. Recent advances in science and manufacturing have allowed this problem to be eliminated in present designs.

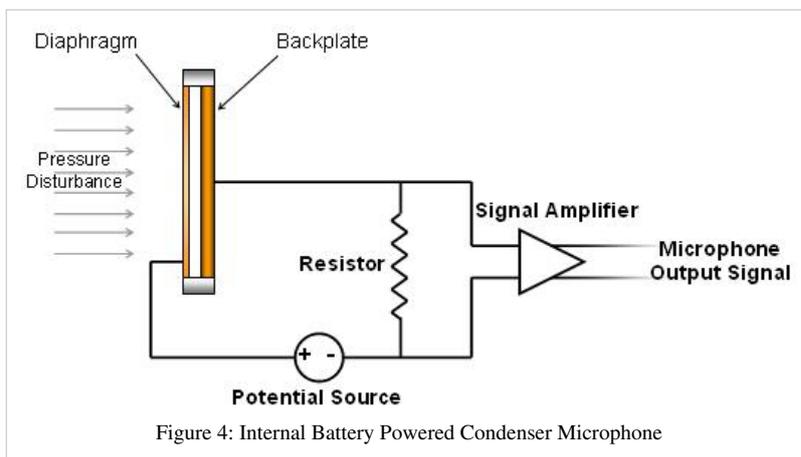


Figure 4: Internal Battery Powered Condenser Microphone

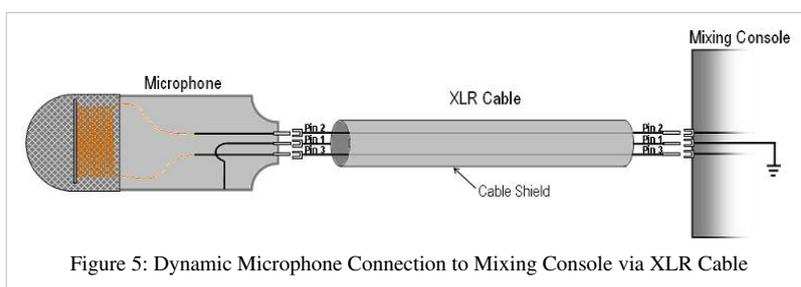


Figure 5: Dynamic Microphone Connection to Mixing Console via XLR Cable

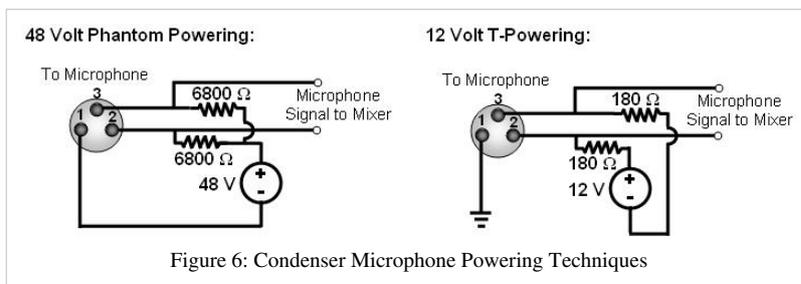


Figure 6: Condenser Microphone Powering Techniques

Conclusion

Two branches of microphones exist in the entertainment industry. Dynamic microphones are found in the moving-coil and ribbon configurations. The movement of the conductor in dynamic microphones induces an electric current which is then transformed into the reproduction of sound. Condenser microphones utilize the properties of capacitors. Creating the charge on the capsule of condenser microphones can be accomplished by battery, phantom powering, T-powering, and by using fixed charge materials in manufacturing.

References

- Sound Recording Handbook. Woram, John M. 1989.
- Handbook of Recording Engineering Fourth Edition. Eargle, John. 2003.

Microphone manufacturer links

- AKG ^[3]
- Audio Technica ^[4]
- Audix ^[5]
- While Bruel & Kjaer ^[6] produces microphones for measurement purposes, DPA ^[7] is the equipment sold for recording purposes
- Electrovoice ^[8]
- Josephson Engineering ^[9]
- Neumann ^[10] (currently a subsidiary of Sennheiser)
- Rode ^[11]
- Schoeps ^[12]
- Sennheiser ^[13]
- Shure ^[14]
- Wharfedale ^[15]

References

- [1] http://en.wikibooks.org/wiki/Acoustic:Piezoelectric_Transducers
- [2] http://en.wikibooks.org/wiki/Acoustic:Acoustic_Transducers_-_The_Loudspeaker
- [3] <http://www.akgusa.com/>
- [4] <http://www.audio-technica.com/cms/site/c35da94027e94819/index.html>
- [5] <http://www.audixusa.com/>
- [6] http://www.bkhome.com/bk_home.asp
- [7] <http://www.dpamicrophones.com/>
- [8] <http://www.electrovoice.com/>
- [9] <http://www.josephson.com/>
- [10] http://www.neumannusa.com/mat_dev/FLift/open.asp
- [11] <http://www.rode.com.au/>
- [12] <http://www.schoeps.de/>
- [13] <http://www.sennheiser.com/>
- [14] <http://www.shure.com/>
- [15] <http://www.wharfedalepro.com/>

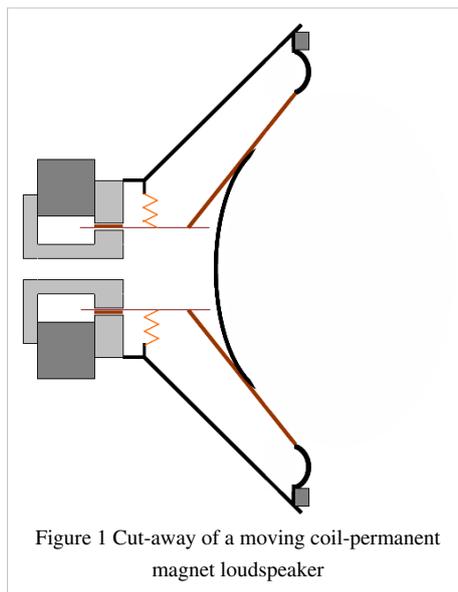
Acoustic Loudspeaker



The purpose of the acoustic transducer is to convert electrical energy into acoustic energy. Many variations of acoustic transducers exist, although the most common is the moving coil-permanent magnet transducer. The classic loudspeaker is of the moving coil-permanent magnet type.

The classic electrodynamic loudspeaker driver can be divided into three key components:

1. The Magnet Motor Drive System
2. The Loudspeaker Cone System
3. The Loudspeaker Suspension



The Magnet Motor Drive System

The main purpose of the Magnet Motor Drive System is to establish a symmetrical magnetic field in which the voice coil will operate. The Magnet Motor Drive System is comprised of a front focusing plate, permanent magnet, back plate, and a pole piece. In figure 2, the assembled drive system is illustrated. In most cases, the back plate and the pole piece are built into one piece called the yoke. The yoke and the front focusing plate are normally made of a very soft cast iron. Iron is a material that is used in conjunction with magnetic structures because the iron is easily saturated when exposed to a magnetic field. Notice in figure 2, that an air gap was intentionally left between the front focusing plate and the yoke. The magnetic field is coupled through the air gap. The magnetic field strength (B) of the air gap is typically optimized for uniformity across the gap. [1]

Figure 2 Permanent Magnet Structure

When a coil of wire with a current flowing is placed inside the permanent magnetic field, a force is produced. B is the magnetic field strength, l is the length of the coil, and I is the current flowing through the coil. The electro-magnetic force is given by the expression of Laplace :

$$d\mathbf{F} = I d\mathbf{l} \times \mathbf{B}$$

\mathbf{B} and $d\mathbf{l}$ are orthogonal, so the force is obtained by integration on the length of the wire (R_e is the radius of a spire, n is the number of spires and \mathbf{e}_x is on the axis of the coil):

$$\mathbf{F} = 2\pi R_e B n I \mathbf{e}_x$$

This force is directly proportional to the current flowing through the coil.

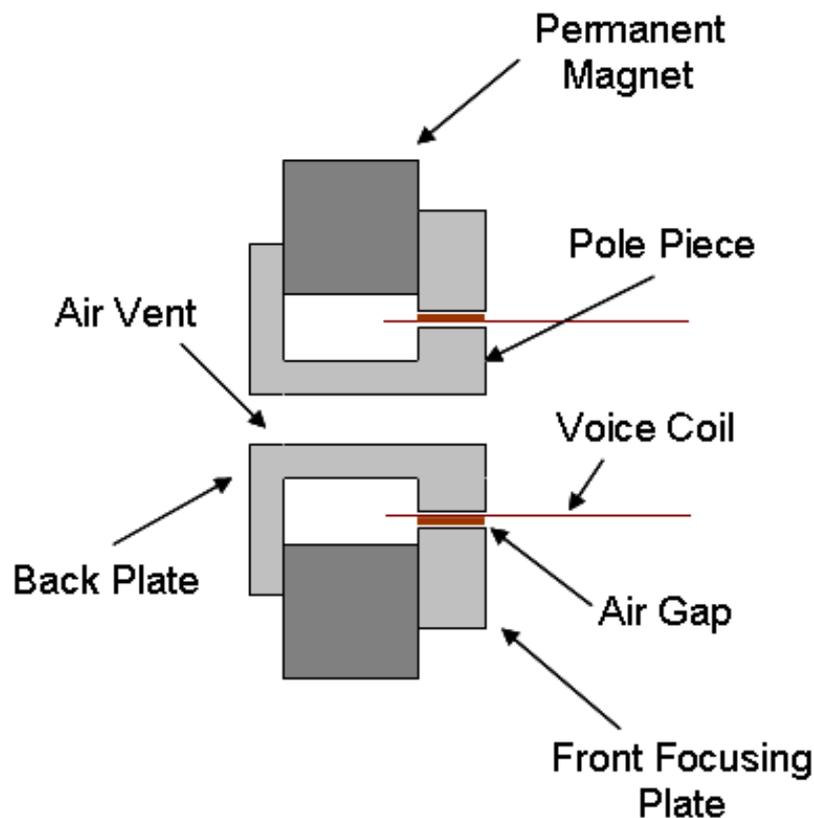


Figure 3 Voice Coil Mounted in Permanent Magnetic Structure

The coil is excited with the AC signal that is intended for sound reproduction, when the changing magnetic field of the coil interacts with the permanent magnetic field then the coil moves back and forth in order to reproduce the input signal. The coil of a loudspeaker is known as the voice coil.

The loudspeaker cone system

On a typical loudspeaker, the cone serves the purpose of creating a larger radiating area allowing more air to be moved when excited by the voice coil. The cone serves a piston that is excited by the voice coil. The cone then displaces air creating a sound wave. In an ideal environment, the cone should be infinitely rigid and have zero mass, but in reality neither is true. Cone materials vary from carbon fiber, paper, bamboo, and just about any other material that can be shaped into a stiff conical shape. The loudspeaker cone is a very critical part of the loudspeaker. Since the cone is not infinitely rigid, it tends to have different types of resonance modes form at different frequencies, which in turn alters and colors the reproduction of the sound waves. The shape of the cone directly influences the directivity and frequency response of the loudspeaker. When the cone is attached to the voice coil, a large gap above the voice coil is left exposed. This could be a problem if foreign particles make their way into the air gap of the voice coil and the permanent magnet structure. The solution to this problem is to place what is known as a dust cap on the cone to cover the air gap. Below a figure of the cone and dust cap are shown.

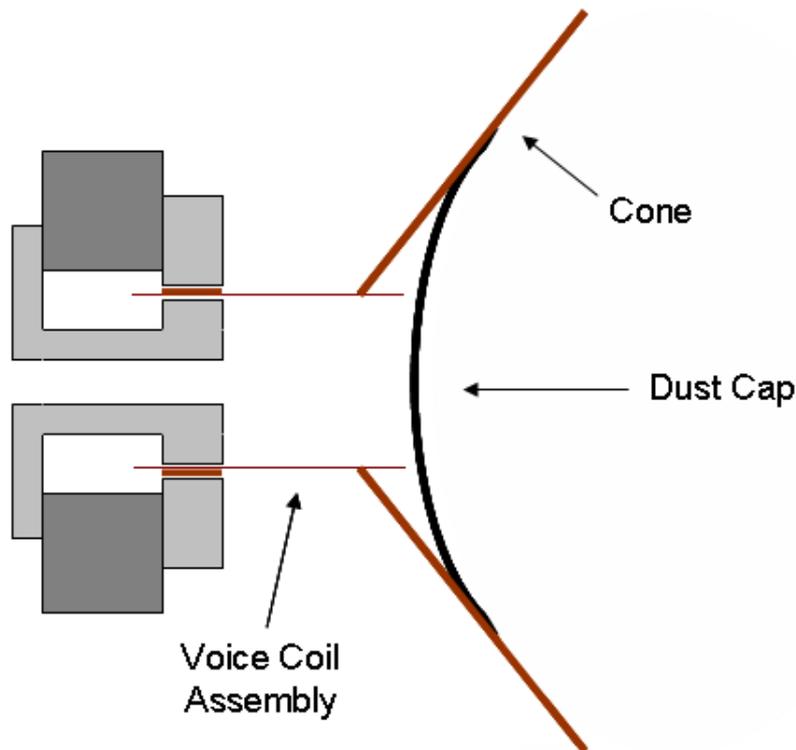


Figure 6 Cone and Dust Cap attached to Voice Coil

The speed of the cone can be expressed with an equation of a mass-spring system with a damping coefficient ξ :

$$m \frac{dv}{dt} + \xi v + k \int v dt = Bli$$

The current intensity i and the speed v can also be related by this equation (U is the voltage, R the electrical resistance and L_b the inductance) :

$$L_b \frac{di}{dt} + Ri = U - Blv$$

By using a harmonic solution, the expression of the speed is :

$$v = \frac{Bli}{\xi + j(m\omega - \frac{k}{\omega})}$$

The electrical impedance can be determined as the ratio of the voltage on the current intensity :

$$Z = \frac{U}{i} = R + jL\omega + \frac{B^2l^2}{\xi + j(m\omega - \frac{k}{\omega})}$$

The frequency response of the loudspeaker can be plotted.

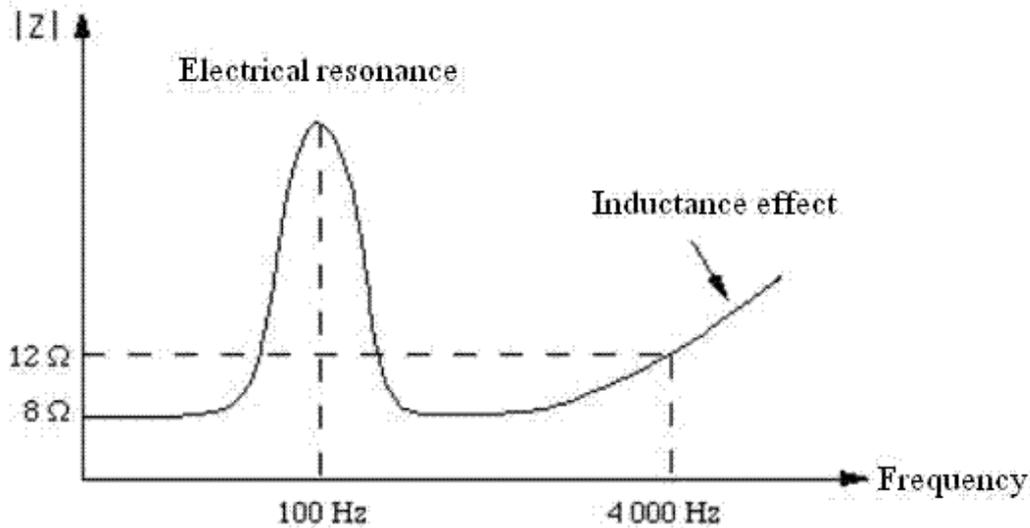


Figure 7 Electrical impedance

A phenomena of electrical resonance is observable around the frequency of 100 Hz. Besides, the inductance of the coil makes the impedance increase from the frequency of 400 Hz. So the range of frequency where the loudspeaker is used is 100 - 4000 Hz

The loudspeaker suspension

Most moving coil loudspeakers have a two piece suspension system, also known as a flexure system. The combination of the two flexures allows the voice coil to maintain linear travel as the voice coil is energized and provide a restoring force for the voice coil system. The two piece system consists of large flexible membrane surrounding the outside edge of the cone, called the surround, and an additional flexure connected directly to the voice coil, called the spider. The surround has another purpose and that is to seal the loudspeaker when mounted in an enclosure. Commonly, the surround is made of a variety of different materials, such as, folded paper, cloth, rubber, and foam. Construction of the spider consists of different woven cloth or synthetic materials that are compressed to form a flexible membrane. The following two figures illustrate where the suspension components are physically at on the loudspeaker and how they function as the loudspeaker operates.

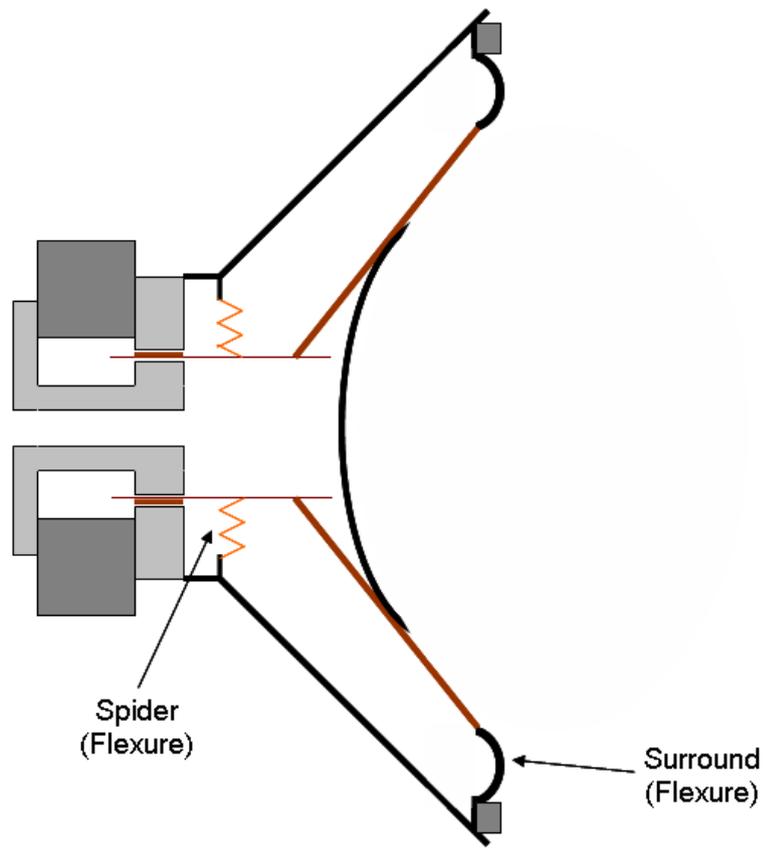


Figure 8 Loudspeaker Suspension System

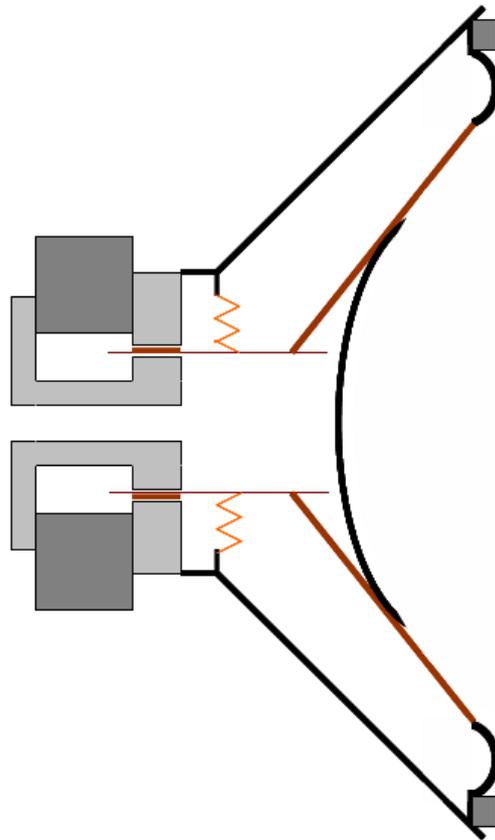


Figure 9 Moving Loudspeaker

Modeling the loudspeaker as a lumped system

Before implementing a loudspeaker into a specific application, a series of parameters characterizing the loudspeaker must be extracted. The equivalent circuit of the loudspeaker is key when developing enclosures. The circuit models all aspects of the loudspeaker through an equivalent electrical, mechanical, and acoustical circuit. Figure 9 shows how the three equivalent circuits are connected. The electrical circuit is comprised of the DC resistance of the voice coil, R_e , the imaginary part of the voice coil inductance, L_e , and the real part of the voice coil inductance, $R_{e_{vc}}$. The mechanical system has electrical components that model different physical parameters of the loudspeaker. In the mechanical circuit, M_m is the electrical capacitance due to the moving mass, C_m is the electrical inductance due to the compliance of the moving mass, and R_m is the electrical resistance due to the suspension system. In the acoustical equivalent circuit, M_A models the air mass and R_A models the radiation impedance[2]. This equivalent circuit allows insight into what parameters change the characteristics of the loudspeaker. Figure 10 shows the electrical input impedance as a function of frequency developed using the equivalent circuit of the loudspeaker.

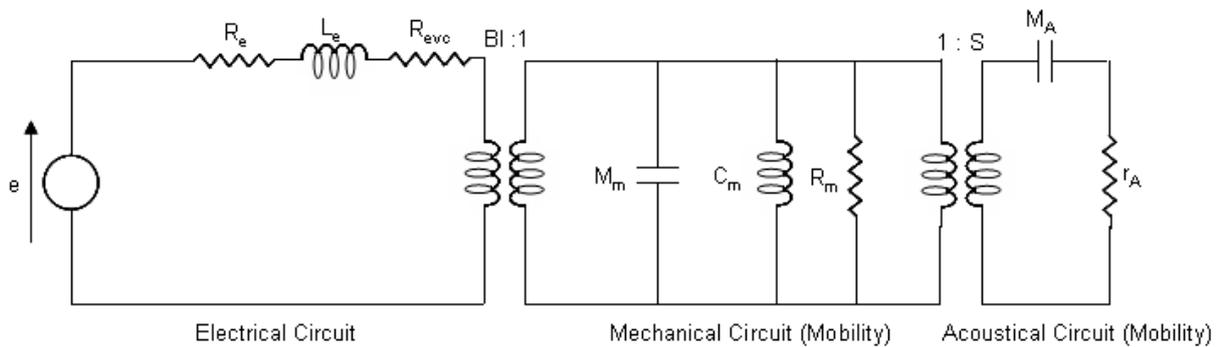


Figure 9 Loudspeaker Analogous Circuit

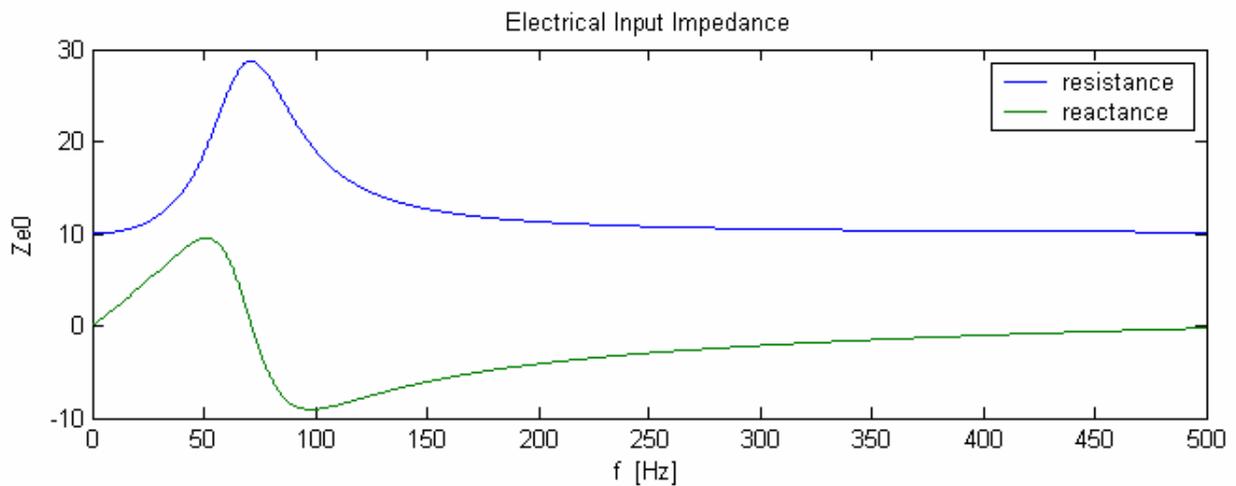


Figure 10 Electrical Input Impedance

The acoustical enclosure

Function of the enclosure

The loudspeaker emits two waves : a front wave and a back wave. With a reflection on a wall, the back wave can be added with the front wave and produces destructive interferences. As a result, the sound pressure level in the room is not uniform. At certain positions, the interaction is additive, and the sound pressure level is higher. On the contrary, certain positions offer destructive interaction between the waves and the sound pressure level is lower.

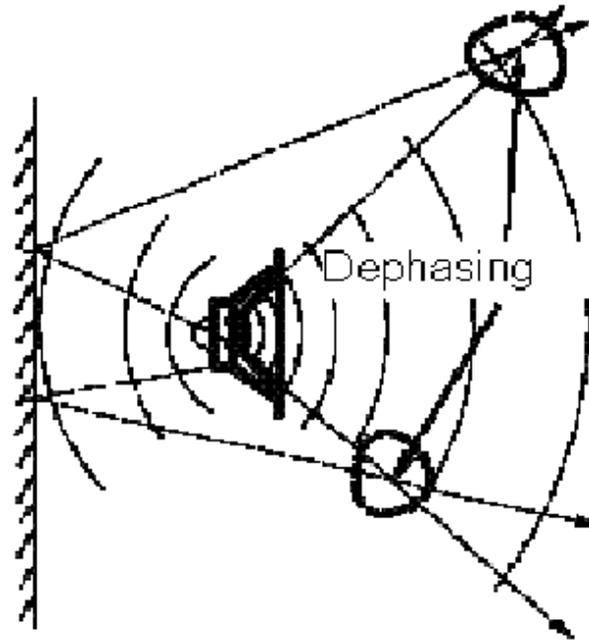


Figure 11 Loudspeaker without baffle producing destructive interferences

The solution is to put a baffle round the loudspeaker in order to prevent the back wave from interfering with the front wave. The sound pressure level is uniform in the room and the quality of the loudspeaker is higher.

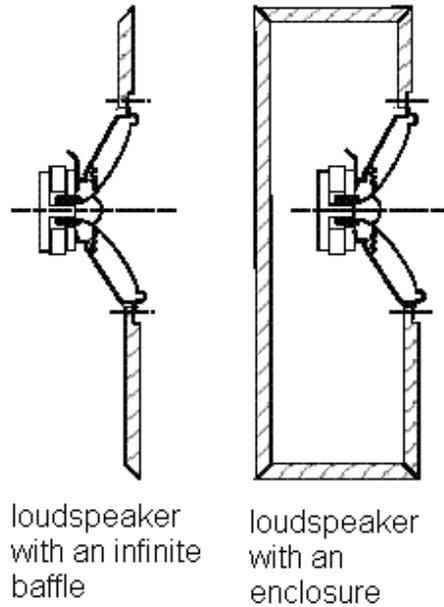


Figure 12 Loudspeakers with infinite baffle and enclosure

Loudspeaker-external fluid interaction

The external fluid exerts a pressure on the membrane of the loudspeaker cone. This additive force can be evaluate as an additive mass and an additive damping in the equation of vibration of the membrane.

$$-(m + MS)\omega^2\xi + i\omega(C + RS)\xi + K\xi = fe^{i\omega t}$$

When the fluid is the air, this additive mass and additive damping are negligible. For example, at the frequency of 1000 Hz, the additive mass is 3g.

Loudspeaker-internal fluid interaction

The volume of air in the enclosure constitutes an additive stiffness. This is called the acoustic load. In low frequencies, this additive stiffness can be four times the stiffness of the loudspeaker cone. The internal air stiffness is very high because of the boundary conditions inside the enclosure. The walls impose a condition of zero airspeed that makes the stiffness increase.

Figure 13 Stiffness of the loudspeaker cone and stiffness of the internal air

The stiffness of the internal air (in red) is fourth time higher than the stiffness of the loudspeaker cone (in blue). That is why the design of the enclosure is relevant in order to improve the quality of the sound and avoid a decrease of the sound pressure level in the room at some frequencies.

References

1. The Loudspeaker Design Cookbook 5th Edition; Dickason, Vance., Audio Amateur Press, 1997.
2. Beranek, L. L. Acoustics. 2nd ed. Acoustical Society of America, Woodbridge, NY. 1993.

Sealed Box Subwoofer Design

Introduction

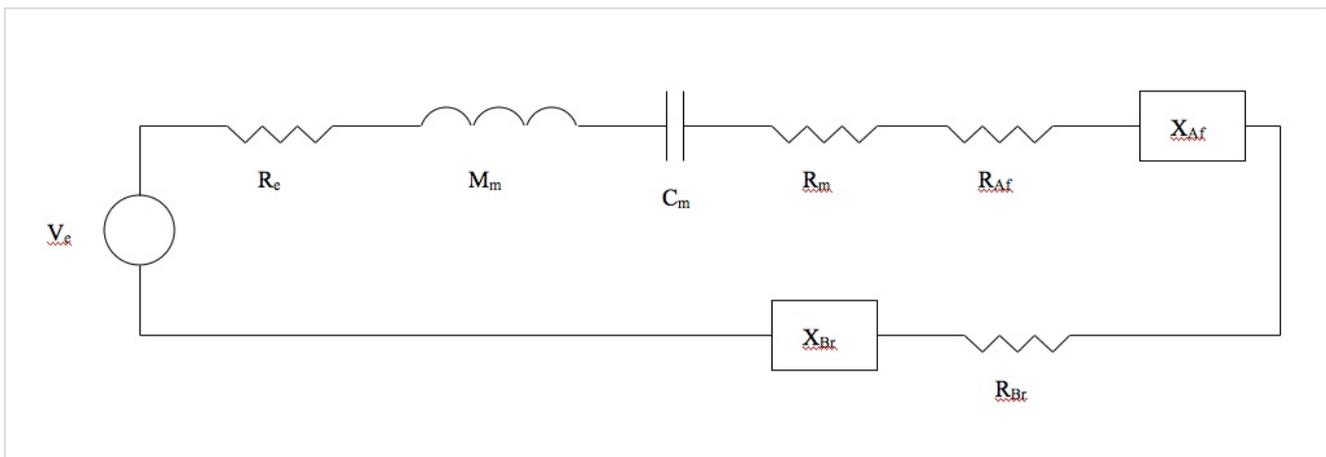
A sealed or closed box baffle is the most basic but often the cleanest sounding sub-woofer box design. The sub-woofer box in its most simple form, serves to isolate the back of the speaker from the front, much like the theoretical infinite baffle. The sealed box provides simple construction and controlled response for most sub-woofer applications. The slow low end roll-off provides a clean transition into the extreme frequency range. Unlike ported boxes, the cone excursion is reduced below the resonant frequency of the box and driver due to the added stiffness provided by the sealed box baffle.

Closed baffle boxes are typically constructed of a very rigid material such as MDF (medium density fiber board) or plywood .75 to 1 inch thick. Depending on the size of the box and material used, internal bracing may be necessary to maintain a rigid box. A rigid box is important to design in order to prevent unwanted box resonance.

As with any acoustics application, the box must be matched to the loudspeaker driver for maximum performance. The following will outline the procedure to tune the box or maximize the output of the sub-woofer box and driver combination.

Closed baffle circuit

The sealed box enclosure for sub-woofers can be modeled as a lumped element system if the dimensions of the box are significantly shorter than the shortest wavelength reproduced by the sub-woofer. Most sub-woofer applications are crossed over around 80 to 100 Hz. A 100 Hz wave in air has a wavelength of about 11 feet. Sub-woofers typically have all dimensions much shorter than this wavelength, thus the lumped element system analysis is accurate. Using this analysis, the following circuit represents a sub-woofer enclosure system.



where all of the following parameters are in the mechanical mobility analog

- V_e - voltage supply
- R_e - electrical resistance
- M_m - driver mass
- C_m - driver compliance

R_m - resistance

R_{Ar} - rear cone radiation resistance into the air

X_{Af} - front cone radiation reactance into the air

R_{Br} - rear cone radiation resistance into the box

X_{Br} - rear cone radiation reactance into the box

Driver parameters

In order to tune a sealed box to a driver, the driver parameters must be known. Some of the parameters are provided by the manufacturer, some are found experimentally, and some are found from general tables. For ease of calculations, all parameters will be represented in the SI units meter/kilogram/second. The parameters that must be known to determine the size of the box are as follows:

f_0 - driver free-air resonance

C_{MS} - mechanical compliance of the driver

S_D - effective area of the driver

Resonance of the driver

The resonance of the driver is usually either provided by the manufacturer or must be found experimentally. It is a good idea to measure the resonance frequency even if it is provided by the manufacturer to account for inconsistent manufacturing processes.

The following diagram shows verga and the setup for finding resonance:

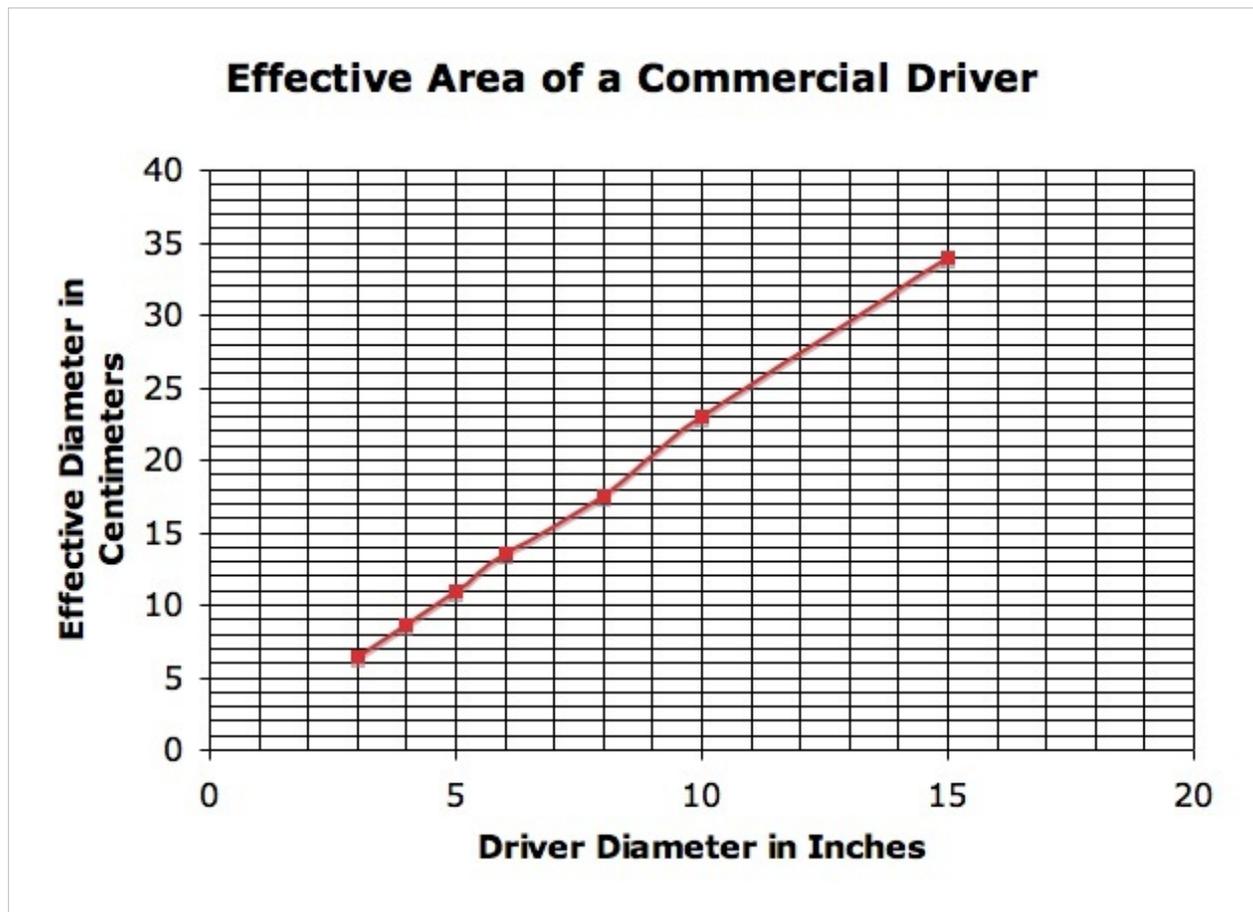
Where voltage V1 is held constant and the variable frequency source is varied until V2 is a maximum. The frequency where V2 is a maximum is the resonance frequency for the driver.

Mechanical compliance

By definition compliance is the inverse of stiffness or what is commonly referred to as the spring constant. The compliance of a driver can be found by measuring the displacement of the cone when known masses are placed on the cone when the driver is facing up. The compliance would then be the displacement of the cone in meters divided by the added weight in newtons.

Effective area of the driver

The physical diameter of the driver does not lead to the effective area of the driver. The effective diameter can be found using the following diagram:



From this diameter, the area is found from the basic area of a circle equation.

Acoustic compliance

From the known mechanical compliance of the cone, the acoustic compliance can be found from the following equation:

$$C_{AS} = C_{MS} S_D^2$$

From the driver acoustic compliance, the box acoustic compliance is found. This is where the final application of the sub-woofer is considered. The acoustic compliance of the box will determine the percent shift upwards of the resonant frequency. If a large shift is desired for high SPL applications, then a large ratio of driver to box acoustic compliance would be required. If a more flattened response is desired for high fidelity applications, then a lower ratio of driver to box acoustic compliance would be required. Specifically, the ratios can be found in the following figure using line (b) as reference.

$$C_{AS} = C_{AB} * r$$

r - driver to box acoustic compliant ratio



Sealed box design

Volume of box

The volume of the sealed box can now be found from the box acoustic compliance. The following equation is used to calculate the box volume

$$V_B = C_{AB} \cdot \gamma$$

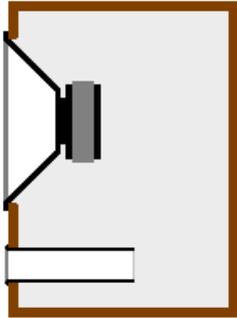
Box dimensions

From the calculated box volume, the dimensions of the box can then be designed. There is no set formula for finding the dimensions of the box, but there are general guidelines to be followed. If the driver was mounted in the center of a square face, the waves generated by the cone would reach the edges of the box at the same time, thus when combined would create a strong diffracted wave in the listening space. In order to best prevent this, the driver should be either be mounted offset of a square face, or the face should be rectangular.

The face of the box which the driver is set in should not be a square.

Miscellaneous Applications

Bass-Reflex Enclosure Design



Bass-reflex enclosures improve the low-frequency response of loudspeaker systems. Bass-reflex enclosures are also called "vented-box design" or "ported-cabinet design". A bass-reflex enclosure includes a vent or port between the cabinet and the ambient environment. This type of design, as one may observe by looking at contemporary loudspeaker products, is still widely used today. Although the construction of bass-reflex enclosures is fairly simple, their design is not simple, and requires proper tuning. This reference focuses on the technical details of bass-reflex design. General loudspeaker information can be found here ^[1].

Effects of the Port on the Enclosure Response

Before discussing the bass-reflex enclosure, it is important to be familiar with the simpler sealed enclosure system performance. As the name suggests, the sealed enclosure system attaches the loudspeaker to a sealed enclosure (except for a small air leak included to equalize the ambient pressure inside). Ideally, the enclosure would act as an acoustical compliance element, as the air inside the enclosure is compressed and rarified. Often, however, an acoustic material is added inside the box to reduce standing waves, dissipate heat, and other reasons. This adds a resistive element to the acoustical lumped-element model. A non-ideal model of the effect of the enclosure actually adds an acoustical mass element to complete a series lumped-element circuit given in Figure 1. For more on sealed enclosure design, see the Sealed Box Subwoofer Design page.

Figure 1. Sealed enclosure acoustic circuit.

In the case of a bass-reflex enclosure, a port is added to the construction. Typically, the port is cylindrical and is flanged on the end pointing outside the enclosure. In a bass-reflex enclosure, the amount of acoustic material used is usually much less than in the sealed enclosure case, often none at all. This allows air to flow freely through the port. Instead, the larger losses come from the air leakage in the enclosure. With this setup, a lumped-element acoustical circuit has the form shown in the diagram below.

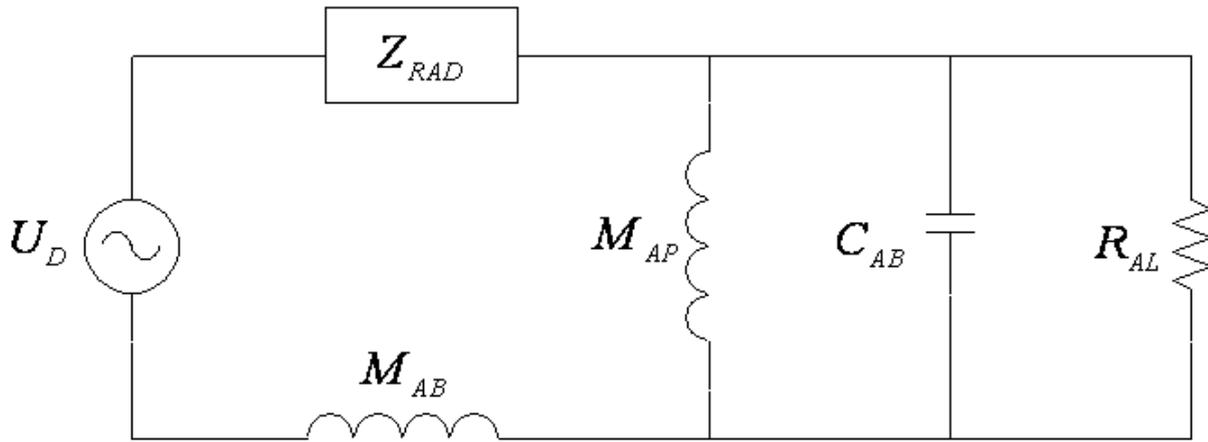


Figure 2. Bass-reflex enclosure acoustic circuit.

In this figure, Z_{RAD} represents the radiation impedance of the outside environment on the loudspeaker diaphragm. The loading on the rear of the diaphragm has changed when compared to the sealed enclosure case. If one visualizes the movement of air within the enclosure, some of the air is compressed and rarified by the compliance of the enclosure, some leaks out of the enclosure, and some flows out of the port. This explains the parallel combination of M_{AP} , C_{AB} , and R_{AL} . A truly realistic model would incorporate a radiation impedance of the port in series with M_{AP} , but for now it is ignored. Finally, M_{AB} , the acoustical mass of the enclosure, is included as discussed in the sealed enclosure case. The formulas which calculate the enclosure parameters are listed in Appendix B. It is important to note the parallel combination of M_{AP} and C_{AB} . This forms a Helmholtz resonator (click here for more information). Physically, the port functions as the “neck” of the resonator and the enclosure functions as the “cavity.” In this case, the resonator is driven from the piston directly on the cavity instead of the typical Helmholtz case where it is driven at the “neck.” However, the same resonant behavior still occurs at the enclosure resonance frequency, f_B . At this frequency, the impedance seen by the loudspeaker diaphragm is large (see Figure 3 below). Thus, the load on the loudspeaker reduces the velocity flowing through its mechanical parameters, causing an anti-resonance condition where the displacement of the diaphragm is a minimum. Instead, the majority of the volume velocity is actually emitted by the port itself instead of the loudspeaker. When this impedance is reflected to the electrical circuit, it is proportional to $1/Z$, thus a minimum in the impedance seen by the voice coil is small. Figure 3 shows a plot of the impedance seen at the terminals of the loudspeaker. In this example, f_B was found to be about 40 Hz, which corresponds to the null in the voice-coil impedance.

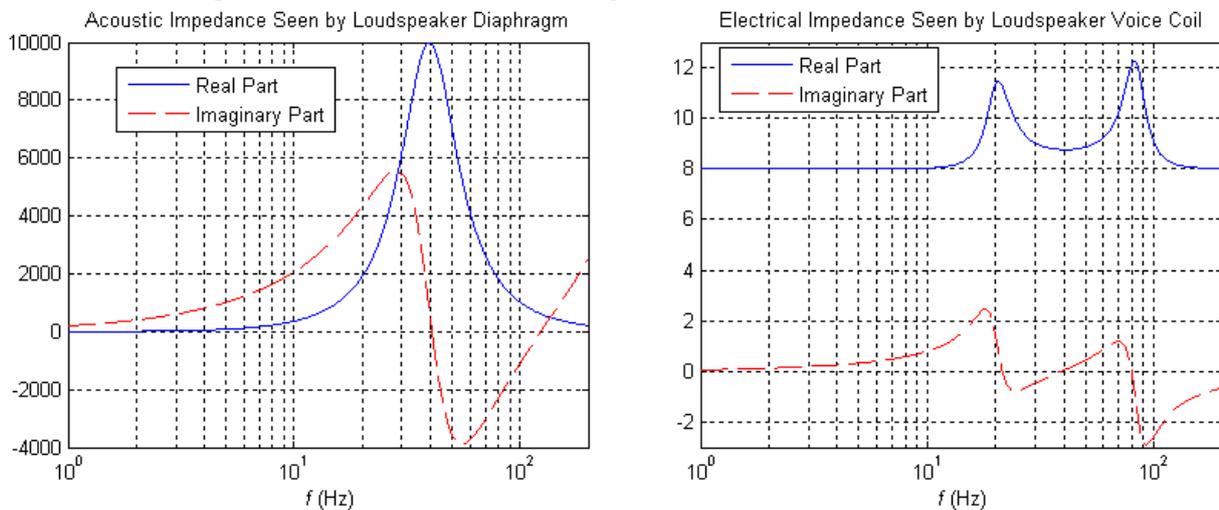


Figure 3. Impedances seen by the loudspeaker diaphragm and voice coil.

Quantitative Analysis of Port on Enclosure

The performance of the loudspeaker is first measured by its velocity response, which can be found directly from the equivalent circuit of the system. As the goal of most loudspeaker designs is to improve the bass response (leaving high-frequency production to a tweeter), low frequency approximations will be made as much as possible to simplify the analysis. First, the inductance of the voice coil, L_E , can be ignored as long as $\omega \ll R_E/L_E$. In a typical loudspeaker, L_E is of the order of 1 mH, while R_E is typically 8 Ω , thus an upper frequency limit is approximately 1 kHz for this approximation, which is certainly high enough for the frequency range of interest. Another approximation involves the radiation impedance, Z_{RAD} . It can be shown [1] that this value is given by the following equation (in acoustical ohms):

$$Z_{RAD} = \frac{\rho_0 c}{\pi a^2} \left[\left(1 - \frac{J_1(2ka)}{ka} \right) + j \frac{H_1(2ka)}{ka} \right]$$

Where $J_1(x)$ and $H_1(x)$ are types of Bessel functions. For small values of ka ,

$$J_1(2ka) \approx ka \quad \text{and} \quad H_1(2ka) \approx \frac{8(ka)^2}{3\pi} \Rightarrow Z_{RAD} \approx j \frac{8\rho_0\omega}{3\pi^2 a} = jM_{A1}$$

Hence, the low-frequency impedance on the loudspeaker is represented with an acoustic mass M_{A1} [1]. For a simple analysis, R_E , M_{MD} , C_{MS} , and R_{MS} (the transducer parameters, or *Thiele-Small* parameters) are converted to their acoustical equivalents. All conversions for all parameters are given in Appendix A. Then, the series masses, M_{AD} , M_{A1} , and M_{AB} , are lumped together to create M_{AC} . This new circuit is shown below.

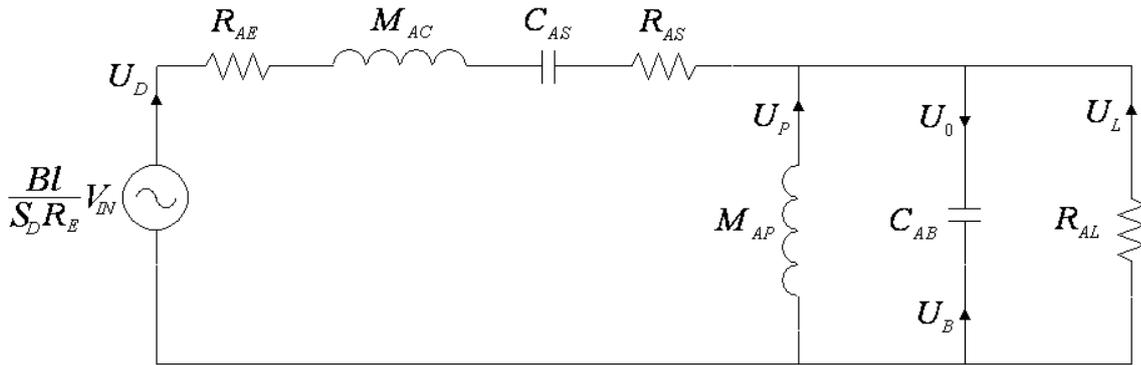


Figure 4. Low-Frequency Equivalent Acoustic Circuit

Unlike sealed enclosure analysis, there are multiple sources of volume velocity that radiate to the outside environment. Hence, the diaphragm volume velocity, U_D , is not analyzed but rather $U_0 = U_D + U_P + U_L$. This essentially draws a “bubble” around the enclosure and treats the system as a source with volume velocity U_0 . This “lumped” approach will only be valid for low frequencies, but previous approximations have already limited the analysis to such frequencies anyway. It can be seen from the circuit that the volume velocity flowing *into* the enclosure, $U_B = -U_0$, compresses the air inside the enclosure. Thus, the circuit model of Figure 3 is valid and the relationship relating input voltage, V_{IN} to U_0 may be computed.

In order to make the equations easier to understand, several parameters are combined to form other parameter names. First, ω_B and ω_S , the enclosure and loudspeaker resonance frequencies, respectively, are:

$$\omega_B = \frac{1}{\sqrt{M_{AP} C_{AB}}} \quad \omega_S = \frac{1}{\sqrt{M_{AC} C_{AS}}}$$

Based on the nature of the derivation, it is convenient to define the parameters ω_0 and h , the Helmholtz tuning ratio:

$$\omega_0 = \sqrt{\omega_B \omega_S} \quad h = \frac{\omega_B}{\omega_S}$$

A parameter known as the *compliance ratio* or *volume ratio*, α , is given by:

$$\alpha = \frac{C_{AS}}{C_{AB}} = \frac{V_{AS}}{V_{AB}}$$

Other parameters are combined to form what are known as *quality factors*:

$$Q_L = R_{AL} \sqrt{\frac{C_{AB}}{M_{AP}}} \quad Q_{TS} = \frac{1}{R_{AE} + R_{AS}} \sqrt{\frac{M_{AC}}{C_{AS}}}$$

This notation allows for a simpler expression for the resulting transfer function [1]:

$$\frac{U_0}{V_{IN}} = G(s) = \frac{(s^3/\omega_0^4)}{(s/\omega_0)^4 + a_3(s/\omega_0)^3 + a_2(s/\omega_0)^2 + a_1(s/\omega_0) + 1}$$

where

$$a_1 = \frac{1}{Q_L \sqrt{h}} + \frac{\sqrt{h}}{Q_{TS}} \quad a_2 = \frac{\alpha + 1}{h} + h + \frac{1}{Q_L Q_{TS}} \quad a_3 = \frac{1}{Q_{TS} \sqrt{h}} + \frac{\sqrt{h}}{Q_L}$$

Development of Low-Frequency Pressure Response

It can be shown [2] that for $ka < 1/2$, a loudspeaker behaves as a spherical source. Here, a represents the radius of the loudspeaker. For a 15" diameter loudspeaker in air, this low frequency limit is about 150 Hz. For smaller loudspeakers, this limit increases. This limit dominates the limit which ignores L_E , and is consistent with the limit that models Z_{RAD} by M_{A1} .

Within this limit, the loudspeaker emits a volume velocity U_0 , as determined in the previous section. For a simple spherical source with volume velocity U_0 , the far-field pressure is given by [1]:

$$p(r) \simeq j\omega \rho_0 U_0 \frac{e^{-jkr}}{4\pi r}$$

It is possible to simply let $r = 1$ for this analysis without loss of generality because distance is only a function of the surroundings, not the loudspeaker. Also, because the transfer function magnitude is of primary interest, the exponential term, which has a unity magnitude, is omitted. Hence, the pressure response of the system is given by [1]:

$$\frac{p}{V_{IN}} = \frac{\rho_0 s U_0}{4\pi V_{IN}} = \frac{\rho_0 B l}{4\pi S_D R_E M_{AS}} H(s)$$

Where $H(s) = sG(s)$. In the following sections, design methods will focus on $|H(s)|^2$ rather than $H(s)$, which is given by:

$$|H(s)|^2 = \frac{\Omega^8}{\Omega^8 + (a_3^2 - 2a_2)\Omega^6 + (a_2^2 + 2 - 2a_1a_3)\Omega^4 + (a_1^2 - 2a_2)\Omega^2 + 1} \quad \Omega = \frac{\omega}{\omega_0}$$

This also implicitly ignores the constants in front of $|H(s)|$ since they simply scale the response and do not affect the shape of the frequency response curve.

Alignments

A popular way to determine the ideal parameters has been through the use of alignments. The concept of alignments is based upon filter theory. Filter development is a method of selecting the poles (and possibly zeros) of a transfer function to meet a particular design criterion. The criteria are the desired properties of a magnitude-squared transfer function, which in this case is $|H(s)|^2$. From any of the design criteria, the poles (and possibly zeros) of $|H(s)|^2$ are found, which can then be used to calculate the numerator and denominator. This is the "optimal" transfer

function, which has coefficients that are matched to the parameters of $|H(s)|^2$ to compute the appropriate values that will yield a design that meets the criteria.

There are many different types of filter designs, each which have trade-offs associated with them. However, this design is limited because of the structure of $|H(s)|^2$. In particular, it has the structure of a fourth-order high-pass filter with all zeros at $s = 0$. Therefore, only those filter design methods which produce a low-pass filter with only poles will be acceptable methods to use. From the traditional set of algorithms, only Butterworth and Chebyshev low-pass filters have only poles. In addition, another type of filter called a quasi-Butterworth filter can also be used, which has similar properties to a Butterworth filter. These three algorithms are fairly simple, thus they are the most popular. When these low-pass filters are converted to high-pass filters, the $s \rightarrow 1/s$ transformation produces s^8 in the numerator.

More details regarding filter theory and these relationships can be found in numerous resources, including [5].

Butterworth Alignment

The Butterworth algorithm is designed to have a *maximally flat* pass band. Since the slope of a function corresponds to its derivatives, a flat function will have derivatives equal to zero. Since as flat of a pass band as possible is optimal, the ideal function will have as many derivatives equal to zero as possible at $s = 0$. Of course, if all derivatives were equal to zero, then the function would be a constant, which performs no filtering.

Often, it is better to examine what is called the *loss function*. Loss is the reciprocal of gain, thus

$$|\hat{H}(s)|^2 = \frac{1}{|H(s)|^2}$$

The loss function can be used to achieve the desired properties, then the desired gain function is recovered from the loss function.

Now, applying the desired Butterworth property of maximal pass-band flatness, the loss function is simply a polynomial with derivatives equal to zero at $s = 0$. At the same time, the original polynomial must be of degree eight (yielding a fourth-order function). However, derivatives one through seven can be equal to zero if [3]

$$|\hat{H}(\Omega)|^2 = 1 + \Omega^8 \Rightarrow |H(\Omega)|^2 = \frac{1}{1 + \Omega^8}$$

With the high-pass transformation $\Omega \rightarrow 1/\Omega$,

$$|H(\Omega)|^2 = \frac{\Omega^8}{\Omega^8 + 1}$$

It is convenient to define $\Omega = \omega/\omega_{3dB}$, since $\Omega = 1 \Rightarrow |H(s)|^2 = 0.5$ or -3 dB. This definition allows the matching of coefficients for the $|H(s)|^2$ describing the loudspeaker response when $\omega_{3dB} = \omega_0$. From this matching, the following design equations are obtained [1]:

$$a_1 = a_3 = \sqrt{4 + 2\sqrt{2}} \quad a_2 = 2 + \sqrt{2}$$

Quasi-Butterworth Alignment

The quasi-Butterworth alignments do not have as well-defined of an algorithm when compared to the Butterworth alignment. The name “quasi-Butterworth” comes from the fact that the transfer functions for these responses appear similar to the Butterworth ones, with (in general) the addition of terms in the denominator. This will be illustrated below. While there are many types of quasi-Butterworth alignments, the simplest and most popular is the 3rd order alignment (QB3). The comparison of the QB3 magnitude-squared response against the 4th order Butterworth is shown below.

$$|H_{QB3}(\omega)|^2 = \frac{(\omega/\omega_{3dB})^8}{(\omega/\omega_{3dB})^8 + B^2(\omega/\omega_{3dB})^2 + 1} \quad |H_{B4}(\omega)|^2 = \frac{(\omega/\omega_{3dB})^8}{(\omega/\omega_{3dB})^8 + 1}$$

Notice that the case $B = 0$ is the Butterworth alignment. The reason that this QB alignment is called 3rd order is due to the fact that as B increases, the slope approaches 3 dec/dec instead of 4 dec/dec, as in 4th order Butterworth. This phenomenon can be seen in Figure 5.

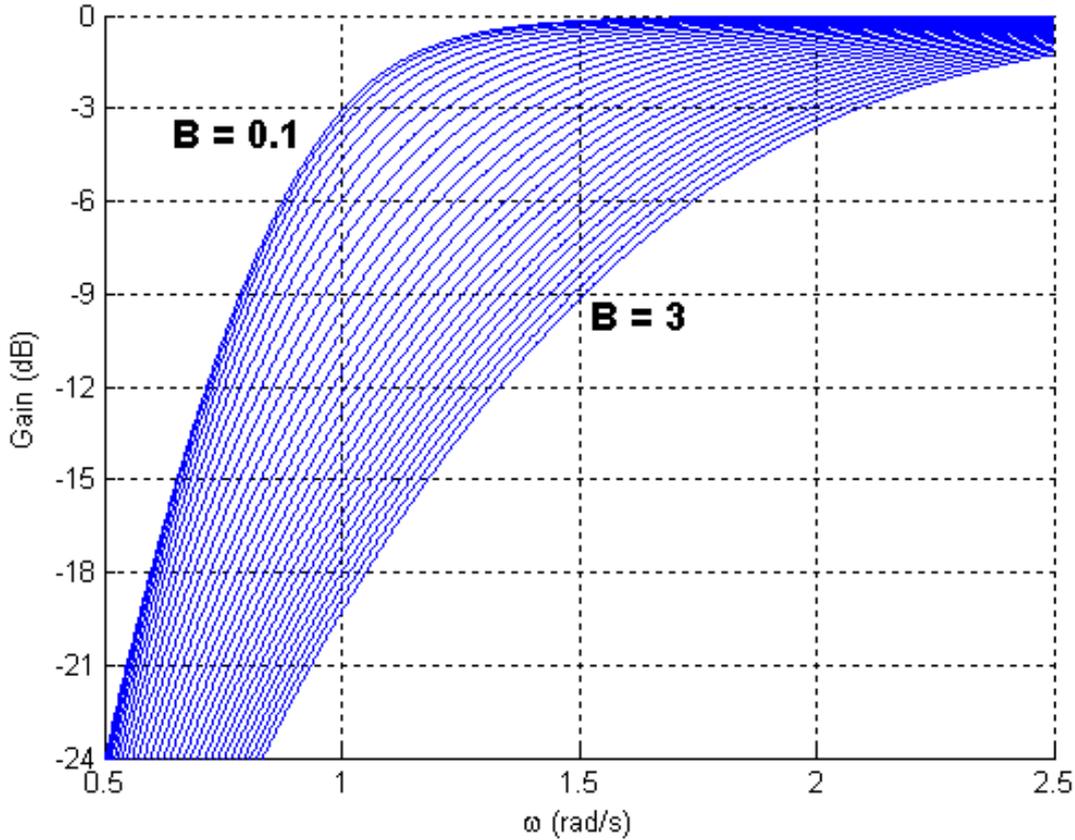


Figure 5: 3rd-Order Quasi-Butterworth Response for $0.1 \leq B \leq 3$

Equating the system response $|H(s)|^2$ with $|H_{QB3}(s)|^2$, the equations guiding the design can be found [1]:

$$B^2 = a_1^2 - 2a_2 \quad a_2^2 + 2 = 2a_1a_3 \quad a_3 = \sqrt{2a_2} \quad a_2 > 2 + \sqrt{2}$$

Chebyshev Alignment

The Chebyshev algorithm is an alternative to the Butterworth algorithm. For the Chebyshev response, the maximally-flat passband restriction is abandoned. Now, a *ripple*, or fluctuation is allowed in the pass band. This allows a steeper transition or roll-off to occur. In this type of application, the low-frequency response of the loudspeaker can be extended beyond what can be achieved by Butterworth-type filters. An example plot of a Chebyshev high-pass response with 0.5 dB of ripple against a Butterworth high-pass response for the same ω_{3dB} is shown below.

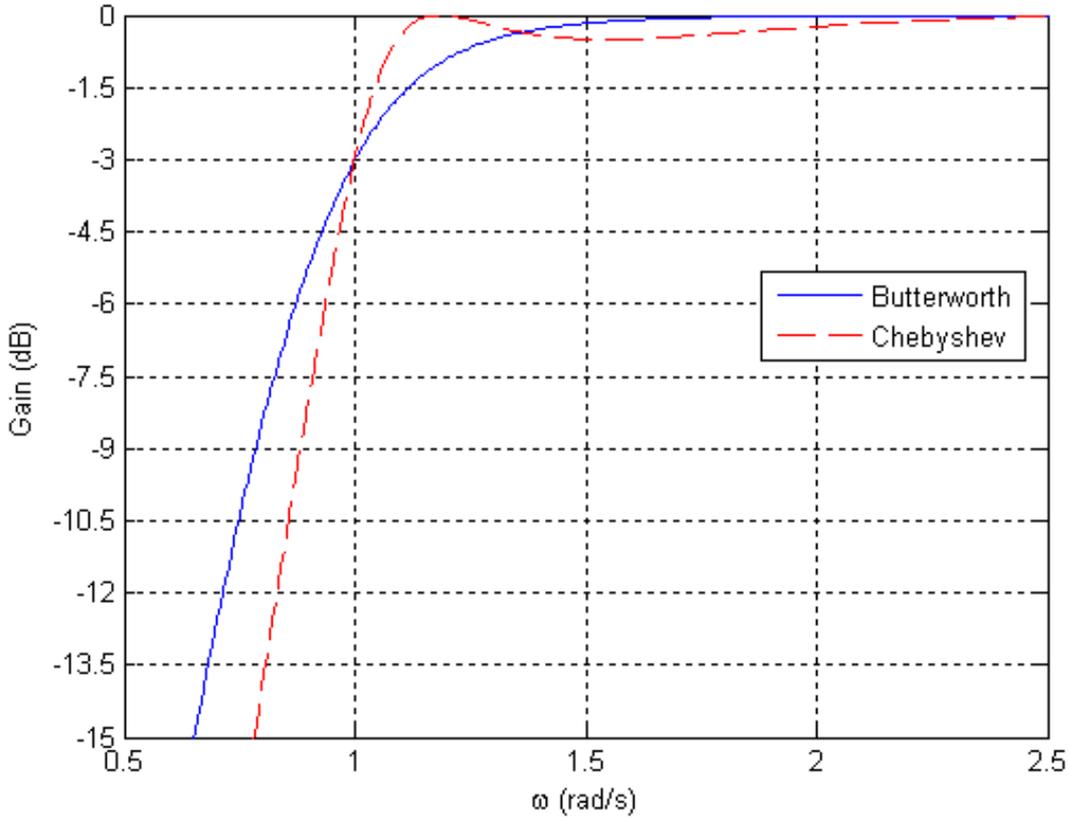


Figure 6: Chebyshev vs. Butterworth High-Pass Response.

The Chebyshev response is defined by [4]:

$$|\hat{H}(j\Omega)|^2 = 1 + \epsilon^2 C_n^2(\Omega)$$

$C_n(\Omega)$ is called the *Chebyshev polynomial* and is defined by [4]:

$$C_n(\Omega) = \begin{cases} \cos[n\cos^{-1}(\Omega)] & |\Omega| < 1 \\ \cosh[n\cosh^{-1}(\Omega)] & |\Omega| > 1 \end{cases}$$

Fortunately, Chebyshev polynomials satisfy a simple recursion formula [4]:

$$C_0(x) = 1 \quad C_1(x) = x \quad C_n(x) = 2xC_{n-1} - C_{n-2}$$

For more information on Chebyshev polynomials, see the Wolfram Mathworld: Chebyshev Polynomials [2] page.

When applying the high-pass transformation to the 4th order form of $|\hat{H}(j\Omega)|^2$, the desired response has the form [1]:

$$|H(j\Omega)|^2 = \frac{1 + \epsilon^2}{1 + \epsilon^2 C_4^2(1/\Omega)}$$

The parameter ϵ determines the ripple. In particular, the magnitude of the ripple is $10\log[1 + \epsilon^2]$ dB and can be chosen by the designer, similar to B in the quasi-Butterworth case. Using the recursion formula for $C_n(x)$,

$$C_4\left(\frac{1}{\Omega}\right) = 8\left(\frac{1}{\Omega}\right)^4 - 8\left(\frac{1}{\Omega}\right)^2 + 1$$

Applying this equation to $|H(j\Omega)|^2$ [1],

$$\Rightarrow |H(\Omega)|^2 = \frac{\frac{1+\epsilon^2}{64\epsilon^2}\Omega^8}{\frac{1+\epsilon^2}{64\epsilon^2}\Omega^8 + \frac{1}{4}\Omega^6 + \frac{5}{4}\Omega^4 - 2\Omega^2 + 1}$$

$$\Omega = \frac{\omega}{\omega_n} \quad \omega_n = \frac{\omega_{3dB}}{2} \sqrt{2 + \sqrt{2 + 2\sqrt{2 + \frac{1}{\epsilon^2}}}}$$

Thus, the design equations become [1]:

$$\omega_0 = \omega_n \sqrt[8]{\frac{64\epsilon^2}{1 + \epsilon^2}} \quad k = \tanh \left[\frac{1}{4} \sinh^{-1} \left(\frac{1}{\epsilon} \right) \right] \quad D = \frac{k^4 + 6k^2 + 1}{8}$$

$$a_1 = \frac{k\sqrt{4 + 2\sqrt{2}}}{\sqrt[4]{D}}, \quad a_2 = \frac{1 + k^2(1 + \sqrt{2})}{\sqrt{D}} \quad a_3 = \frac{a_1}{\sqrt{D}} \left[1 - \frac{1 - k^2}{2\sqrt{2}} \right]$$

Choosing the Correct Alignment

With all the equations that have already been presented, the question naturally arises, "Which one should I choose?" Notice that the coefficients a_1 , a_2 , and a_3 are not simply related to the parameters of the system response. Certain combinations of parameters may indeed invalidate one or more of the alignments because they cannot realize the necessary coefficients. With this in mind, general guidelines have been developed to guide the selection of the appropriate alignment. This is very useful if one is designing an enclosure to suit a particular transducer that cannot be changed.

The general guideline for the Butterworth alignment focuses on Q_L and Q_{TS} . Since the three coefficients a_1 , a_2 , and a_3 are a function of Q_L , Q_{TS} , h , and α , fixing one of these parameters yields three equations that uniquely determine the other three. In the case where a particular transducer is already given, Q_{TS} is essentially fixed. If the desired parameters of the enclosure are already known, then Q_L is a better starting point. In the case that the rigid requirements of the Butterworth alignment cannot be satisfied, the quasi-Butterworth alignment is often applied when Q_{TS} is not large enough.. The addition of another parameter, B , allows more flexibility in the design.

For Q_{TS} values that are too large for the Butterworth alignment, the Chebyshev alignment is typically chosen. However, the steep transition of the Chebyshev alignment may also be utilized to attempt to extend the bass response of the loudspeaker in the case where the transducer properties can be changed.

In addition to these three popular alignments, research continues in the area of developing new algorithms that can manipulate the low-frequency response of the bass-reflex enclosure. For example, a 5th order quasi-Butterworth alignment has been developed [6]. Another example [7] applies root-locus techniques to achieve results. In the modern age of high-powered computing, other researchers have focused their efforts in creating computerized optimization algorithms that can be modified to achieve a flatter response with sharp roll-off or introduce quasi-ripples which provide a boost in sub-bass frequencies [8].

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Appendix A: Equivalent Circuit Parameters

Name	Electrical Equivalent	Mechanical Equivalent	Acoustical Equivalent
Voice-Coil Resistance	R_E	$R_{ME} = \frac{(Bl)^2}{R_E}$	$R_{AE} = \frac{(Bl)^2}{R_E S_D^2}$
Driver (Speaker) Mass	See C_{MEC}	M_{MD}	$M_{AD} = \frac{M_{MD}}{S_D^2}$
Driver (Speaker) Suspension Compliance	$L_{CES} = (Bl)^2 C_{MS}$	C_{MS}	$C_{AS} = S_D^2 C_{MS}$
Driver (Speaker) Suspension Resistance	$R_{ES} = \frac{(Bl)^2}{R_{MS}}$	R_{MS}	$R_{AS} = \frac{R_{MS}}{S_D^2}$
Enclosure Compliance	$L_{CEB} = \frac{(Bl)^2 C_{AB}}{S_D^2}$	$C_{MB} = \frac{C_{AB}}{S_D^2}$	C_{AB}
Enclosure Air-Leak Losses	$R_{EL} = \frac{(Bl)^2}{S_D^2 R_{AL}}$	$R_{ML} = S_D^2 R_{AL}$	R_{AL}
Acoustic Mass of Port	$C_{MEP} = \frac{S_D^2 M_{AP}}{(Bl)^2}$	$M_{MP} = S_D^2 M_{AP}$	M_{AP}
Enclosure Mass Load	See C_{MEC}	See M_{MC}	M_{AB}
Low-Frequency Radiation Mass Load	See C_{MEC}	See M_{MC}	M_{A1}
Combination Mass Load	$C_{MEC} = \frac{S_D^2 M_{AC}}{(Bl)^2}$ $= \frac{S_D^2 (M_{AB} + M_{A1}) + M_{MD}}{(Bl)^2}$	$M_{MC} = S_D^2 (M_{AB} + M_{A1}) + M_{MD}$	$M_{AC} = M_{AD} + M_{AB} + M_{A1}$ $= \frac{M_{MD}}{S_D^2} + M_{AB} + M_{A1}$

Polymer-Film Acoustic Filters



Introduction

Acoustic filters are used in many devices such as mufflers, noise control materials (absorptive and reactive), and loudspeaker systems to name a few. Although the waves in simple (single-medium) acoustic filters usually travel in gases such as air and carbon-monoxide (in the case of automobile mufflers) or in materials such as fiberglass, polyvinylidene fluoride (PVDF) film, or polyethylene (Saran Wrap), there are also filters that couple two or three distinct media together to achieve a desired acoustic response. General information about basic acoustic filter design can be perused at the following wikibook page [Acoustic Filter Design & Implementation ^[1]]. The focus of this article will be on acoustic filters that use multilayer air/polymer film-coupled media as its acoustic medium for sound waves to propagate through; concluding with an example of how these filters can be used to detect and extrapolate audio frequency information in high-frequency "carrier" waves that carry an audio signal. However, before getting into these specific type of acoustic filters, we need to briefly discuss how sound waves interact with the medium(media) in which it travels and how these factors can play a role when designing acoustic filters.

Changes in Media Properties Due to Sound Wave Characteristics

As with any system being designed, the filter response characteristics of an acoustic filter are tailored based on the frequency spectrum of the input signal and the desired output. The input signal may be infrasonic (frequencies below human hearing), sonic (frequencies within human hearing range), or ultrasonic (frequencies above human hearing range). In addition to the frequency content of the input signal, the density, and, thus, the characteristic impedance of the medium (media) being used in the acoustic filter must also be taken into account. In general, the characteristic impedance Z_0 for a particular medium is expressed as...

$$Z_0 = \pm \rho_0 c \text{ (Pa} \cdot \text{s/m)}$$

where

$$\begin{aligned} \pm \rho_0 &= \text{(equilibrium) density of medium } \text{(kg/m}^3\text{)} \\ c &= \text{speed of sound in medium } \text{(m/s)} \end{aligned}$$

The characteristic impedance is important because this value simultaneously gives an idea of how fast or slow particles will travel as well as how much mass is "weighting down" the particles in the medium (per unit area or volume) when they are excited by a sound source. The speed in which sound travels in the medium needs to be taken into consideration because this factor can ultimately affect the time response of the filter (i.e. the output of the filter may not radiate or attenuate sound fast or slow enough if not designed properly). The intensity I_A of a sound wave

is expressed as...

$$I_A = \frac{1}{T} \int_0^T p u \, dt = \pm \frac{P^2}{2\rho_0 c} \quad (W/m^2).$$

I_A is interpreted as the (time-averaged) rate of energy transmission of a sound wave through a unit area normal to the direction of propagation, and this parameter is also an important factor in acoustic filter design because the characteristic properties of the given medium can change relative to intensity of the sound wave traveling through it. In other words, the reaction of the particles (atoms or molecules) that make up the medium will respond differently when the intensity of the sound wave is very high or very small relative to the size of the control area (i.e. dimensions of the filter, in this case). Other properties such as the elasticity and mean propagation velocity (of a sound wave) can change in the acoustic medium as well, but focusing on frequency, impedance, and/or intensity in the design process usually takes care of these other parameters because most of them will inevitably be dependent on the aforementioned properties of the medium.

Why Coupled Acoustic Media in Acoustic Filters?

In acoustic transducers, media coupling is employed in acoustic transducers to either increase or decrease the impedance of the transducer, and, thus, control the intensity and speed of the signal acting on the transducer while converting the incident wave, or initial excitation sound wave, from one form of energy to another (e.g. converting acoustic energy to electrical energy). Specifically, the impedance of the transducer is augmented by inserting a solid structure (not necessarily rigid) between the transducer and the initial propagation medium (e.g. air). The reflective properties of the inserted medium is exploited to either increase or decrease the intensity and propagation speed of the incident sound wave. It is the ability to alter, and to some extent, control, the impedance of a propagation medium by (periodically) inserting (a) solid structure(s) such as thin, flexible films in the original medium (air) and its ability to concomitantly alter the frequency response of the original medium that makes use of multilayer media in acoustic filters attractive. The reflection factor and transmission factor \hat{R} and \hat{T} , respectively, between two media, expressed as...

$$\hat{R} = \frac{\text{pressure of reflected portion of incident wave}}{\text{pressure of incident wave}} = \frac{\rho c - Z_{in}}{\rho c + Z_{in}}$$

and

$$\hat{T} = \frac{\text{pressure of transmitted portion of incident wave}}{\text{pressure of incident wave}} = 1 + \hat{R},$$

are the tangible values that tell how much of the incident wave is being reflected from and transmitted through the junction where the media meet. Note that Z_{in} is the (total) input impedance seen by the incident sound wave upon just entering an air-solid acoustic media layer. In the case of multiple air-columns as shown in Fig. 2, Z_{in} is the aggregate impedance of each air-column layer seen by the incident wave at the input. Below in Fig. 1, a simple illustration explains what happens when an incident sound wave propagating in medium (1) and comes in contact with medium (2) at the junction of the both media ($x=0$), where the sound waves are represented by vectors.

As mentioned above, an example of three successive air-solid acoustic media layers is shown in Fig. 2 and the electroacoustic equivalent circuit for Fig. 2 is shown in Fig. 3 where $L = \rho_s h_s =$ (density of solid material)(thickness of solid material) = unit-area (or volume) mass, $Z = \rho c =$ characteristic acoustic impedance of medium, and $\beta = k = \omega/c =$ wavenumber. Note that in the case of a multilayer, coupled acoustic medium in an acoustic filter, the impedance of each air-solid section is calculated by using the following general purpose impedance ratio equation (also referred to as transfer matrices)...

$$\frac{Z_a}{Z_0} = \frac{\left(\frac{Z_b}{Z_0}\right) + j \tan(kd)}{1 + j \left(\frac{Z_b}{Z_0}\right) \tan(kd)}$$

where Z_b is the (known) impedance at the edge of the solid of an air-solid layer (on the right) and Z_a is the (unknown) impedance at the edge of the air column of an air-solid layer.

Effects of High-Intensity, Ultrasonic Waves in Acoustic Media in Audio Frequency Spectrum

When an ultrasonic wave is used as a carrier to transmit audio frequencies, three audio effects are associated with extrapolating the audio frequency information from the carrier wave: (a) beating effects, (b) parametric array effects, and (c) radiation pressure.

Beating occurs when two ultrasonic waves with distinct frequencies f_1 and f_2 propagate in the same direction, resulting in amplitude variations which consequently make the audio signal information go in and out of phase, or "beat", at a frequency of $f_1 - f_2$.

Parametric array [[2]] effects occur when the intensity of an ultrasonic wave is so high in a particular medium that the high displacements of particles (atoms) per wave cycle changes properties of that medium so that it influences parameters like elasticity, density, propagation velocity, etc. in a non-linear fashion. The results of parametric array effects on modulated, high-intensity, ultrasonic waves in a particular medium (or coupled media) is the generation and propagation of audio frequency waves (not necessarily present in the original audio information) that are generated in a manner similar to the nonlinear process of amplitude demodulation commonly inherent in diode circuits (when diodes are forward biased).

Another audio effect that arises from high-intensity ultrasonic beams of sound is a static (DC) pressure called radiation pressure. Radiation pressure is similar to parametric array effects in that amplitude variations in the signal give rise to audible frequencies via amplitude demodulation. However, unlike parametric array effects, radiation pressure fluctuations that generate audible signals from amplitude demodulation can occur due to any low-frequency modulation and not just from pressure fluctuations occurring at the modulation frequency ω_M or beating frequency $f_1 - f_2$.

An Application of Coupled Media in Acoustic Filters

Figs. 1 - 3 were all from a research paper entitled New Type of Acoustics Filter Using Periodic Polymer Layers for Measuring Audio Signal Components Excited by Amplitude-Modulated High_Intensity Ultrasonic Waves [3] submitted to the Audio Engineering Society (AES) by Minoru Todo, Primary Innovator at Measurement Specialties, Inc., in the October 2005 edition of the AES Journal. Figs. 4 and 5 below, also from this paper, are illustrations of test setups referred to in this paper. Specifically, Fig. 4 is a test setup used to measure the transmission (of an incident ultrasonic sound wave) through the acoustic filter described by Figs. 1 and 2. Fig. 5 is a block diagram of the test setup used for measuring radiation pressure, one of the audio effects mentioned in the previous section. It turns out that out of all of the audio effects mentioned in the previous section that are caused by high-intensity ultrasonic waves propagating in a medium, sound waves produced from radiated pressure are the hardest to detect when microphones and preamplifiers are used in the detection/receiver system. Although nonlinear noise artifacts occur due to overloading of the preamplifier present in the detection/receiver system, the bulk of the nonlinear noise comes from the inherent nonlinear noise properties of microphones. This is true because all microphones, even specialized measurement microphones designed for audio spectrum measurements that have sensitivity well beyond the threshold of hearing, have nonlinearities artifacts that (periodically) increase in magnitude with respect to increase at ultrasonic frequencies. These nonlinearities essentially mask the radiation pressure generated because the magnitude of these nonlinearities are orders of magnitude greater than the radiation pressure. The acoustic (low-pass) filter referred to in this paper was designed in order to filter out the "detrimental" ultrasonic wave that was inducing high nonlinear noise artifacts in the measurement microphones. The high-intensity, ultrasonic wave was producing radiation pressure (which is audible) within the initial acoustic medium (i.e. air). By filtering out the ultrasonic wave, the measurement microphone would only detect the audible

radiation pressure that the ultrasonic wave was producing in air. Acoustic filters like these could possibly be used to detect/receive any high-intensity, ultrasonic signal that may carry audio information which may need to be extrapolated with an acceptable level of fidelity.

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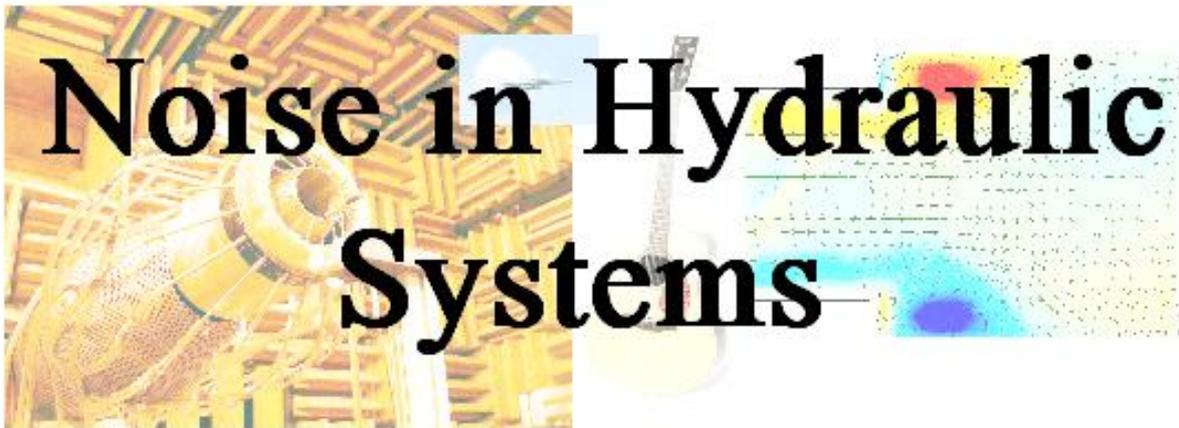
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Created by Valdez L. Gant

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Noise in Hydraulic Systems



Noise in Hydraulic Systems

Hydraulic systems are the most preferred source of power transmission in most of the industrial and mobile equipments due to their power density, compactness, flexibility, fast response and efficiency. The field hydraulics and pneumatics is also known as 'Fluid Power Technology'. Fluid power systems have a wide range of applications which include industrial, off-road vehicles, automotive systems, and aircraft. But, one of the main problems with the hydraulic systems is the noise generated by them. The health and safety issues relating to noise have been recognized for many years and legislation is now placing clear demands on manufacturers to reduce noise levels [1]. Hence, noise reduction in hydraulic systems demands lot of attention from the industrial as well as academic researchers. It needs a good understanding of how the noise is generated and propagated in a hydraulic system in order to reduce it.

Sound in fluids

The speed of sound in fluids can be determined using the following relation.

$$c = \sqrt{\frac{K}{\rho}}$$

where K - fluid bulk modulus, ρ - fluid density, c - velocity of sound

Typical value of bulk modulus range from **2e9 to 2.5e9 N/m²**. For a particular oil, with a density of **889 kg/m³**,

$$\text{speed of sound } c = \sqrt{\frac{2e9}{889}} = 1499.9m/s$$

Source of Noise

The main source of noise in hydraulic systems is the pump which supplies the flow. Most of the pumps used are positive displacement pumps. Of the positive displacement pumps, axial piston swash plate type is mostly preferred due to their reliability and efficiency.

The noise generation in an axial piston pump can be classified under two categories (i) fluidborne noise and (ii) Structureborne noise

Fluidborne Noise (FBN)

Among the positive displacement pumps, highest levels of FBN are generated by axial piston pumps and lowest levels by screw pumps and in between these lie the external gear pump and vane pump [1]. The discussion in this page is mainly focused on **axial piston swash plate type pumps**. An axial piston pump has a fixed number of displacement chambers arranged in a circular pattern separated from each other by an angular pitch equal to $\phi = \frac{360}{n}$ where n is the number of displacement chambers. As each chamber discharges a specific volume of fluid,

the discharge at the pump outlet is sum of all the discharge from the individual chambers. The discontinuity in flow between adjacent chambers results in a kinematic flow ripple. The amplitude of the kinematic ripple can be theoretical determined given the size of the pump and the number of displacement chambers. The kinematic ripple is the main cause of the fluidborne noise. The kinematic ripples is a theoretical value. The actual **flow ripple** at the pump outlet is much larger than the theoretical value because the **kinematic ripple** is combined with a **compressibility component** which is due to the fluid compressibility. These ripples (also referred as flow pulsations) generated at the pump are transmitted through the pipe or flexible hose connected to the pump and travel to all parts of the hydraulic circuit.

The pump is considered an ideal flow source. The pressure in the system will be decided by resistance to the flow or otherwise known as system load. The flow pulsations result in pressure pulsations. The pressure pulsations are superimposed on the mean system pressure. Both the **flow and pressure pulsations** easily travel to all part of the circuit and affect the performance of the components like control valve and actuators in the system and make the component vibrate, sometimes even resonate. This vibration of system components adds to the noise generated by the flow pulsations. The transmission of FBN in the circuit is discussed under transmission below.

A typical axial piston pump with 9 pistons running at 1000 rpm can produce a sound pressure level of more than 70 dBs.

Structureborne Noise (SBN)

In swash plate type pumps, the main sources of the structureborne noise are the fluctuating forces and moments of the swash plate. These fluctuating forces arise as a result of the varying pressure inside the displacement chamber. As the displacing elements move from suction stroke to discharge stroke, the pressure varies accordingly from few bars to few hundred bars. These pressure changes are reflected on the displacement elements (in this case, pistons) as forces and these force are exerted on the swash plate causing the swash plate to vibrate. This vibration of the swash plate is the main cause of **structureborne noise**. There are other components in the system which also vibrate and lead to structureborne noise, but the swash is the major contributor.

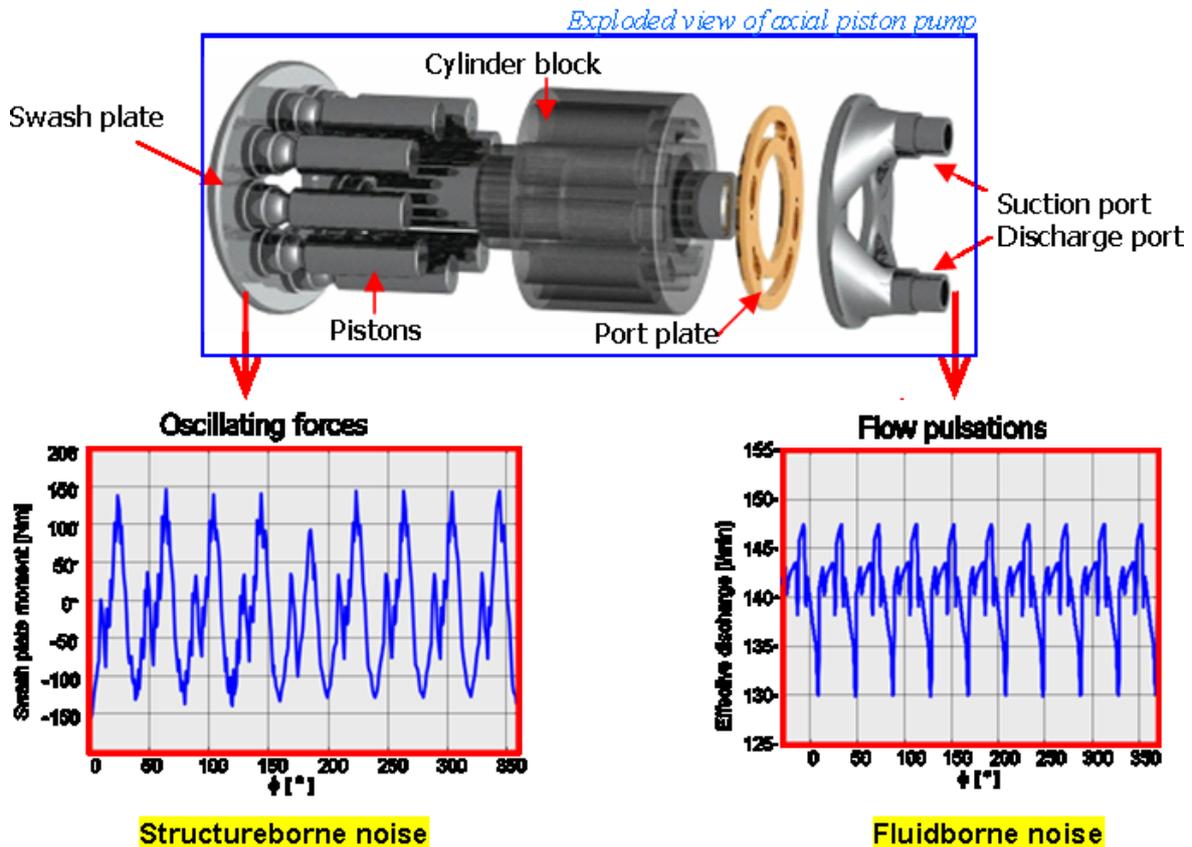


Fig. 1 shows an exploded view of axial piston pump. Also the flow pulsations and the oscillating forces on the swash plate, which cause FBN and SBN respectively are shown for one revolution of the pump.

Transmission

FBN

The transmission of FBN is a complex phenomenon. Over the past few decades, considerable amount of research had gone into mathematical modeling of pressure and flow transient in the circuit. This involves the solution of wave equations, with piping treated as a distributed parameter system known as a transmission line [1] & [3].

Lets consider a simple pump-pipe-loading valve circuit as shown in Fig. 2. The pressure and flow ripple at any location in the pipe can be described by the relations:

$$P = Ae^{-ikx} + Be^{+ikx} \dots\dots(1)$$

$$Q = \frac{1}{Z_0}(Ae^{-ikx} - Be^{+ikx}) \dots\dots(2)$$

where A and B are frequency dependent complex coefficients which are directly proportional to pump (source) flow ripple, but also functions of the source impedance Z_s , characteristic impedance of the pipe Z_0 and the

termination impedance Z_t . These impedances, usually vary as the system operating pressure and flow rate changes, can be determined experimentally.

For complex systems with several system components, the pressure and flow ripples are estimated using the transformation matrix approach. For this, the system components can be treated as lumped impedances (a throttle valve or accumulator), or distributed impedances (flexible hose or silencer). Various software packages are available today to predict the pressure pulsations.

SBN

The transmission of SBN follows the classic source-path-noise model. The vibrations of the swash plate, the main cause of SBN, are transferred to the pump casing which encloses all the rotating group in the pump including displacement chambers (also known as cylinder block), pistons, and the swash plate. The pump case, apart from vibrating itself, transfers the vibration down to the mount on which the pump is mounted. The mount then passes the vibrations down to the main mounted structure or the vehicle. Thus the SBN is transferred from the swash plate to the main structure or vehicle via pumpcasing and mount.

Some of the machine structures, along the path of transmission, are good at transmitting this vibrational energy and they even resonate and reinforce it. By converting only a fraction of 1% of the pump structureborne noise into sound, a member in the transmission path could radiate more ABN than the pump itself [4].

Airborne noise (ABN)

Both FBN and SBN, impart high fatigue loads on the system components and make them vibrate. All of these vibrations are radiated as **airborne noise** and can be heard by a human operator. Also, the flow and pressure pulsations make the system components such as a control valve to resonate. This vibration of the particular component again radiates airborne noise.

Noise reduction

The reduction of the noise radiated from the hydraulic system can be approached in two ways.

(i) **Reduction at Source** - which is the reduction of noise at the pump. A large amount of open literature are available on the reduction techniques with some techniques focusing on reducing FBN at source and others focusing on SBN. Reduction in FBN and SBN at the source has a large influence on the ABN that is radiated. Even though, a lot of progress had been made in reducing the FBN and SBN separately, the problem of noise in hydraulic systems is not fully solved and lot need to be done. The reason is that the FBN and SBN are interrelated, in a sense that, if one tried to reduce the FBN at the pump, it tends to affect the SBN characteristics. Currently, one of the main researches in noise reduction in pumps, is a systematic approach in understanding the coupling between FBN and SBN and targeting them simultaneously instead of treating them as two separate sources. Such a unified approach, demands not only well trained researchers but also sophisticated computer based mathematical model of the pump which can accurately output the necessary results for optimization of pump design. The amplitude of fluid pulsations can be reduced, at the source, with the use of an hydraulic attenuator(5).

(ii) **Reduction at Component level** - which focuses on the reduction of noise from individual component like hose, control valve, pump mounts and fixtures. This can be accomplished by a suitable design modification of the component so that it radiates least amount of noise. Optimization using computer based models can be one of the ways.

Hydraulic System noise

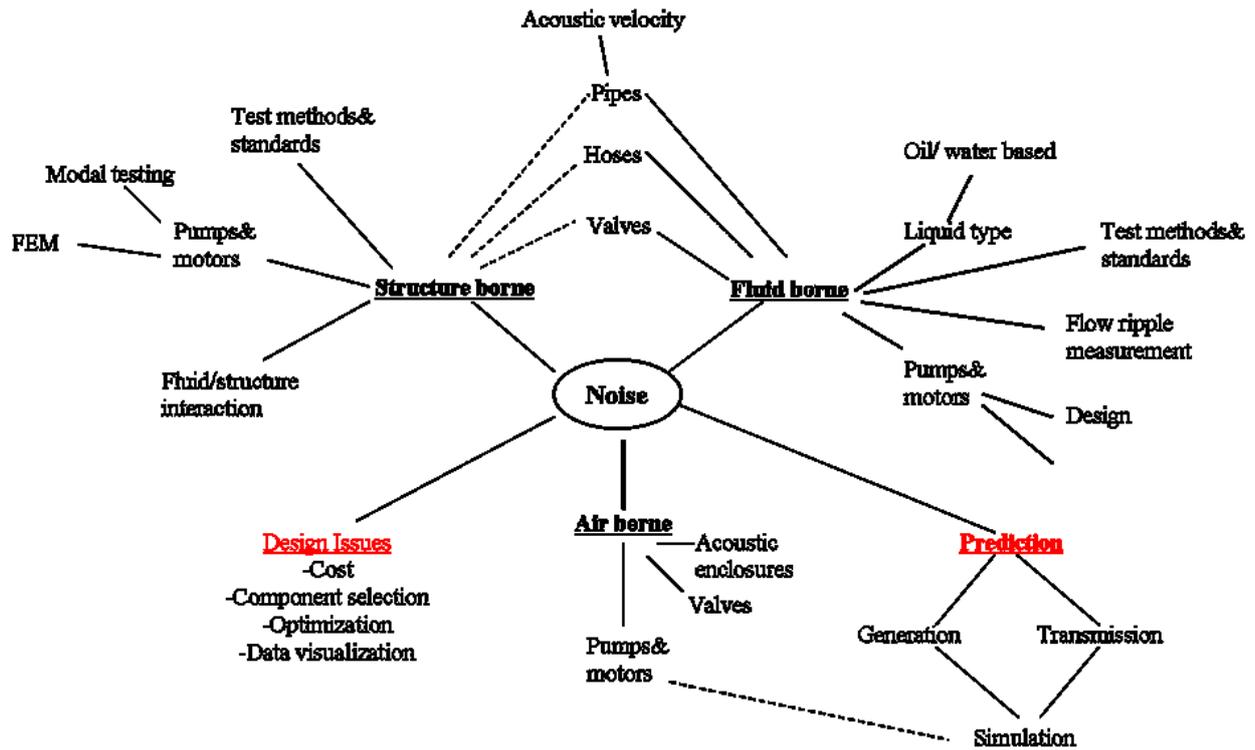


Fig.3 Domain of hydraulic system noise generation and transmission (Figure recreated from [1])

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Noise from Cooling Fans



Proposal

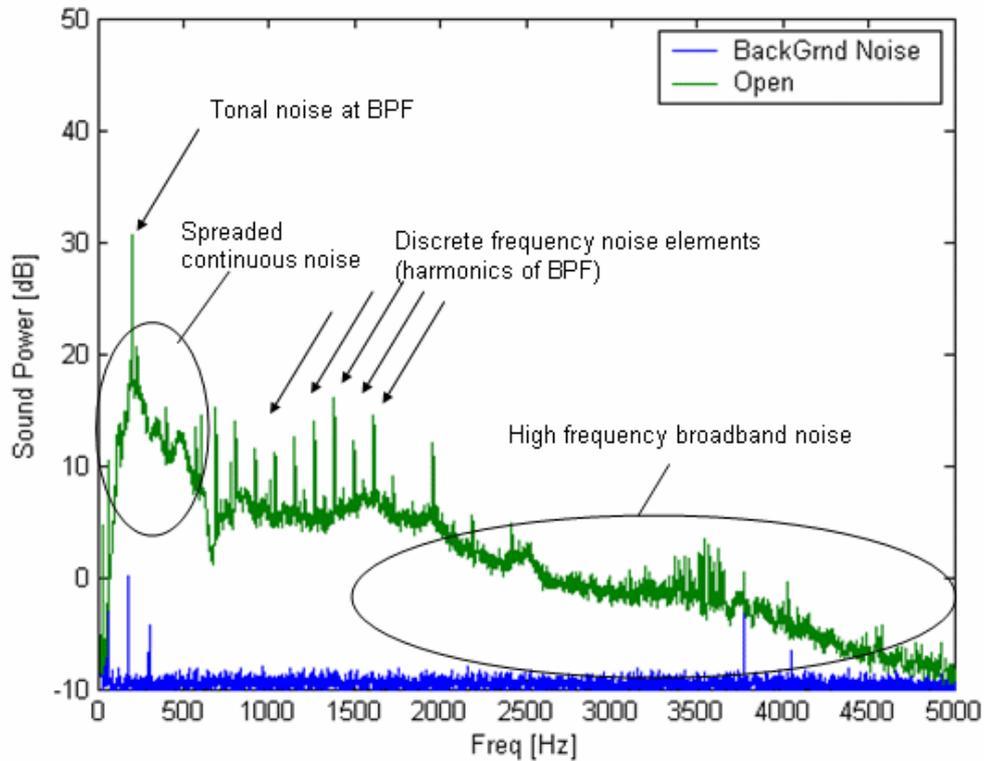
As electric/electronic devices get smaller and functional, the noise of cooling device becomes important. This page will explain the origins of noise generation from small axial cooling fans used in electronic goods like desktop/laptop computers. The source of fan noises includes aerodynamic noise as well as operating sound of the fan itself. This page will be focused on the aerodynamic noise generation mechanisms.

Introduction

If one opens a desktop computer, they may find three (or more) fans. For example, a fan is typically found on the heat sink of the CPU, in the back panel of the power supply unit, on the case ventilation hole, on the graphics card, and even on the motherboard chipset if it is a recent one. Computer noise which annoys many people is mostly due to cooling fans, if the hard drive(s) is fairly quiet. When Intel Pentium processors were first introduced, there was no need to have a fan on the CPU, however, contemporary CPUs cannot function even for several seconds without a cooling fan. As CPU densities increase, the heat transfer for nominal operation requires increased airflow, which causes more and more noise. The type of fans commonly used in desktop computers are axial fans, and centrifugal blowers in laptop computers. Several fan types are shown here (pdf format) ^[1]. Different fan types have different characteristics of noise generation and performance. The axial flow fan is mainly considered in this page.

Noise Generation Mechanisms

The figure below shows a typical noise spectrum of a 120 mm diameter electronic device cooling fan. One microphone was used at a point 1 m from the upstream side of the fan. The fan has 7 blades, 4 struts for motor mounting and operates at 13V. Certain amount of load is applied. The blue plot is background noise of anechoic chamber, and the green one is sound loudness spectrum when the fan is running.



(*BPF = Blade Passing Frequency)

Each noise elements shown in this figure is caused by one or more of following generation mechanisms.

Blade Thickness Noise - Monopole (But very weak)

Blade thickness noise is generated by volume displacement of fluid. A fan blade has its thickness and volume. As the rotor rotates, the volume of each blade displaces fluid volume, then they consequently fluctuate pressure of near field, and noise is generated. This noise is tonal at the running frequency and generally very weak for cooling fans, because their RPM is relatively low. Therefore, thickness of fan blades hardly affects to electronic cooling fan noise. (This kind of noise can become severe for high speed turbomachines like helicopter rotors.)

Tonal Noise by Aerodynamic Forces - Dipole

Uniform Inlet Flow (Negligible)

The sound generation due to uniform and steady aerodynamic force has a characteristic very similar to the blade thickness noise. It is very weak for low speed fans, and depends on fan RPM. Since some steady blade forces are necessary for a fan to do its duty even in an ideal condition, this kind of noise is impossible to avoid. It is known that this noise can be reduced by increasing the number of blades.

Non-uniform Inlet Flow

Non-uniform (still steady) inlet flow causes non-uniform aerodynamic forces on blades as their angular positions change. This generates noise at blade passing frequency and its harmonics. It is one of the major noise sources of electronic cooling fans.

Rotor-Casing interaction

If the fan blades are very close to a structure which is not symmetric, unsteady interaction forces to blades are generated. Then the fan experiences a similar running condition as lying in non-uniform flow field. See Acoustics/Rotor Stator Interactions for details.

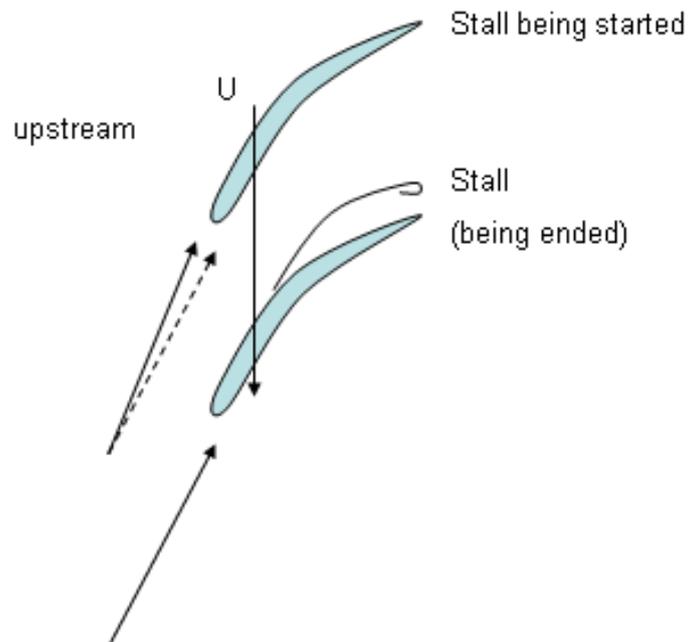
Impulsive Noise (Negligible)

This noise is caused by the interaction between a blade and blade-tip-vortex of the preceding blade, and not severe for cooling fans.

Rotating Stall

[Click here to read the definition and an aerodynamic description of stall.](#)

The noise due to stall is a complex phenomenon that occurs at low flow rates. For some reason, if flow is locally disturbed, it can cause stall on one of the blades. As a result, the upstream passage on this blade is partially blocked. Therefore, the mean flow is diverted away from this passage. This causes increasing of the angle of attack on the closest blade at the upstream side of the originally stalled blade, the flow is again stalled there. On the other hand, the other side of the first blade is un-stalled because of reduction of flow angle.



repeatedly, the stall cell turns around the blades at about 30~50% of the running frequency, and the direction is opposite to the blades. This series of phenomenon causes unsteady blade forces, and consequently generates noise and vibrations.

Non-uniform Rotor Geometry

Asymmetry of rotor causes noise at the rotating frequency and its harmonics (not blade passing frequency obviously), even when the inlet flow is uniform and steady.

Unsteady Flow Field

Unsteady flow causes random forces on the blades. It spreads the discrete spectrum noises and makes them continuous. In case of low-frequency variation, the spread continuous spectral noise is around rotating frequency, and narrowband noise is generated. The stochastic velocity fluctuations of inlet flow generates broadband noise spectrum. The generation of random noise components is covered by the following sections.

Random Noise by Unsteady Aerodynamic Forces

Turbulent Boundary Layer

Even in the steady and uniform inlet flow, there exist random force fluctuations on the blades. That is from turbulent blade boundary layer. Some noise is generated for this reason, but dominant noise is produced by the boundary layer passing the blade trailing edge. The blade trailing edges scatter the non-propagating near-field pressure into a propagatable sound field.

Incident Turbulent

Velocity fluctuations of the intake flow with a stochastic time history generate random forces on blades, and a broadband spectrum noise.

Vortex Shedding

For some reason, a vortex can separate from a blade. Then the circulating flow around the blade starts to be changed. This causes non-uniform forces on blades, and noises. A classical example for this phenomenon is 'Karman vortex street' ^[2]. (some images and animations ^[3].) Vortex shedding mechanism can occur in a laminar boundary layer of low speed fan and also in a turbulent boundary layer of high frequency fan.

Flow Separation

Flow separation causes stall explained above. This phenomenon can cause random noise, which spreads all the discrete spectrum noises, and turns the noise into broadband.

Tip Vortex

Since cooling fans are ducted axial flow machines, the annular gap between the blade tips and the casing is important parameter for noise generation. While rotating, there is another flow through the annular gap due to pressure difference between upstream and downstream of fan. Because of this flow, tip vortex is generated through the gap, and broadband noise increases as the annular gap gets bigger.

Installation Effects

Once a fan is installed, even though the fan is well designed acoustically, unexpected noise problem can come up. It is called as installation effects, and two types are applicable to cooling fans.

Effect of Inlet Flow Conditions

A structure that affects the inlet flow of a fan causes installation effects. For example Hoppe & Neise [3] showed that with and without a bellmouth nozzle at the inlet flange of 500mm fan can change the noise power by 50dB (This application is for much larger and noisier fan though).

Acoustic Loading Effect

This effect is shown on duct system applications. Some high performance graphic cards apply duct system for direct exhaustion.

The sound power generated by a fan is not only a function of its impeller speed and operating condition, but also depends on the acoustic impedances of the duct systems connected to its inlet and outlet. Therefore, fan and duct system should be matched not only for aerodynamic noise reasons but also because of acoustic considerations.

Closing Comment

Noise reduction of cooling fans has some restrictions:

1. Active noise control is not economically effective. 80mm cooling fans are only 5~10 US dollars. It is only applicable for high-end electronic products.
2. Restricting certain aerodynamic phenomenon for noise reduction can cause serious performance reduction of the fan. Increasing RPM of the fan is of course much more dominant factor for noise.

Different stories of fan noise are introduced at some of the linked sites below like active RPM control or noise comparison of various bearings used in fans.

Links to Interesting Sites about Fan Noise

Some practical issue of PC noise are presented at the following sites.

Cooling Fan Noise Comparison - Sleeve Bearing vs. Ball Bearing (pdf format) ^[4]

Brief explanation of fan noise origins and noise reduction suggestions ^[5]

Effect of sweep angle comparison ^[6]

Comparisons of noise from various 80mm fans ^[7]

Noise reduction of a specific desktop case ^[8]

Noise reduction of another specific desktop case ^[9]

Informal study for noise from CPU cooling fan ^[10]

Informal study for noise from PC case fans ^[11]

Active fan speed optimizers for minimum noise from desktop computers ^[12]

- Some general fan noise reduction techniques ^[13]
- Brüel & Kjær - acoustic testing device company. Various applications are presented in "applications" tap ^[14]
- "Fan Noise Solutions" ^[15] by cpeema
- How To Build A Computer/Quiet PC

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Piezoelectric Transducers



Introduction

Piezoelectricity from the Greek word "piezo" means pressure electricity. Certain crystalline substances generate electric charges under mechanical stress and conversely experience a mechanical strain in the presence of an electric field. The piezoelectric effect describes a situation where the transducing material senses input mechanical vibrations and produces a charge at the frequency of the vibration. An AC voltage causes the piezoelectric material to vibrate in an oscillatory fashion at the same frequency as the input current.

Quartz is the best known single crystal material with piezoelectric properties. Strong piezoelectric effects can be induced in materials with an ABO₃, Perovskite crystalline structure. 'A' denotes a large divalent metal ion such as lead and 'B' denotes a smaller tetravalent ion such as titanium or zirconium.

For any crystal to exhibit the piezoelectric effect, its structure must have no center of symmetry. Either a tensile or compressive stress applied to the crystal alters the separation between positive and negative charge sites in the cell causing a net polarization at the surface of the crystal. The polarization varies directly with the applied stress and is direction dependent so that compressive and tensile stresses will result in electric fields of opposite voltages.

Vibrations & Displacements

Piezoelectric ceramics have non-centrosymmetric unit cells below the Curie temperature and centrosymmetric unit cells above the Curie temperature. Non-centrosymmetric structures provide a net electric dipole moment. The dipoles are randomly oriented until a strong DC electric field is applied causing permanent polarization and thus piezoelectric properties.

A polarized ceramic may be subjected to stress causing the crystal lattice to distort changing the total dipole moment of the ceramic. The change in dipole moment due to an applied stress causes a net electric field which varies linearly with stress.

Dynamic Performance

The dynamic performance of a piezoelectric material relates to how it behaves under alternating stresses near the mechanical resonance. The parallel combination of C2 with L1, C1, and R1 in the equivalent circuit below control the transducers reactance which is a function of frequency.

Frequency Response

The graph below shows the impedance of a piezoelectric transducer as a function of frequency. The minimum value at f_n corresponds to the resonance while the maximum value at f_m corresponds to anti-resonance. ^{Superscript text}*Italic text*

Resonant Devices

Non resonant devices may be modeled by a capacitor representing the capacitance of the piezoelectric with an impedance modeling the mechanically vibrating system as a shunt in the circuit. The impedance may be modeled as a capacitor in the non resonant case which allows the circuit to reduce to a single capacitor replacing the parallel combination.

For resonant devices the impedance becomes a resistance or static capacitance at resonance. This is an undesirable effect. In mechanically driven systems this effect acts as a load on the transducer and decreases the electrical output. In electrically driven systems this effect shunts the driver requiring a larger input current. The adverse effect of the static capacitance experienced at resonant operation may be counteracted by using a shunt or series inductor resonating with the static capacitance at the operating frequency.

Applications

Mechanical Measurement

Because of the dielectric leakage current of piezoelectrics they are poorly suited for applications where force or pressure have a slow rate of change. They are, however, very well suited for highly dynamic measurements that might be needed in blast gauges and accelerometers.

Ultrasonic

High intensity ultrasound applications utilize half wavelength transducers with resonant frequencies between 18 kHz and 45 kHz. Large blocks of transducer material is needed to generate high intensities which makes manufacturing difficult and is economically impractical. Also, since half wavelength transducers have the highest stress amplitude in the center the end sections act as inert masses. The end sections are often replaced with metal plates possessing a much higher mechanical quality factor giving the composite transducer a higher mechanical quality factor than a single-piece transducer.

The overall electro-acoustic efficiency is:

$$\begin{aligned} Q_{m0} &= \text{unloaded mechanical quality factor} \\ Q_E &= \text{electric quality factor} \\ Q_L &= \text{quality factor due to the acoustic load alone} \end{aligned}$$

The second term on the right hand side is the dielectric loss and the third term is the mechanical loss.

Efficiency is maximized when:

$$\text{then: } \eta_{\max} = 1 - \frac{2}{k_{\text{eff}} \sqrt{Q_E Q_{m0}}} \quad \left(k_{\text{eff}} \sqrt{Q_E Q_{m0}} \ll 1 \right)$$

The maximum ultrasonic efficiency is described by: $I_{w\max} = \frac{1}{2} (\omega_0 u_l)_{\max}^2 \rho_{tw} v_w (W/m^2)$

Applications of ultrasonic transducers include:

- Welding of plastics
- Atomization of liquids
- Ultrasonic drilling
- Ultrasonic cleaning
- Ultrasonic foils in the paper machine wet end for more uniform fibre distribution
- Ultrasound
- Non-destructive testing
- etc.

More Information and Source of Information

MorganElectroCeramics ^[1]

Ultra Technology ^[2]

References

[1] <http://morganelectroceramics.com>

[2] <http://www.ultratechnology.se>

Generation and Propagation of Thunder

Thunder is the sound made by lightning. Depending on the nature of the lightning and distance of the listener, thunder can range from a sharp, loud crack to a long, low rumble (brontide). The sudden increase in pressure and temperature from lightning produces rapid expansion of the air surrounding and within a bolt of lightning. In turn, this expansion of air creates a sonic shock wave which produces the sound of thunder, often referred to as a *clap*, *crack*, or *peal of thunder*. The distance of the lightning can be calculated by the listener depending on when the sound is heard vs. the vision of the lightning strike.



A lightning bolt.

Etymology

The *d* in Modern English *thunder* (from earlier Old English *þunor*) is epenthetic, and is now found as well in Modern Dutch *donder* (cp Middle Dutch *donre*, and Old Norse *þorr*, Old Frisian *þuner*, Old High German *donar* descended from Proto-Germanic **þunraz*). In Latin the term was *tonare* "to thunder". The name of the Germanic god Thor comes from the Old Norse word for thunder.

NOTE:The text above is taken from the wikipedia entry ^[1].

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[1] <http://en.wikipedia.org/wiki/Thunder>

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